

9. Quantum Statistics, 3.5

12. Atomic Nucleus, 1

lecture 34, November 17, 2017

housekeeping

Coming attractions

Next week:

lecture M and T

chapter 12 homework due 11/22...HW workshop 11/21

Wed before Thanksgiving is a lost cause, so the schedule always indicated no class

End game:

watch the schedule link as I'll do some adjustment

exam #3 is Friday, December 1

I've not given any quizzes...have you noticed? I'll add that percentage to the homework fraction



today

statistical physics - Bose Condensation

Atomic nucleus - introduction



still MORE BOSONS Helium

Consider now a gas of bosons \rightarrow not spin 1, so 1 dof.

$$g(E)dE = \frac{8\sqrt{2} \pi m^{3/2}}{h^3} E^{1/2} dE \quad \dots \text{but take back that factor of 2}$$

\downarrow

$$g_B(E)dE = \frac{4\sqrt{2} \pi m_B^{3/2}}{h^3} E^{1/2} dE$$

$$n_B(E)dE = g_B(E) F_{BE}(E) dE$$

$$F_{BE}(E) = \frac{1}{B e^{E/\mu} - 1}$$

\swarrow
 $e^{-\mu/kT}$

$$N = \int_0^{\infty} n(E) dE = V \cdot \frac{4\sqrt{2} \pi m_B^{3/2}}{h^3} \int_0^{\infty} \frac{E^{1/2} dE}{e^{(E-\mu)/kT} - 1}$$

$$N = \int_0^{\infty} n(E) dE = V \cdot \frac{4\sqrt{2} \pi m_B^{3/2}}{h^3} \int_0^{\infty} \frac{E^{1/2} dE}{e^{(E-\mu)/kT} - 1}$$

N can't be negative!

$e^{(E-\mu)/kT}$ must be > 0 -- E is > 0 -- so $(T) < 0 \exists T$

change variables $x = E/kT$ \neq let $T \downarrow$

$$N = V \cdot \frac{4\sqrt{2} \pi m^{3/2} h^{3/2} T^{3/2}}{h^3} \int_0^{\infty} \frac{x^{1/2} dx}{e^{x-\mu/kT} - 1}$$

↑
constant as
 $T \downarrow$

↑
goes ↓ as $T \downarrow$

↓
must ↑ as $T \downarrow \Rightarrow |\mu(T)| \downarrow$
 $\Rightarrow \mu(T) \rightarrow 0$ as $T \downarrow$

But: problem.

At some point

$\mu(T) = 0 \rightarrow$ at $T_c \leftarrow$ "critical"

$$N = V \cdot \frac{4\sqrt{2}}{h^3} \pi m^{3/2} h^{3/2} T^{3/2} \int_0^{\infty} \frac{x^{3/2}}{e^{x-\mu/hT} - 1} dx$$

lnho: $\int_0^{\infty} \frac{x^{3/2}}{e^x - 1} dx = 2.315 \dots$ ie, when $\mu(T_c) = 0$

$$N = \frac{V \cdot 4\sqrt{2}}{h^3} \pi m^{3/2} h^{3/2} T^{3/2} (2.315)$$

Problem:

↑
cavit contracts $T \downarrow$

↑
cavit compensate to keep N constant

Below T_c ? Trouble. Einstein to the rescue in 1924

Einstein's way out: separate out the ground state, $E=0$

$$N = \int_0^{\infty} n(E) dE = V \cdot \frac{4\sqrt{2}}{h^3} \pi m_B^{3/2} \int_0^{\infty} \frac{E^{1/2} dE}{e^{(E-\mu)/kT} - 1}$$

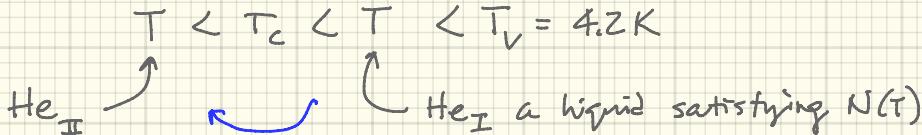
} only applies for particles not in g.s.: $E > 0$, not $E=0$

Below $T_c \dots$

$$N = N_0 + N_n$$

↑ ground state of He_{II}
↑ normal He: He_I

as $T \downarrow \dots$ N_n is depleted and N_0 increases.



