## hi

Day 27, 19.04.2018
Particle Physics 2

## housekeeping

The end game: next slide Particle Physics:

Readings: Oerter and Hobson


Hobson_PP.pdf is chapter 17 out of Hobson
Homework \#12 is all from MasteringPhysics - normal due date Feynman Diagram rules

3 movies in the lecture slide directory - you'll need them for homework and the final
they are: primitiveDiagrams_X.mp4 where $X=0,1,2$

## last 2 weeks \& final

Homework \#13 will be assigned 4/21 and due 4/28 - normal rotation
On-line final exam will be assigned Sunday, 4/29 and due Tuesday night, May 1
will cover material since midterm plus the last week of class
There is 1 more 10 point quiz (stay tuned)...
only the shadow knows when
Remember when I was sick?
been trying to catch up, but not going to make it. Hence:
Final Exam day:

1. You'll arrive at 0745 on May 4, here. I know.
2. I'll provide bagels. You supply liquids.
3. We'll have a quiz.
4. I'll finish with about a 1 hour grand finale, Ialapalooza, mind-bending lecture
5. You'll do your Feynman Diagram Project
6. There will be no poster project this year


## honors project began

https://qstbb.pa.msu.edu/storage/Homework_Projects/honors_project_2018/
contains:
the first instructions: the plan \& tutorial
the second instructions - v2 uploaded, added a missing student
the data, assigned by name in the second instructions - see next

## dates:

complete first part, March 16
analyze data by April 24 and hand in complete writeup at the final exam

## the data

## should have been in zipped format

## rather, somehow they were unzipped in some process

fixed: now
https://qstbb.pa.msu.edu/ storage/Homework_Projects/ honors_project_2018/

|  | Last modified | Size |  |
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|  |  |  |  |

## I need a Section 2

to test the Z-path uploading machinery and instructions

## particles in time

An anti-electron...coming into an initial state to a node:


Yes, this makes sense
is the same thing as
An electron coming out of an initial state (?)


Nope, this makes no sense...time-backwards

An anti-electron...coming out of a final state:


Yes, this makes sense
is the same thing as
An electron coming into a final state (?)


Nope, this makes no
sense...time-backwards
primitive

## diagrams

## are general

but this is completely general...for any charged fermion:

$f$ could be electron, positron, proton, antiproton...and more - any electrically charged fermion.

Their diagrams are identical.

## Primitive Diagram Scorecard

your first entry

Primitive Diagrams


## beta decay

the inaugural non-QED interaction
Weak Force

## Fermi

## Theory of

## Beta

## Decay

## uses the Dirac

 ideas of quantum electrodynamics
## particle creation and annihilation



Fig. 5. Energy distribution curve of the beta-rays.
$\mathrm{m}_{\text {neutron }}>\mathrm{m}_{\text {proton }}$
a smidgen.

a free neutron has a lifetime of about 11 minutes. He sent the paper to Nature, but it was rejected:
"it contained speculations which were too remote from reality"

from his original paper for different nuclear species parameters

# exchange force 

the modern view:
if there's a force...there's a field
if there's a field...there's a particle

## we know

## one

## force..so

## far

## electromagnetism

 electricitymagnetism
united by Relativity remember?


The modern idea:
The force of electromagnetism is "propagated" by the photon.

Multiple names: "propogator"
"Intermediate Vector Boson"

I'll call the photon: the "Messenger Field for Electromagnetism"

## charge independence

the force that holds the protons and neutrons together
is the same between $n-n, p-p, n-p$
Strong Force
but only over a very short range...
the STRONG force
overwhelms the electromagnetic force
uncertainty
certainly
to the
rescue
brilliant
observation by
Yukawa
maybe there's a quantum that is active only over the size of a nucleus: "U"
another exchange force/particle?

So: $p \rightarrow n+U$ ?


Suppose U travels at c within a nucleus... $\Delta t=\Delta x / c$
Then Uncertainty could estimate U's mass... $\Delta E \Delta t=h / 4 \pi$

$$
m_{U}=\Delta E / c^{2}
$$

$$
m_{U} \stackrel{?}{\approx} 100 \times 10^{6} \mathrm{eV}=100 \mathrm{MeV}
$$

## the

## Yukawa



## particle

## is the pion



These coupling strengths are large - strong.

In technical terms we call this...the strong interaction.

