1. Special Relativity, 2 lecture 2, September 1, 2017

housekeeping

remember to check the course page:

chipbrock.org

and sign up for the feedburner reminders

I'll post something tonight...tell me if you don't get a message on Tuesday





Maxwell's Equations + 1

1.
$$\oint_{S} \vec{E} \cdot d\vec{S} = \frac{Q}{\epsilon_0}$$

Gauss' Law for electrostatics

$$\textbf{2.} \quad \oint_{S} \vec{B} \cdot d\vec{S} = 0$$

3.
$$\oint_{\ell} \vec{E} \cdot d\vec{\ell} = \mu_0 I + \mu_0 \epsilon_0 \frac{d\Phi_E}{dt}$$

Ampere's Law

$$4. \quad \oint_{\ell} \vec{B} \cdot d\vec{\ell} = -\frac{d\Phi_E}{dt}$$

Faraday's Law

Lorentz Force

 $\epsilon_0 = 8.854 \times 10^{-12} \text{ C}^2 \text{N}^{-1} \text{m}^{-2}$ $\mu_0 = 4\pi \times 10^{-7} \text{ JA}^{-2} \text{m}^{-1}$

5. $\vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$

vacuum permittivity vacuum permeability

$$\frac{1}{4\pi\epsilon_0} = k_e = 8.9 \times 10^9 \ \rm Nm^2 C^{-2}$$

tracking

theory and experiment circa 1900



differential form of Maxwell's Equations

In general:
$$\vec{E} = E_x \hat{i} + E_y \hat{j} + E_y \hat{h}$$

and no
 $\vec{E} = -\left(\frac{\partial V}{\partial x} \hat{i} + \frac{\partial V}{\partial y} \hat{j} + \frac{\partial V}{\partial z} \hat{h}\right)$
Standard calculus operator: gradient
 $\vec{\nabla} = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{h} \frac{\partial}{\partial z}$
 $h_0 = \vec{E} = \vec{\nabla} V$
Also... divergence
 $\vec{\nabla} \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$
and... $cvrl$
 $\vec{\nabla} \times \vec{A} = \hat{i} \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_z}{\partial z}\right) + \hat{j} \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x}\right)$
 $+ \hat{h} \left(\frac{\partial A_1}{\partial x} - \frac{\partial A_x}{\partial y}\right)$

2 Take a truy volume C×,7,3 5 Imagine an E(x, y, 3) C bach tace Ø ds = - î dydz dy which changes at surface vectors the front face point out from X -> X + dx $dS = + \lambda dy dy$ $\vec{E}(x,y,z) + \left(\frac{\partial \vec{E}}{\partial x}\right) dx$ bach change ÷

hooh at the flax through the volume \$ E.ds due to these two faces. $\int \vec{E} \cdot d\vec{S} = \vec{E} \cdot (-\hat{\lambda} \, dy \, dz) + (\vec{E} + \frac{\partial \vec{E}}{\partial x} \, dx) \cdot (\hat{\lambda} \, dy \, dz)$ bach, trant cancel = $dxdydz = \frac{1}{2} (\vec{E} \cdot \hat{z})$ other faces? = dxdydz DEx Dx similar. $\phi \vec{E} \cdot d\vec{s} = dx dy dz \left(\begin{array}{c} \partial E_x \\ \partial X \end{array} \right) = dx dy dz \left(\begin{array}{c} \partial E_x \\ \partial X \end{array} \right) = \frac{\partial x dy dz}{\partial y} \left(\overrightarrow{\nabla} \cdot \overrightarrow{E} \right)$ DEZ

Suppose then had been a charge inside? 9 = p dxdydz [†] chny-density OK: $\oint \vec{E} \cdot \vec{dS} = 9$ dx dy dz (₹.Ē) = dxdydzp 6. $\overrightarrow{\nabla}$. $\overrightarrow{E} = \underbrace{f}_{E_{o}}$ Maxwell's Equation for Gouss' Low

$$\vec{\nabla} \cdot \vec{E} = \int_{c_0}^{c_0} \vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \cdot \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \cdot c_0 \cdot \frac{\partial \vec{E}}{\partial t} + \mu_0 \cdot \vec{J}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \cdot c_0 \cdot \frac{\partial \vec{E}}{\partial t} + \mu_0 \cdot \vec{J}$$