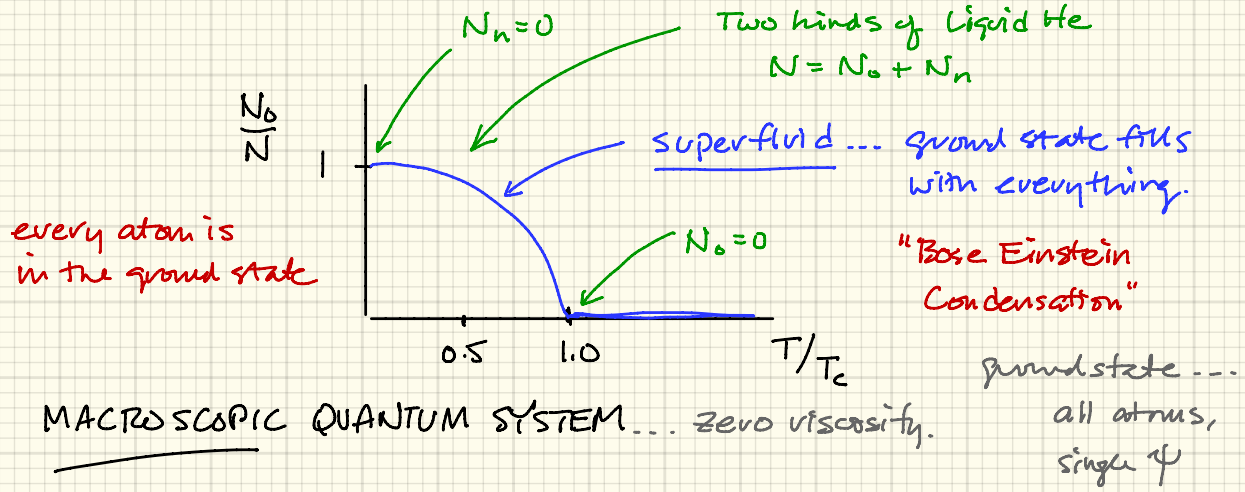
something different

phase transition

$$N = \frac{V \cdot 4\sqrt{2} \pi m^{3/2} h^{3/2} T^{3/2} (2.315)^3}{h^3} \quad \leftarrow \text{can solve for } T_c$$

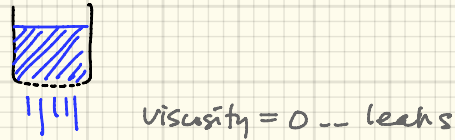
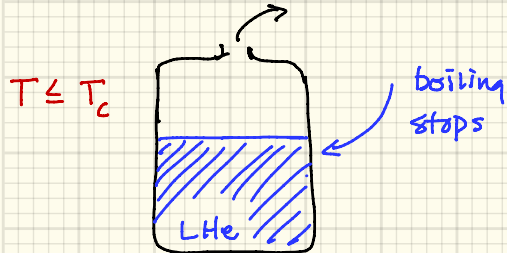
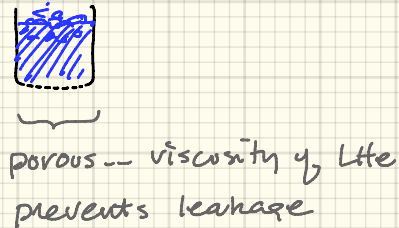
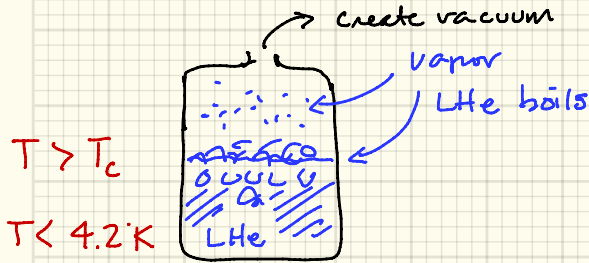
$$M = m_{\text{He}} = 2m_p + 2m_n$$

$$T_c = \frac{h^2}{2m_{\text{He}}k} \left[\frac{N}{V} \frac{1}{2\pi (2.315)} \right]^{2/3} \sim 3\text{K} \quad (\text{actually } 2.17\text{K})$$