## If we ignore electromagnetism...the proton \& the neutron are

 very much alike - we can treat them as being the same particle
## neutrons

## and

protons
act like they are identical particles
the electric charge?
as a force...Yukawa's force is 100 times the electromagnetic

For nuclear forces: treat p and n as identical and differing only by a "quantum number" called "Isospin"
$I$

$$
\begin{aligned}
& \text { ( } \\
& \mathbf{N} \\
& \text { "nucleon" }
\end{aligned}
$$

A neutron... is a "nucleon" with "isospin down" A proton... is a "nucleon" with "isospin up"

They go together...within the strong, nuclear force.
How?
refers to:
entomology:
example:
either a proton or a neutron
from "nucleus"...the "-on" tends to be a particle name
"nucleon force"
strange things in cosmic rays
thick photographic substrates


## by 1950 the forces were identified

"strong"
as evidenced by the pion (refined later)
"electromagnetic"
as evidenced by the exchange of photons among electrically charged particles
"weak"
as originally evidenced by neutron beta decay, and subsequently pion, muon, and other hadronic decays
"gravitational"
the weakest of all...quantum theory of gravity still a mystery
three
forces now of vastly different strengths

Electromagnetic force 0.007


Weak force 0.000001


## FAMILIES

Nature prefers
like-particles


Lepton
Families
electrons and a neutrino
muons and a neutrino


## These sorts of patterns are a huge deal.

$$
\begin{array}{r}
\mathrm{Q} \\
0 \\
-1
\end{array} \quad\binom{\nu_{e}}{e} \quad\binom{\nu_{\mu}}{\mu} \quad\binom{\nu_{\tau}}{\tau}
$$




#### Abstract








taus and a neutrino

(2) -

$\square$

## by 1955




## 100's of them

## things wer

me what's so "elementary" about that?


## hadron

refers to:
any particle that interacts via the Strong Force
entomology:
example:
$\alpha \bar{\rho} \rho o ́ \sigma$ "hadros" "large", "massive"
proton and neutron not electron, not photon
jargon alert: lepton
refers to:
entomology:
example:
originally, an electron, muon, neutrino
" $\lambda \varepsilon \pi t$ tós" (leptos), "fine, small, thin"
electron, muon, neutrino, tau!

## The Particle Zoo?


jargon alert: particle quantum numbers
refers to:
quantities that are inherently a part of particles, which are conserved in interactions or decays
entomology:
example:
historical to Bohr and Schroedinger
electric charge, baryon number, lepton number, isospin
something like these will never happen:
so, you'll always see:
total electric charge at the beginning equals total charges at the end


## Strangeness,

## S

## strangeness

 seems to come in pairs
## assign "strangeness" empirically.



$$
\pi^{-}+p \rightarrow \Lambda^{0}+K^{0}
$$

$$
\begin{array}{lllll}
\mathrm{S}: & 0 & 0 & -1 & +1
\end{array}
$$

Strong interaction


$$
\begin{array}{lllll}
\mathrm{S}: & 0 & 0 & -1 & 0
\end{array}
$$

and yet you do see:

$$
\Lambda \rightarrow p+\pi^{-}
$$

$$
\begin{array}{llll}
\mathrm{S}: & -1 & 0 & 0
\end{array}
$$

Weak
interaction

Production of a subset of all baryons seems to require them to come in pairs.
Strong interactions conserve Strangeness
Decay of those same baryons...notsomuch Weak interactions change Strangeness by 1 unit

## the dominant Baryons

$\left.\begin{array}{|c|c|c|c|c|c|c|c|}\hline \text { Particle } & \text { Symbol } & \begin{array}{c}\text { Rest Mass } \\ \mathrm{MeV} / \mathbf{c}^{2}\end{array} & \text { spin } & \mathbf{Q} & \mathbf{B} & \mathbf{s} & \text { Lifetime }\end{array} \begin{array}{c}\text { dominant decay } \\ \text { modes }\end{array}\right]$

## the dominant Mesons

| Particle | Symbol | antiparticle | Rest <br> Mass <br> $\mathrm{MeV} / \mathrm{c}^{2}$ | spin | Q | B | S | Lifetime | dominant decay modes |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Pion | $\pi^{+}$ | $\pi^{-}$ | 139.6 | 0 | +1 | 0 | 0 | $2.6 \times 10^{-8}$ | $\mu^{+} \nu_{\mu}$ |
| Pi-zero | $\pi^{0}$ | $\pi^{0}$ | 135 | 0 | 0 | 0 | 0 | 920 | $2 \gamma$ |
| Kaon | $K^{+}$ | $K^{-}$ | 493.7 | 0 | +1 | 0 | +1 | $1.24 \times 10^{-8}$ | $\mu^{+} \nu_{\mu}, \pi^{+} \pi^{0}$ |
| K-short | $K_{S}^{0}$ | $K_{S}^{0}$ | 497.7 | 0 | 0 | 0 | +1 | $0.89 \times 10^{-10}$ | $\pi^{+} \pi^{-}, 2 \pi^{0}$ |
| K-long | $K_{L}^{0}$ | $K_{L}^{0}$ | 497.7 | 0 | 0 | 0 | +1 | $5.2 \times 10^{-8}$ | $\pi^{ \pm} \ell^{\mp} \nu_{\ell}$ |
| Eta | $\eta^{0}$ | $\eta^{0}$ | 548.8 | 0 | 0 | 0 | 0 | $<10^{-18}$ | $2 \gamma, \pi^{+} \pi^{-} \pi^{0}$ |
| Eta-prime | $\eta^{0 \prime}$ | $\eta^{0 \prime}$ | 958 | 1 | 0 | 0 | 0 | ... | $\pi^{+} \pi^{-} \eta$ |
| Rho | $\rho^{+}$ | $\rho^{-}$ | 770 | 1 | +1 | 0 | 0 | $0.4 \times 10^{-23}$ | $\pi^{+} \pi^{-}, 2 \pi^{0}$ |
| Rho-naught | $\rho^{0}$ | $\rho^{0}$ | 770 | 1 | 0 | 0 | 0 | $0.4 \times 10^{-23}$ | $\pi^{+} \pi^{-}$ |
| Omega | $\omega^{0}$ | $\omega^{0}$ | 782 | 1 | 0 | 0 | 0 | $0.8 \times 10^{-22}$ | $\pi^{+} \pi^{-} \pi^{0}$ |
| Phi | $\phi$ | $\phi$ | 1020 | 1 | 0 | 0 | 0 | $20 \times 10^{-23}$ | $K^{+} K^{-}, K^{0} \bar{K}^{0}$ |