FREE SPACE $\rightarrow \infty \text{ sources.}$

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$\vec{\nabla} \cdot \vec{E} = -\frac{\partial \vec{B}}{\partial t} \longrightarrow \text{ thue changing } \vec{B}? \text{ get } \vec{L} \cdot \vec{E}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \cdot c_0 \cdot \vec{2E} \longrightarrow \text{ thue changing } \vec{E}? \text{ get } \vec{L} \cdot \vec{B}$$

FxB = Moe. DE pith out a partiallar component: X $\frac{\partial B_3}{\partial y} = \frac{\partial B_1}{\partial 3} = \mu_0 \epsilon_0 \partial E_X$ remember this now: $\vec{\nabla} \times \vec{E} = -\vec{\partial} \vec{B} \longrightarrow \vec{\nabla} \times \vec{\nabla} \times \vec{E} = \vec{\nabla} \times \left(-\vec{\partial} \vec{B} \right)$ anidentity $= \overline{\nabla} \left(\overline{\nabla} \cdot \overline{E} \right) - \nabla^{2} \overline{E}$ $= O \quad Gruss'$ Low : $\nabla^2 \vec{E} = \frac{2}{2\star} (\vec{\nabla} \times \vec{E})$ $\vec{\nabla}^{2}\vec{E} = p_{0} \epsilon_{0} \vec{\delta}\vec{E} \longrightarrow \vec{\nabla}^{2}\vec{E} - p_{0} \epsilon_{0} \vec{\delta}\vec{E} = 0$ $(iii) \quad \vec{\delta}^{2}\vec{E} + \vec{\delta}^{2}\vec{E} + \vec{\delta}^{2}\vec{E} \qquad \text{another for } \vec{B} \text{ also}$ $\vec{\delta}_{X^{2}} + \vec{\delta}_{Y^{1}}^{2} + \vec{\delta}_{Y^{2}}^{2}$

$$\nabla^{2}\vec{E} - \mu \circ \vec{G} \stackrel{\vec{\partial}\vec{E}}{\partial t^{2}} = 0 \qquad \text{loch at} \qquad \frac{\partial^{2}\vec{\Phi}}{\partial x^{2}} - \frac{1}{u^{2}} \frac{\partial^{2}\vec{\Phi}}{\partial t^{2}} = 0$$
another for \vec{B} also $\rightarrow a$ work equation for ϕ with velocity v .
 $\vec{\Phi}$
 a wave equation for \vec{E}
with velocity $v = \sqrt{\frac{1}{\mu \circ 6}}$
ONE OF THOSE WOW MOMENTS IN THEORETICAL PHYSICS
 $\mu = 4\pi \times 10^{7} \text{ N} \cdot \text{C}^{2} \text{ s}^{2}$
 $G_{0} = 8.8542 \times 10^{12} \text{ C}^{2} \text{N}^{-1} \text{m}^{-2}$
So $\sqrt{\frac{1}{\mu \circ 6}} = \sqrt{\frac{1}{(4\pi \times 10^{7} \text{ N} \text{ s}^{2})(8.8542 \times 10^{12} \text{ c}^{2} \text{ m})}$
 $v = 2.948 \times 10^{8} \text{ m/s} = c$
and $|\vec{B}| = |\vec{E}|/c$

$$E(x,b)$$

$$E(x,b)$$

$$E(x,b)$$

$$E(x,b)$$

$$E(x,v,t)$$

Summary 1.
$$\oint \vec{E} \cdot d\vec{S} = \frac{q}{E_0}$$

 $Z \cdot \oint \vec{E} \cdot d\vec{S} = 0$
 $3. \oint \vec{E} \cdot d\vec{L} = poI + poEo d\vec{E}_{\vec{E}}$
 $4. \oint \vec{E} \cdot d\vec{L} = -d\vec{E}_{\vec{E}}$
These are for "both" regions of space and extended sources..."Integral form"
Useful are versions which are infinitesimal \rightarrow " differential form"
Useful: potential
 $E_L = -\frac{dV}{dL}$ $L = x.y. \text{ or } z$
 $V = V(x, y, z)$
Now : experiments on light.