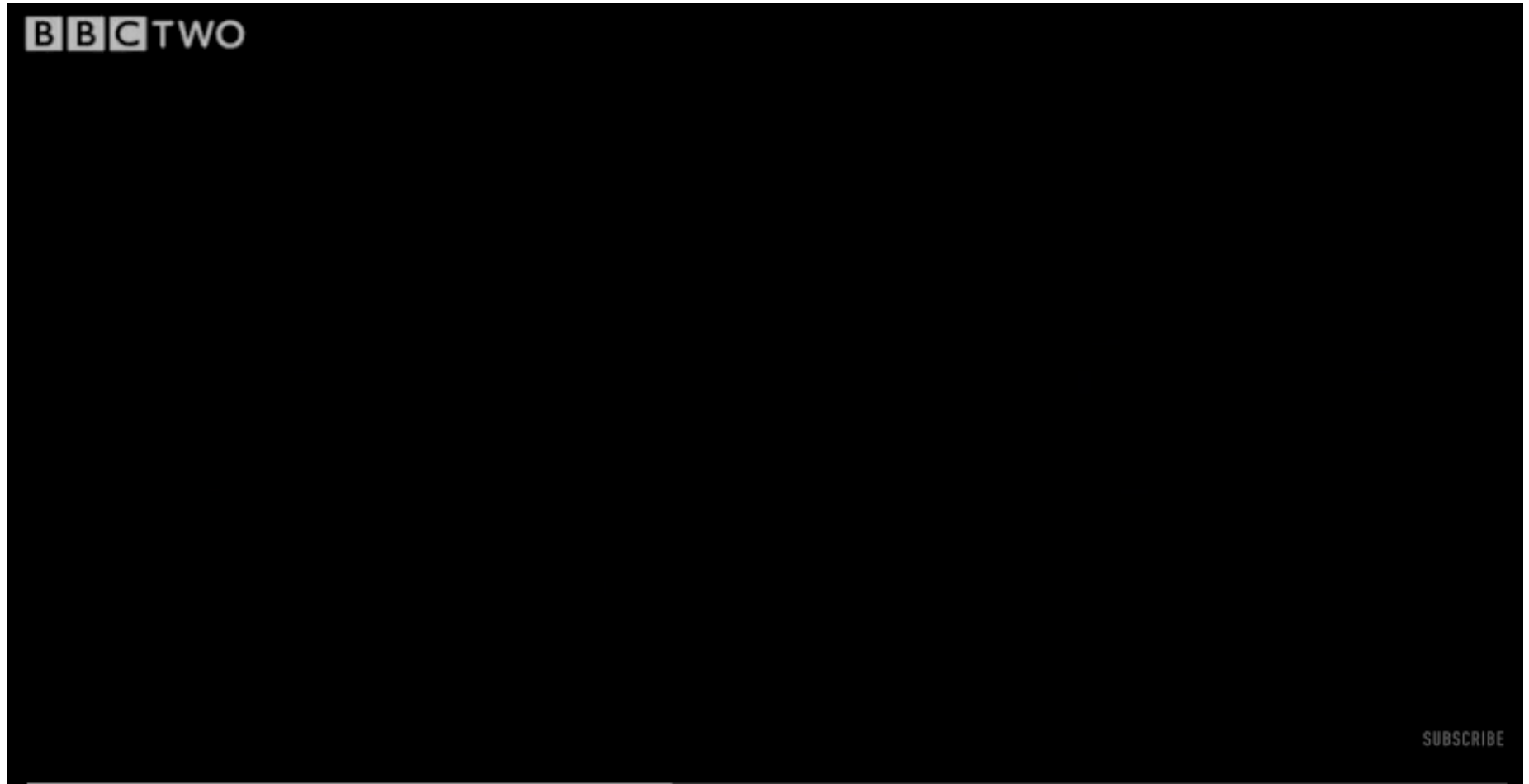
Liquid Helium condensate is --- unusual.

He atoms interact with one another --- closely spaced



movie →

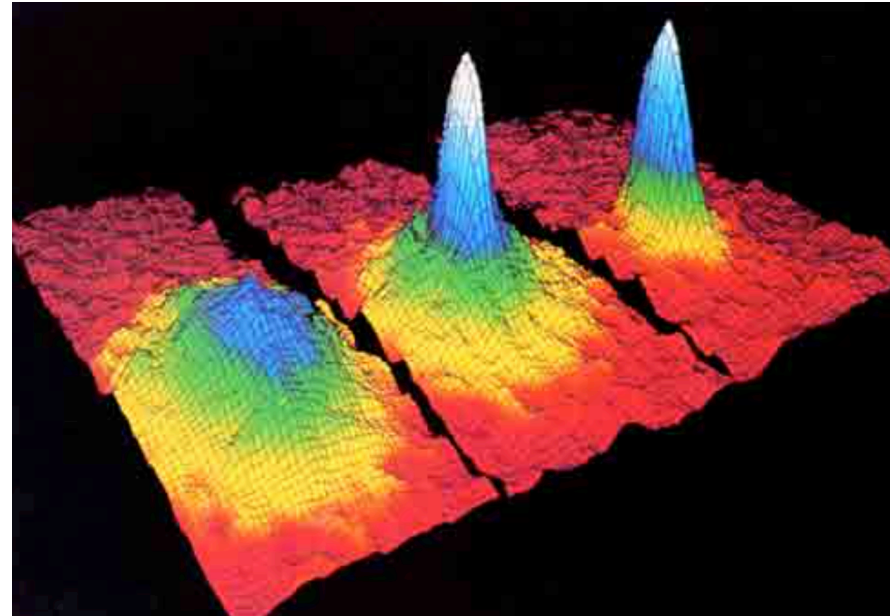
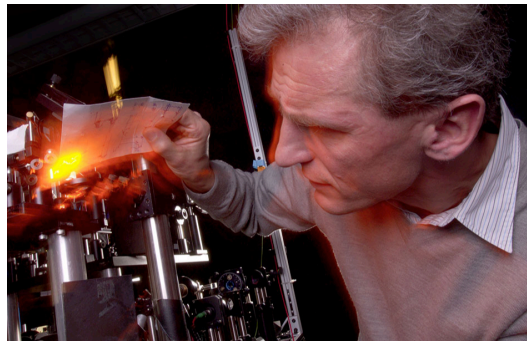
liquid Helium