## patterns emerged

to Murray Gell-Mann \& (independently) Yuval Ne’eman in 1964

$$
\begin{gathered}
\Delta^{0} \Xi^{-} \Delta^{-} \\
\Sigma^{-} \Lambda^{0} \Xi^{*} p \\
\Delta^{+} \Sigma^{0}{ }^{*} \Sigma^{*-} \\
\Sigma^{*+} \Delta^{++} n \\
\Xi^{*} \Sigma^{+} \Sigma^{* 0}
\end{gathered}
$$



## family arrangements



## quarks

the mathematical description of such patterns

## 1964



## Murray Gell-Mann

1929 -
theoretician
Nobel Laureate 1969

Yale at age of 15. PhD from MIT at age of 22.

Speaks at least 13 languages fluently. Studies linguistics now, among other things.

Unraveled many of the organization puzzles of the particle zoo:
strangeness
an empirical mass formula relating them

Worries a lot now about the nature of physical law.

A not-so-good TED lecture on mathematical Beauty in physics...link below.

Not known for his humility.

Gell-Mann found that the patterns work

Gell-Mann's original pattern for quarks. Changed...
if every particle is composed of smaller bits
with fractional electric charge:
charge of up quark:
charge of down quark: charge of strange quark:
$+2 / 3$ e
$-1 / 3$ e
$-1 / 3$ e

# fundamental particles, circa...now 

quarks and leptons
hadrons are composite: made of quarks
electrons and cousins are fundamental on their own

Baryons \& Mesons differ by quark-content Baryons are made of 3 quarks

Mesons are made of 1 quark and 1 antiquark

## Quarks

## 1964 version

## fundamental fermions

in same league as electrons and neutrinos

S


| Quark | Symbol | Rest <br> Mass <br> MeV/c | spin | Q | B | S |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| up | $u$ | $1.7-3.3$ | $1 / 2$ | $+2 / 3$ | $1 / 3$ | 0 |
| down | $d$ | $4.1-5.8$ | $1 / 2$ | $-1 / 3$ | $1 / 3$ | 0 |
| strange | $s$ | 101 | $1 / 2$ | $-1 / 3$ | $1 / 3$ | -1 |

## piece 'em together:

## proton

electric charge $=+1$

| Quark | Symbol | Rest <br> Mass <br> MeV/c | spin | Q | B | s |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| up | $u$ | $1.7-3.3$ | $1 / 2$ | $+2 / 3$ | $1 / 3$ | 0 |
| down | $d$ | $4.1-5.8$ | $1 / 2$ | $-1 / 3$ | $1 / 3$ | 0 |
| strange | $s$ | 101 | $1 / 2$ | $-1 / 3$ | $1 / 3$ | -1 |

## piece 'em together:

## proton electric charge $=+1$

| Quark | Symbol | Rest <br> Mass <br> MeV/c | spin | Q | B | s |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| up | $u$ | $1.7-3.3$ | $1 / 2$ | $+2 / 3$ | $1 / 3$ | 0 |
| down | $d$ | $4.1-5.8$ | $1 / 2$ | $-1 / 3$ | $1 / 3$ | 0 |
| strange | $s$ | 101 | $1 / 2$ | $-1 / 3$ | $1 / 3$ | -1 |

## piece 'em together:

## neutron

## electric charge $=0$

| Quark | Symbol | Rest <br> Mass <br> MeV/c | spin | Q | B | S |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| up | $u$ | $1.7-3.3$ | $1 / 2$ | $+2 / 3$ | $1 / 3$ | 0 |
| down | $d$ | $4.1-5.8$ | $1 / 2$ | $-1 / 3$ | $1 / 3$ | 0 |
| strange | $s$ | 101 | $1 / 2$ | $-1 / 3$ | $1 / 3$ | -1 |