1. Speed of light 1676 Ole Romer 93 Mmiles Ē S EN 9 here? here? clock with measured measured reflecting light to E TT+11 min $C_{T} = 11$ min knew how long zz min Pn light for IO to to cross Earth eclipse Narton or Huggens -> 16 min total $C = \frac{93 \times 10^6}{8 \min} \left(\frac{1 \min}{60 \text{ s}} \right) = 200, \text{mm} \text{mi/s} \sim 186, \text{growmi/s}$



Analogy: you say the star is the direction of telescope light : vain earth orbit: trach 4 telescope: hat \mathcal{D} Jan July January July going <going ---> 20mph James Bradley 1728 > 5mph Observed angle =) earth maing through ether ymsee 5 you could ETHER determine

Fizean's Ether-Drag Experiment 2. Motion relative to the ether : Speed with current speed of C IS CW hight m against curvent stationory **谷**口 日 water $C_W = \frac{c}{n} + fv$ if ether dragged $C_A = C_A - fv$ by water ._ $f \neq 1$ f \$0 H. Fizeau 1851: f ≠ I (=) some drag of Ether by water ? Ś ETHER









1887

Albert Michelson (1852-1931)

and

Edward Morley



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The Nobel Prize in Physics 1907



Albert Abraham Michelson Prize share: 1/1

The Nobel Prize in Physics 1907 was awarded to Albert A. Michelson "for his optical precision instruments and the spectroscopic and metrological investigations carried out with their aid".

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style: "The Nobel Prize in Physics 1907". *Nobelprize.org.* Nobel Media AB 2014. Web. 28 http://www.nobelprize.org/nobel_prizes/physics/laureates/1907/> The prize went to Michelson for the instrument: the Michelson interferometer

the same idea as the gravitational wave instrument

We remember him for the most important null measurement in 200 years

"Michelson Morley

trying to measure the speed of Earth relative to Ether

measure the fringes in light interfering from the two paths...then rotate the instrument 90 degrees - and do it again.

The differences between the two configurations is related to the time difference



Experiments"

This technique was perfected by cowboy, Albert Michelson and eventually his sidekick, Edward Morley at Case Western Reserve in Cleveland between 1880 and 1888



If the beams get back out of phase...one traveled through the ether differently from the other.

neat simulation

http://www.kcvs.ca/site/projects/physics_files/specialRelativity/michelsonMorley/mmExperiment.swf





С water, u U, ether C, hight · swimmer, v , u = speed of ether wind V=C ZLZ Ξ $\frac{2L_{1}}{c} - \frac{1}{\mu^{2}} \frac{1}{\mu^{2}}$ 6 С 24 $t_2 =$ -t. ZLz C 1 - u2/c2 W2/22 $v_e = 3 \times 10^4 \, \text{m/s}$ C= 3×10° m/s { Ue/c~104 Taylor Expans $1 + \frac{1}{2} \frac{1}{2}$ $\frac{2}{c}$ 1 + $\frac{1}{c}$ + . - . + <mark>2</mark>). $) - \frac{2L_2}{2} \left(1 + \frac{1}{2} \frac{u^2}{2^2} \right)$ $\equiv \Delta t(o)$ ting

$$\Delta t(a) = \frac{2L}{C} \left(1 + \frac{w}{C}\right) - \frac{2L_2}{C} \left(1 + \frac{1}{2} \frac{w}{C^2}\right)$$

$$= \frac{2L_1}{C} - \frac{2L_2}{C} + \frac{2L_1}{C^3} - \frac{L_2}{C^3} \frac{w^2}{C^3}$$

$$\Delta t(a) = \frac{2}{C} \left(L_1 - L_2\right) + \frac{w^2}{C^3} \left(2L_1 - L_2\right) = t_1(a) - t_2(a)$$
how be clever hard to measure precisely.

Rotate
$$Even + \frac{1}{2} = \frac{1$$

$$\Delta t(b) = \frac{2}{c} (L_1 - L_2) + \frac{u^2}{c^3} (2L_1 - L_2) = t_1(b) - t_2(b)$$

$$\Delta t(W_2) = \frac{1}{c} (W_2) - \frac{1}{c^3} (W_2)$$

$$= \frac{2}{c} (L_1 - L_2) + \frac{u^2}{c^3} (L_1 - 2L_2)$$

$$\Delta T = \Delta t(b) - \Delta t(W_2) \qquad !$$

$$= \frac{2}{c} (L_1 - L_2) + \frac{u^2}{c^3} (2L_1 - L_2) - \frac{2}{c} (L_1 - L_2) - \frac{u^2}{c^3} (L_1 - 2L_2)$$

$$= \frac{u^2}{c^3} (2L_1 - L_1 - L_2 + 2L_2)$$

$$\Delta T = \frac{u^2}{c^3} (L_1 + L_2) \qquad \rightarrow s \qquad c = \frac{5}{\delta T}$$

$$S = c \Delta T$$

$$S = \frac{u^2}{c^2} (L_1 + L_2) \checkmark$$