<https://www.youtube.com/watch?v=9FudzqfpLLs>

BEC - 1995

Rb atoms...cooled to 50nK



nuclear physics hits:

- 1896 radioactivity discovered, Becquerel
- 1898 Radium/Polonium, Curies
- 1899 alpha/beta rays, Rutherford
- 1902 natural "transmutation," Rutherford and Soddy
- 1909 alpha rays = helium nuclei, Rutherford and Royds
- 1911 nucleus of + charge, Rutherford, Marsden, Geiger
- 1920 proton discovered, Rutherford
- 1921 neutron predicted, Rutherford
- 1921 laboratory transmutation, Rutherford and Chadwick
- 1929 proton accelerator, (Rutherford) Cockcroft and Walton
- 1930 neutrino predicted, Pauli
- 1931 Deuterium discovered, Urey
- 1932 neutron discovered, (Rutherford) Chadwick
- 1933 proton-neutron model of nucleus, Heisenberg
- 1933 artificial radioactivity, I. Curie and Joliot
- 1934 pion predicted, Yukawa
- 1935 neutrino model, Fermi
- 1938 nuclear fission, Hahn and Strassmann
- 1942 controlled nuclear fusion, Fermi
- 1947 pion and muon discovered, Occhialini and Powell



NUCLEAR PHYSICS

Just the facts, ma'am:

Nuclei \rightarrow neutrons & protons (to first approximation)

"protons" discovered and named by ... Rutherford, of course

"neutrons" predicted by and discovered in the lab of ... Rutherford ... of course

Notation

Z "Atomic number" = # protons in nucleus

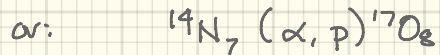
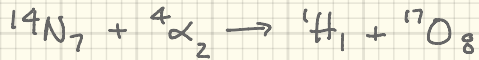
N "neutron number" = ... # of neutrons

A "mass number" = $Z + N$

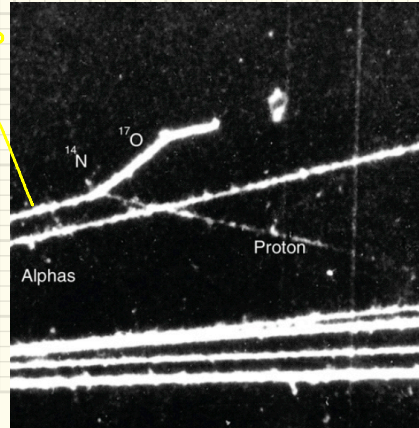
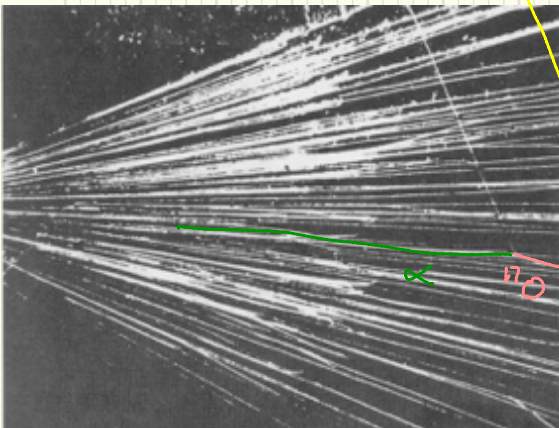
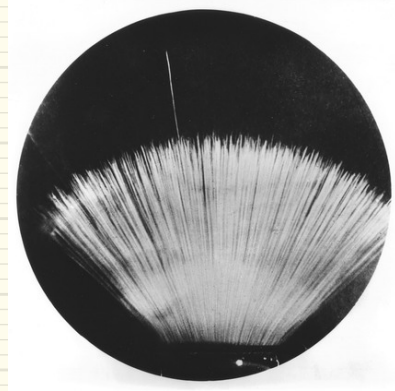
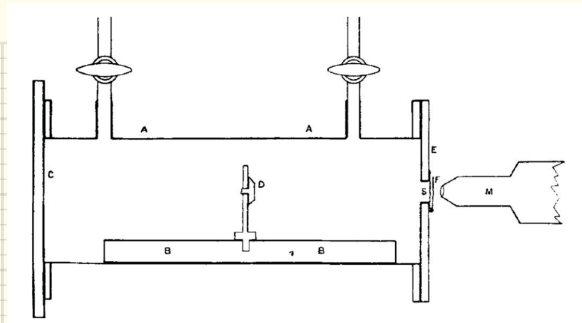
Symbol $\begin{matrix} A \\ Z \end{matrix} X_N$... often $\begin{matrix} A \\ X \end{matrix}$