## piece 'em together:

## neutron

## $+2 / 3$

electric charge $=0$
$-1 / 3$
$-1 / 3$

| Quark | Symbol | Rest <br> Mass <br> MeV/c | spin | Q | B | s |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| up | $u$ | $1.7-3.3$ | $1 / 2$ | $+2 / 3$ | $1 / 3$ | 0 |
| down | $d$ | $4.1-5.8$ | $1 / 2$ | $-1 / 3$ | $1 / 3$ | 0 |
| strange | $s$ | 101 | $1 / 2$ | $-1 / 3$ | $1 / 3$ | -1 |

they all fit

$\operatorname{spin} 3 / 2$

like a glove

S


## discovered at Brookhaven within a year

 the "Omega minus" was discovered at Brookhaven National Lab S

$$
-1-\frac{1}{2} \quad 0 \quad \frac{1}{2} \quad 1
$$

I


## most famous bubble chamber picture in history, 1964



FIG. 2. Photograph and line diagram of event showing dec

The event in question is shown in Fig. 2, and the pertinent measured quantities are given in Table I. Our interpretation of this event is

$$
\begin{aligned}
& K^{-}+p \rightarrow \Omega^{-}+\begin{array}{l}
+K^{+}+K^{0} \\
\begin{array}{l}
\Xi^{0}+\pi^{-} \\
\Lambda^{0}+\pi^{0}
\end{array}
\end{array}
\end{aligned}
$$

## particle: <br> Omega minus

symbol:
$\Omega^{-}$
charge:
mass:
spin:
category:
-1
1672.45 MeV/c²

3/2
Fermion, baryon, $\mathrm{I}=\mathrm{0}, \mathrm{B}=1, \mathrm{~S}=-3$

## the dominant Baryons

| Particle | Symbol | Rest Mass <br> $\mathbf{M e V} / \mathbf{c}^{2}$ | spin | $\mathbf{Q}$ | $\mathbf{B}$ | $\mathbf{s}$ | Lifetime | dominant decay <br> modes |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| proton | $p$ | 938.3 | $1 / 2$ | +1 | +1 | 0 | $>10^{31} \mathrm{y}$ |  |
| neutron | $n$ | 939.6 | $1 / 2$ | 0 | +1 | 0 | 920 | $p e^{-} \bar{\nu}_{e}$ |

## mesons

|  | Quark | Symbol | $\begin{gathered} \text { Rest } \\ \text { Mass } \\ \text { MeV/c² } \end{gathered}$ | spin | Q | B | S |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | up | $u$ | 1.7-3.3 | 1/2 | +2/3 | $1 / 3$ | 0 |  |  |
|  | down | $d$ | 4.1-5.8 | 1/2 | -1/3 | 1/3 | 0 |  |  |
|  | strange | $s$ | 101 | 1/2 | -1/3 | 1/3 | -1 |  |  |
| The pion |  | Particle | Symbol | antiparticle | Rest Mass $\mathrm{MeV} / \mathrm{c}^{2}$ | spin | Q | B | S |
|  |  | Pion | $\pi^{+}$ | $\pi^{-}$ | 139.6 | 0 | +1 | 0 | 0 |
|  |  | $\pi^{+}=\left(\begin{array}{lll}u & \& & \bar{d}\end{array}\right)$ |  |  |  | has the right stuff. |  |  |  |
| Q: |  | +1 | $+2 / 3+-(-1 / 3)$ |  |  |  |  |  |  |
| B: |  | 0 | $1 / 3+-(1 / 3)$ |  |  |  |  |  |  |
| S: |  | 0 |  | 0 |  |  |  |  |  |