eg: Nitrogen: $^{14}\text{N}_7$, Oxygen $^{16}\text{O}_8$, α $^4\alpha_2$ or $^4\text{He}_2$

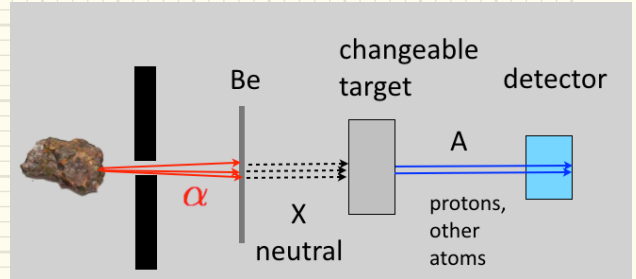
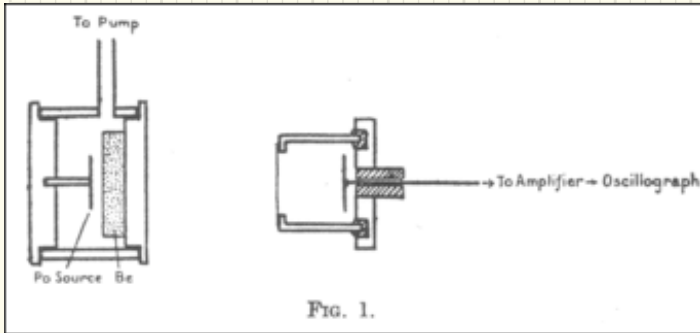
Rutherford's Discovery measurement



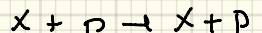
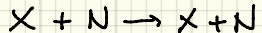
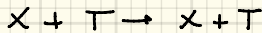
spoke of "H atoms" later named
 "proton" \Rightarrow "first one"
 speculated that a neutral neutron existed



Neutron - 1932



Chadwick compared



$$\left. \begin{array}{l} X + T \rightarrow X + T \\ X + N \rightarrow X + N \\ X + p \rightarrow X + p \end{array} \right\} m_T = 14 \text{ and } 1$$

max recoil velocity

$$v_T = \frac{2m_x}{m_x + m_T}$$

v_{X_0} ? so ratios:

$$\frac{v_H}{v_N} = \frac{m_x + 14}{m_x + 1} = \frac{3.3 \times 10^6 \text{ m/s}}{4.7 \times 10^6 \text{ m/s}} \Rightarrow m_x \approx 1.1 (x m_p)$$

actually found:

$$m_n = 938 \pm 1.8 \text{ MeV}/c^2$$

then

$$m_n = 939.57 \text{ MeV}/c^2$$

Now:

$$m_p = 938.27 \text{ MeV}/c^2$$

$$m_n = 939.57 \text{ MeV}/c^2 *$$

$$m_e = 0.511 \text{ MeV}/c^2$$

would:

$${}^1_1\text{p} = 1.007276 \text{ u}$$

$${}^1_0\text{n} = 1.008665 \text{ u}$$

$${}^0_{-1}\text{e}$$

but almost never.

Generically: "nucleon" = proton or neutron

"A" is the # of nucleons in a nucleus

Mass convention:

unified mass unit, u or "Atomic Mass Unit"

$$\text{defined for } {}^{12}\text{C} = 12 \text{ u exactly.} \Rightarrow 1 \text{ u} = 1.660559 \times 10^{-27} \text{ kg} \\ = 931.5 \text{ MeV}/c^2$$

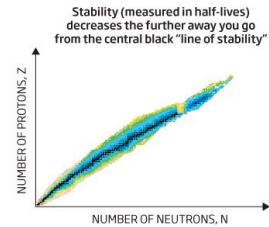
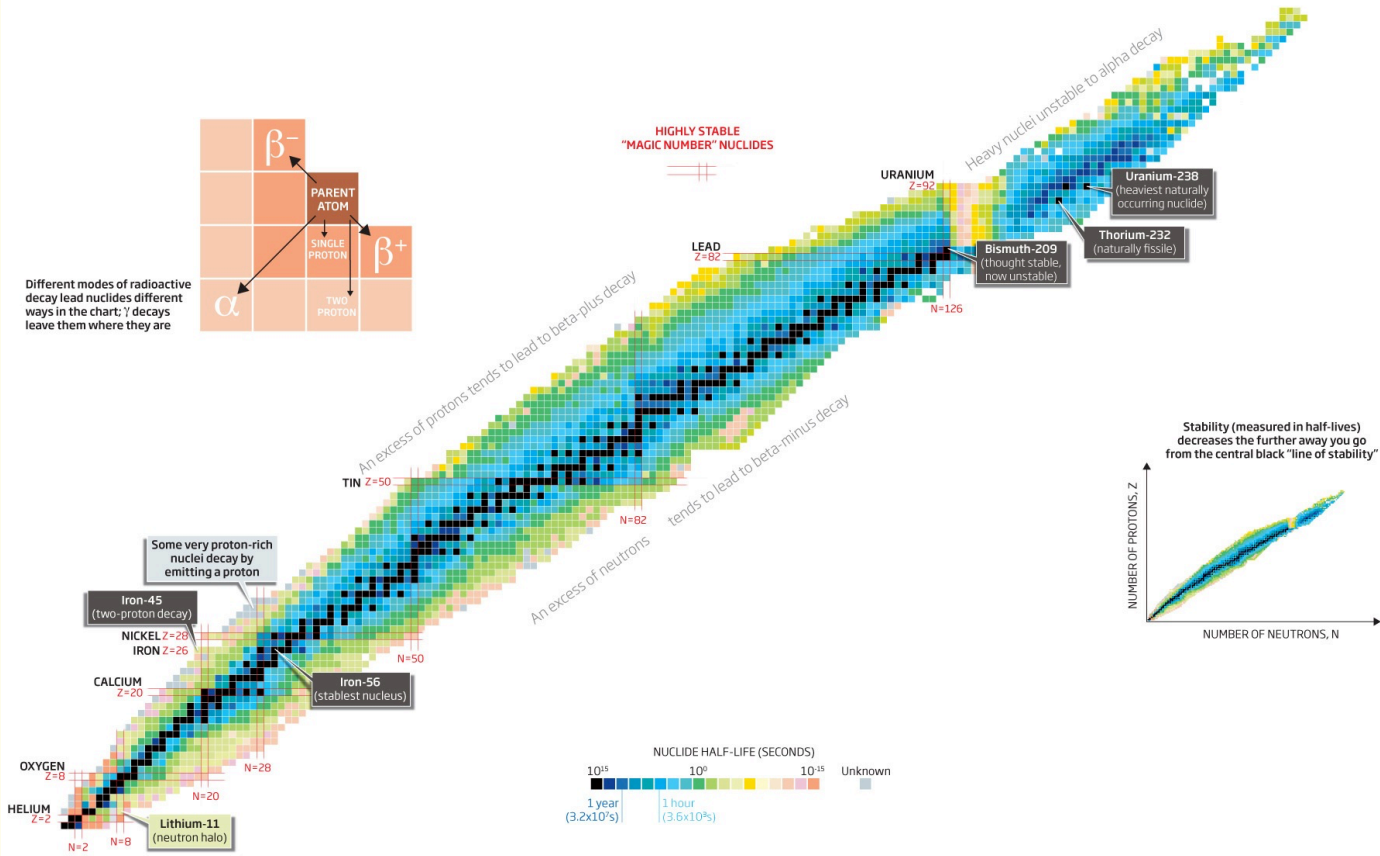
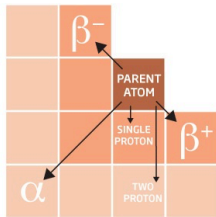
* took a while

Isotope: elements with different # neutrons & same # protons
(um... same element, huh)

Natural abundances vary

Carbon	${}^1_6\text{C}_5$	(= ${}^1_6\text{C} = {}^1_6\text{C}_6$)	trace %
	${}^{12}_6\text{C}_6$		98.69%
	${}^{13}_6\text{C}_7$		1.1%
	${}^{14}_6\text{C}_8$		trace %