## a little different

a similar thing happens for the mesons
$\operatorname{spin} 1$


## meson quark <br> content


spin 1


I

## the dominant Mesons

| Particle | Symbol | anti- <br> particle | Rest <br> MeV/c | spin | $\mathbf{Q}$ | $\mathbf{B}$ | $\mathbf{s}$ | Lifetime | dominant decay <br> modes | quark content |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Pion | $\pi^{+}$ | $\pi^{-}$ | 139.6 | 0 | +1 | 0 | 0 | $2.6 \times 10^{-8}$ | $\mu^{+} \nu_{\mu}$ | $u \bar{d}$ |
| Pi-zero | $\pi^{0}$ | $\pi^{0}$ | 135 | 0 | 0 | 0 | 0 | 920 | $2 \gamma$ | $\frac{1}{\sqrt{2}(u \bar{u}+d \bar{d})}$ |
| Kaon | $K^{+}$ | $K^{-}$ | 493.7 | 0 | +1 | 0 | +1 | $1.24 \times 10^{-8}$ | $\mu^{+} \nu_{\mu}, \pi^{+} \pi^{0}$ | $u \bar{s}$ |
| K-short | $K_{S}^{0}$ | $K_{S}^{0}$ | 497.7 | 0 | 0 | 0 | +1 | $0.89 \times 10^{-10}$ | $\pi^{+} \pi^{-}, 2 \pi^{0}$ | $d \bar{s}, s \bar{d}$ |
| K-long | $K_{L}^{0}$ | $K_{L}^{0}$ | 497.7 | 0 | 0 | 0 | +1 | $5.2 \times 10^{-8}$ | $\pi^{ \pm} \ell^{\mp} \nu_{\ell}$ | $d \bar{s}, s \bar{d}$ |
| Eta | $\eta^{0}$ | $\eta^{0}$ | 548.8 | 0 | 0 | 0 | 0 | $<10^{-18}$ | $2 \gamma, \pi^{+} \pi^{-} \pi^{0}$ | $u \bar{u}, d \bar{d}, s \bar{s}$ |
| Eta-prime | $\eta^{0 \prime}$ | $\eta^{0 \prime}$ | 958 | 1 | 0 | 0 | 0 | $\ldots$ | $\pi^{+} \pi^{-} \eta$ | $u \bar{u}, d \bar{d}, s \bar{s}$ |
| Rho | $\rho^{+}$ | $\rho^{-}$ | 770 | 1 | +1 | 0 | 0 | $0.4 \times 10^{-23}$ | $\pi^{+} \pi^{-}, 2 \pi^{0}$ | $u \bar{d}$ |
| Rho-naught | $\rho^{0}$ | $\rho^{0}$ | 770 | 1 | 0 | 0 | 0 | $0.4 \times 10^{-23}$ | $\pi^{+} \pi^{-}$ | $u \bar{u}, d \bar{d}$ |
| Omega | $\omega^{0}$ | $\omega^{0}$ | 782 | 1 | 0 | 0 | 0 | $0.8 \times 10^{-22}$ | $\pi^{+} \pi^{-} \pi^{0}$ | $u \bar{u}, d \bar{d}$ |
| Phi | $\phi$ | $\phi$ | 1020 | 1 | 0 | 0 | 0 | $20 \times 10^{-23}$ | $K^{+} K^{-}, K^{0} \bar{K}{ }^{0}$ | $s \bar{s}$ |

## spins work out

Keep track of quark spins

$$
\begin{array}{ll}
\operatorname{spin}+1 / 2 & q \uparrow \\
\operatorname{spin}-1 / 2 & q \downarrow
\end{array}
$$

for example, a couple of baryons:

$$
p \quad u \uparrow u \downarrow d \uparrow \quad \text { total spin: } 1 / 2
$$

$$
\Delta^{+} \quad u \uparrow u \uparrow d \uparrow \quad \text { total spin: } 3 / 2
$$

for example, a couple of mesons:

$$
\begin{array}{lll}
\pi^{+} & u \uparrow \bar{d} \downarrow & \text { total spin: } 0 \\
\rho^{+} & u \uparrow \bar{d} \uparrow & \text { total spin: } 1
\end{array}
$$

## there are still

A model of "quark molecules"...

## 100's more baryons and mesons

what's up with that? you're asking

Molecules can have vibrational and rotational excited states...

So can quarks.
$N^{*}$ is a state with the same quark content as a proton
$N^{*}$ but which has a high orbital angular momentum

$$
d \quad u \quad u
$$

Other states can be well-modeled by assuming relative vibrational modes..

$$
d
$$

## you can

 tell a particle physicist by the books that we carry"I laughed, I cried"


## now the

## jargon

Hadrons: particles made of quarks.

## gets a little more straightforward



Mesons: particles made of 1 quark and 1 antiquark.

## a variety of consequences

One could begin to understand particle decays and reactions in terms of pseudo-Feynman diagrams* like this:

$$
\begin{array}{ll}
\pi^{+}+p \rightarrow \pi^{+}+p \quad \begin{array}{l}
\text { Fermi had produced "resonances" } \\
\\
\\
\\
\\
\text { "hat suggested that something was } \\
\text { in between" the initial and final } \\
\text { states }
\end{array}
\end{array}
$$

$$
\pi^{+}+p \rightarrow \Delta^{++} \rightarrow \pi^{+}+p
$$


scatterings now are thought of diferently
by following the lines...
$\pi^{+}+p \rightarrow \Delta^{++} \rightarrow \pi^{+}+p$
Feynman Diagram, pre-1964:

in quark language:


## how about a strong interaction decay?

a little nonintuitive.
$\Delta^{0} \rightarrow \pi^{-}+p$
the old way:

the quark way:


3 quarks
some quark-creation required!


5 quarks
stay tuned.

## is the world made of actual

 quarks?or is this just a convenient organizing scheme
that's all Gell-Mann thought

But evidence started to accumulate that surprised everyone

# First piece of convincing evidence: 

## we can bang on them

individually...Feynman saw this first.