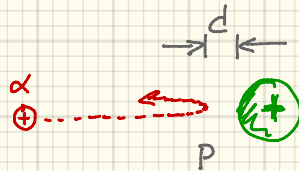
Different modes of radioactive decay lead nuclides different ways in the chart; γ decays leave them where they are



Stability (measured in half-lives) decreases the further away you go from the central black "line of stability"

SIZES AND SHAPES OF NUCLEI

Started with Rutherford & his model



@ turning point, P, $KE(\alpha) \rightarrow PE(\vec{E})$

$$\frac{1}{2} m v^2 = \frac{1}{4\pi\epsilon_0} \frac{Q_\alpha Q_{Au}}{r} = \frac{1}{4\pi\epsilon_0} \frac{e^2 Z_\alpha Z_{Au}}{d}$$

$$d = \frac{4}{4\pi\epsilon_0} \frac{Z_{Au} e^2}{m v^2} \sim 3.2 \times 10^{-14} \text{ m}$$

a maximum size for Au nucleus

For Ag... he found $d \sim 2 \times 10^{-14} \text{ m}$

New unit: femtometer (fm)

$$1 \text{ fm} = 10^{-15} \text{ m}$$



$$R = r_0 A^{1/3} \text{ — Empirical}$$

also called "Fermi" .. fm!

$$r_0 = 1.2 \times 10^{-15} \text{ m} = 1.2 \text{ fm}$$

NUCLEAR DENSITIES

Assume spherical... nuclear matter densities:

$$V = \frac{4}{3}\pi r^3 = \frac{4}{3}\pi r_0^3 A \quad \text{so } V \propto A$$

$$m_N = m_p \approx m_n$$

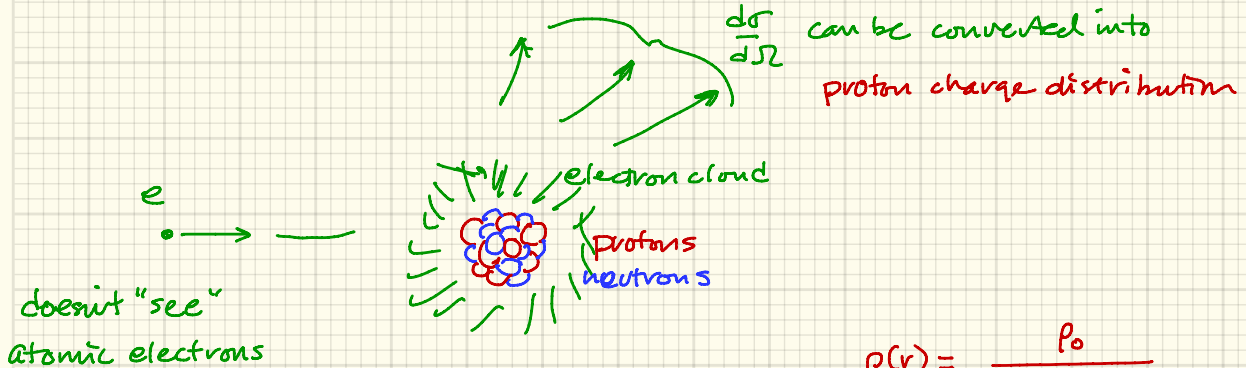
$$\rho = \frac{M}{V} = \frac{Am_N}{\frac{4}{3}\pi r_0^3 A} = \frac{3m_N}{4\pi r_0^3} \sim 2.3 \times 10^{17} \text{ kg/m}^3$$

$2 \times 10^{14} \times \rho(\text{water})$

not bad: density of neutron star $\sim 10^8 \text{ kg/m}^3$

Sophistication of electron beam production

→ long program at Stanford "size" is not a clear and distinct concept.
 would "map" the charge distribution of nuclei.