## remember.

the crucial thing in order to "see" something?
wavelength has to be about the size of the object
larger the momentum
the smaller the spatial resolving capability
scattering of an electron from a nucleus
slow electron, long wavelength photon

"sees" the whole nucleus
scattering of an electron from a nucleus
fast electron, medium-short wavelength photon
(e)

"sees" an individual proton in the nucleus
scattering of an electron from a nucleus
very fast electron, very-short wavelength photon
(e)

"sees" an individual quark in a proton or neutron
That's how we became convinced in 1969 -
the same sort of backwards scattering as Rutherford's

## Share this:

## The Nobel Prize in Physics 1990



Jerome I. Friedman Prize share: $1 / 3$


Henry W. Kendall Prize share: $1 / 3$


Photo: T. Nakashima Richard E. Taylor Prize share: 1/3

The Nobel Prize in Physics 1990 was awarded jointly to Jerome I. Friedman, Henry W. Kendall and Richard E. Taylor "for their pioneering investigations concerning deep inelastic scattering of electrons on protons and bound neutrons, which have been of essential importance for the development of the quark model in particle physics".

Photos: Copyright © The Nobel Foundation

## Share this:

[^0]particle: up quark
symbol:
charge:
mass:
spin:
category:
$+2 / 3$
1.7 to $3.3 \mathrm{MeV} / \mathrm{c}^{2}$

1/2
Fermion, $\mathrm{I}=+1 / 2, \mathrm{~B}=1 / 3, \mathrm{~S}=0$
particle: down quark
symbol:
charge:
mass:
spin:
category:
$-1 / 3$
4.1 to $5.8 \mathrm{MeV} / \mathrm{c}^{2}$

1/2
Fermion, $\mathrm{l}=-1 / 2, \mathrm{~B}=1 / 3, \mathrm{~S}=0$
particle: strange quark
symbol: s
charge: $\quad-1 / 3$
mass:
spin:
category:
$101 \mathrm{MeV} / \mathrm{c}^{2}$
1/2
Fermion, $\mathrm{I}=-1 / 2, \mathrm{~B}=1 / 3, \mathrm{~S}=-1$

# shifting gears 

the weak interaction needs a boson
the quantum relativistic field theory theme song:


## this kind of magic:

If there is a force...there's a field


If there's a field, there's a quantum to go with it.

Because Nature is Clumpy.

## for the electromagnetic interaction:

 the force is the electromagnetic force the field is $E \& B$ the clumpiness - the quantum - is:The photon: $\gamma$


## Well, the Weak Force

 must have a field...yadda yadda yadda

If there is a force...there's a field


If there's a field, there's a quantum to go with it.

Because Nature is Clumpy.
for weak interaction:
the field must be a weak field...\& Massive \& electrically charged
the clumpiness -the quantum - must be something else.
here's a weak interaction

## neutron beta

decay

changes electric charge
the weak interaction here changes the bottom and the top of these doublets

Manipulate the graph in the now familiar way:


## the muon

## decay is

## the same

 sort ofin that second way of looking at it:


$$
\binom{\nu_{e}}{e}>\text { and } \quad\binom{\nu_{\mu}}{\mu}
$$

## do it

## again?

## can a "photon" be forced to exist that governs the weak interaction?

It was a dream that the electromagnetic interaction

could have a weak interaction counterpart.


Feynman and Murray Gell-Mann worked out a consistent theory based on the idea of a "heavy" photon with electric charge.
"W" for "Weak"

Notice that $f$ and $f^{\prime}$ and $\mathrm{W}^{ \pm}$all have to have their electric charges assigned so that electric charge is conserved.

## temporary

 entries
## into your

 table of primitive diagrams
so, a new primitive diagram

## for the Weak Interaction



## keep <br> track of the <br> charge <br> flow

there are 2 W
charged states


$$
\begin{aligned}
& n \rightarrow p+W^{-} \rightarrow p+e^{-}+\bar{\nu}_{e} \\
\mathrm{Q}: & 0=+1+-1=+1+-1+0=0
\end{aligned}
$$

So: $\quad W^{-}$lowers the electrical charge by 1

$$
W^{+} \text {raises the electrical charge by } 1
$$



## here is

## where

those weak "doublets"
come in

The particle doublets that we know so far:

making these transitions is the W Boson's job.


Notice, that all of these transitions change the electric charge as well as the particle type

call a generic lepton, " $\ell$ "

$$
\begin{aligned}
& \binom{\nu_{\ell}}{\ell} \longleftrightarrow W \\
& \ell=e, \mu, \tau
\end{aligned}
$$

## "deep inelastic scattering"

hitting quarks individually
of course in a statistical fashion
neutrinos do it too...

analyses of these reactions,

$$
\nu N \rightarrow \mu X \quad e N \rightarrow e X
$$


confirm the point-like (?) nature of quarks
confirm their apparent loose-binding within nucleons (in a second)
confirm their fractional electric charges!
so, a new
primitive
diagram

## for the Weak

 Interaction with quarks, to go with the leptons

and in the quark interpretation: the reason $W$ does: $\binom{p}{n} \zeta$ is because it does: $\binom{u}{d} \zeta$

instead of what I had before:

## there are

 still weak
## interactions

## including transitions among quarks

The particle doublets that we know so far:

making these transitions is the W Boson's job.