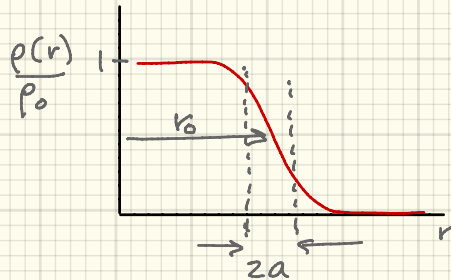


$$\rho(r) = \frac{\rho_0}{(1 + e^{r-r_0/a})}$$

$$r_0 \approx 1.07 A^{1/3} \text{ fm}$$

$$2a \approx 1 \text{ fm}$$

pretty good for $Z > 20$



Looks like F_{FD}
 doesn't it

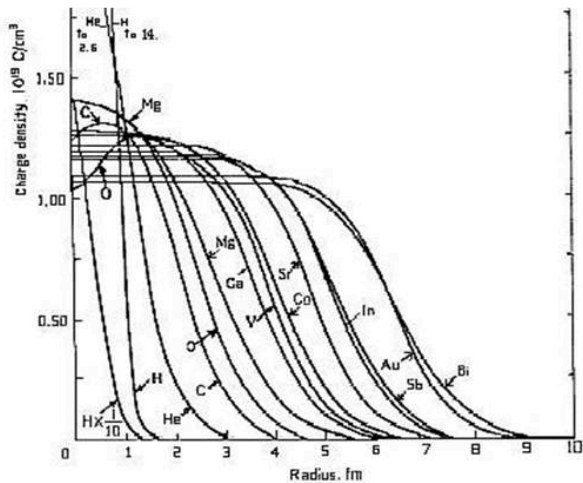
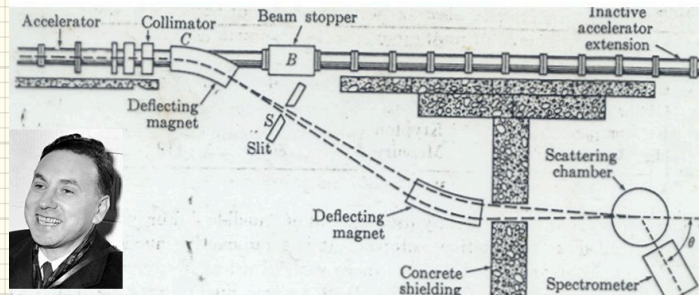
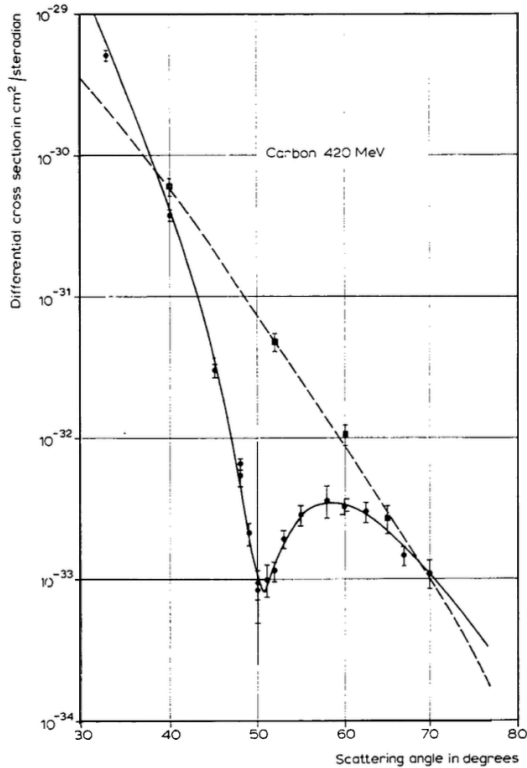


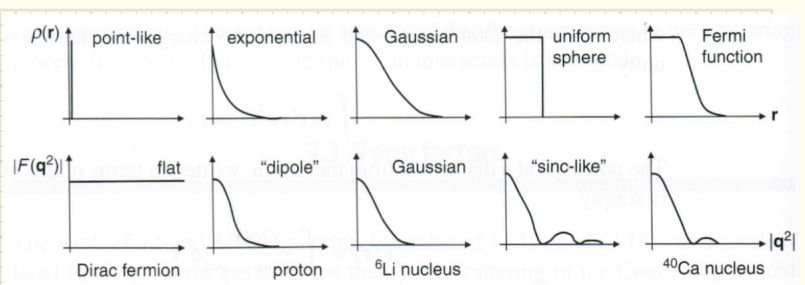
Figure 3- Charge distributions for various nuclei as determined from electron scattering experiments. [From R. Hofstadter, *Ann. Rev. Nucl. Sci.*, 7, 291 (1957).]

Electron scattering at Stanford 1954 - 57





↑ electron diffraction
 θ_1



$$\sin \theta_1 \approx \frac{1.22 \lambda}{D} \leftarrow \text{here, } \lambda_e$$

$D \leftarrow \text{diameter}$

$$D = \frac{1.22 \lambda_e}{\sin \theta_1} \quad \text{for } 450 \text{ MeV } e^-s$$

$$\lambda_e \approx 3 \times 10^{-15} \text{ m}$$

$$D = \frac{(1.22)(3 \times 10^{-15})}{\sin 50^\circ} = 4.8 \times 10^{-15} \text{ m}$$

$$\text{so } r_c = 2.4 \times 10^{-15} \text{ m}$$

$$= 2.4 \text{ fm}$$