Notice, that all of these transitions change the electric charge as well as the particle type

call a generic lepton, " $\ell$ "

$$
\begin{aligned}
& \binom{\nu_{\ell}}{\ell} \longleftrightarrow W \\
& \ell=e, \mu, \tau
\end{aligned}
$$

## there are

## still weak

interactions

## including <br> transitions among

 quarksThe particle doublets in quark language:
Q

$$
\begin{aligned}
& +2 / 3 \\
& -1 / 3
\end{aligned}\binom{u}{d} \longleftrightarrow W\binom{?}{s}\left\langle W \begin{array}{l}
\text { making these transitions } \\
\text { is still the W Boson's job. }
\end{array}\right.
$$

call a generic lepton, " $\ell$ " call a generic quark, " $q$ " $\binom{\nu_{\ell}}{\ell} \longleftrightarrow W\binom{q}{q^{\prime}} \longleftrightarrow W$
$\ell=e, \mu, \tau \quad q=u, d, s$


> or:
call a generic fermion, " $f$ "

$$
\begin{aligned}
& \binom{f}{f^{\prime}} \longleftrightarrow W \\
& f=\ell, q
\end{aligned}
$$

NOW . . . your
second

## entry into

your
table of primitive diagrams

particle: charm quark
symbol:
charge:
mass:
spin:
category:

## C

$+2 / 3$
$1,270 \mathrm{MeV} / \mathrm{c}^{2}$
1/2
Fermion, $\mathrm{l}=0, \mathrm{~B}=1 / 3, \mathrm{~S}=0, \mathrm{C}=+1$

## SO,

## decays

## we've

## seen

just put in the decaying quark and let the other "spectator quarks"
come along for the ride

$$
\pi^{+} \rightarrow \mu^{+}+\nu_{\mu}
$$

responsible for making neutrino beams from proton accelerators


## Strong interaction, again:

The original question about nuclei...
now in play for quarks:
what holds the quarks inside of the baryons and mesons?

## Gross,

## Politzer,

## and

## Wilczek

## 2004

"asymptotic freedom" in strong interactions

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David J. Gross
H. David Politzer

Frank Wilczek

The Nobel Prize in Physics 2004 was awarded jointly to David J. Gross, H. David Politzer and Frank Wilczek "for the discovery of asymptotic freedom in the theory of the strong interaction".

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MLA style: "The Nobel Prize in Physics 2004". Nobelprize.org. 10 Apr 2013
http:/henw.nobelprize.org/nobel_prizes/physics/laureates/2004/
it's the glue that holds everything together virtually

## Predicted the existence of the Strong Messenger

 Particle: the Gluonmy gluon


## third <br> entry <br> into your

## table of primitive diagrams



# there are two amazing things 

about gluons

## thing 1

they self-interact

a photon propagates the electromagnetic force...but it does not have an electric charge

the gluon propagates the strong force...and it DOES have a "strong charge"

This has significant consequences...almost magical

## fourth and fifth entries

 into your table of primitive diagrams

## thing 2

their force field is the opposite of electromagnetism, or gravity


## ah, but the gluon is odd

## pull ${ }^{6} \mathrm{em}$ <br> apart

## called

quark confinement


We don't

## see

individual quarks or
gluons
they make more quarks and gluons
and interact very quickly into a cascade of particles
"quark-gluon jets"


## in ATLAS



## '"hard"' quark production


particle: gluon
symbol: $\quad g$
charge: 0
mass: 0
spin:
category:
1
Strong Vector Boson
three
forces now of vastly different strengths

Electromagnetic force 0.007


Weak force 0.000001



## proton

symbol:
charge:
mass:
spin:
category:


## down quark

symbol:
charge:
mass:
spin:


Fermion, $\mathrm{I}=-1 / 2, \mathrm{~B}=1 / 3, \mathrm{~S}=0$
particle:
up quark
symbol:
charge:
mass:
spin:
category:


Fermion, $\mathrm{I}=+1 / 2, \mathrm{~B}=1 / 3, \mathrm{~S}=0$
why does the proton weigh?




Field Energy

## SO:

## $m=$

when you step on the scale

## you measure the earth's attraction

to the gluons' mass-energy in your protons and neutrons and you use the non-quantum Newton's theory to do it
your "weight" is a quantum relativistic field theoretic thing

# here's the elementary particles story 

circa 1975

## the

## messengers

spin 1 Bosons<br>circa 1980

particle: bottom quark
symbol:
charge:
mass:
spin:
category:
b
$-1 / 3$ e
4.5 GeV/c² $=4.5 \mathrm{p}$

1/2
Fermion, quark

## the

## '"top quark'"

## was

## discovered in 1995

by two experiments at Fermilab

## with MSU faculty and students intimately involved

## Observation of the Top Quark

The DO Collaboration reports on a search for the sandard model top quark in pp collisions at $\sqrt{s}=1.8 \mathrm{TeV}$ ut the Fermilab Tevarron with an integruted luminosity of approximately $50 \mathrm{p}^{-1}$. We have searched for $\pi$ production in the dilkptoun and single-lepton decay channels with and without tagging of $b$-quark jets. We observed 17 events with an expected hackground of $3.8 \pm 0.6$ events. The probability for an upward fluctuation of the background to produce the observed signal is $2 \times 10$
(equivalent to 4.6 standard deviations). The kinematic properties of the excess events are consistent (equivalent to 4.6 standard deviations). The kinematic properties of the excess events are consisten
with top quark decay. We conclude that we have observed the top quark and measured its mass to be $199-21$ (stat) $\pm 22$ (syys) $\mathrm{CeV} / \mathrm{c}^{1}$ and its production cross section to be $6.4 \pm 2.2 \mathrm{pb}$
Pacs nembers $14.65 . \mathrm{Hz} .138500 .1385 . \mathrm{Nu}$

```
We establish the existence of the top quark using a }67\mp@subsup{\textrm{pb}}{}{-1}\mathrm{ data sample of $pP collisions at
\sqrt{}{s}-1.8 TeV collected with the Collider Detector at Fermilab (CDF). Employing techniques similar
to those we previously published. we observe a signal consistent with Hf docay to WWb\overline{b}, but
inconsistent with the hackgroand prediction by 4.8%. Adlitional evidence for the top quark is
176 = %(stat) }=10(\mathrm{ sys) GeV/c
PACS numbers 14.65.H2. 1385.08. 13.85.Ni
```

February 24th, 11AM, we submitted our discovery paper to Physical Review Letters

March 2, 1995 the announcement was made at Fermilab


## top quark

symbol:
charge:
mass:
spin:
category:
$+2 / 3$ e
$172.0 \pm 2.2 \mathrm{GeV} / \mathrm{c}^{2}=172 \mathrm{p}$
1/2
Fermion, quark

## quarks \& leptons

## the weak

## interactions

## still operate with

 the increased doublet setsThe complete (circa 2000) particle doublets:

$$
\begin{aligned}
+2 / 3 & \binom{u}{-1 / 3}
\end{aligned}\binom{c}{s} \quad\binom{t}{b}
$$

## the weak

## interactions

still operate with the increased doublet sets

The complete (circa 2000) particle doublets:
Q

$$
+2 / 3 \quad\left(\begin{array}{l}
u \\
-1 / 3 \\
d
\end{array}\right) \longleftrightarrow W\binom{c}{s} \longleftrightarrow W\binom{t}{b} \longleftarrow W
$$

$$
\begin{gathered}
0 \\
-1
\end{gathered}\binom{\nu_{e}}{e} \longleftarrow W\binom{\nu_{\mu}}{\mu} \longleftrightarrow W\binom{\nu_{\tau}}{\tau} \longleftrightarrow W
$$

## the

## modern

## picture

of the elementary particle patterns circa 2000 and still current
the lepton families...lepton "doublets"

$$
\binom{\nu_{e}}{e^{-}} \quad\binom{\nu_{\mu}}{\mu^{-}} \quad\binom{\nu_{\tau}}{\tau^{-}}
$$

and their interactions: $\mathbf{X}$ no, $\boldsymbol{\checkmark}$ yes.

| leptons | $\nu_{e}$ | $e$ | $\nu_{\mu}$ | $\mu$ | $\nu_{\tau}$ | $\tau$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & \text { strong } \\ & \substack{\text { Onog }} \end{aligned}$ | $x$ | $x$ | $x$ | $x$ | $x$ | $x$ |
| $\sim_{\gamma}^{\text {ectumgene }}$ | $x$ | $\checkmark$ | $x$ | $\checkmark$ | $x$ | $\checkmark$ |
| $\sim_{W}^{\text {meak }}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| Lexional | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |

## the

## modern

## picture

of the elementary particle patterns circa 2000
the quark families...quark "doublets"

$$
\binom{u}{d} \quad\binom{c}{s} \quad\binom{t}{b}
$$

and their interactions: $\mathbf{X}$ no, $\boldsymbol{\checkmark}$ yes.

| quarks | $u$ | d | c | $s$ | $t$ | $b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Strong ${ }^{\text {m }}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| electromagnetic <br> nWn 2 | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| $\begin{gathered} \text { veak } \\ m W W \end{gathered}$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |
| Itional | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ | $\checkmark$ |

## The Particle Zoo?



## The Particle Zoo? tamed.



## shifting gears

the weak and electromagnetic forces are one.


## '"phase transitions"

## not a subject of Particle Physics

we thought
but we stole a theory from materials scientists
think about a phase transition

what a physicist sees is a change of symmetry


## there are basically <br> 2 kinds

## 1st Order nucleation

## 2d Order continuous



Boiling starts in various locations inside of liquid water

Other kinds of phase transitions happen uniformly throughout the substance.
you
probably

## are mostly

familiar
with:
freezing
melting
boiling

These "2nd Order," phase transitions are continuouseverywhere:
crystallization
changes of density
magnetism
superconductivity
superfluidity
plasma transition
electron gases
Bose gases


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