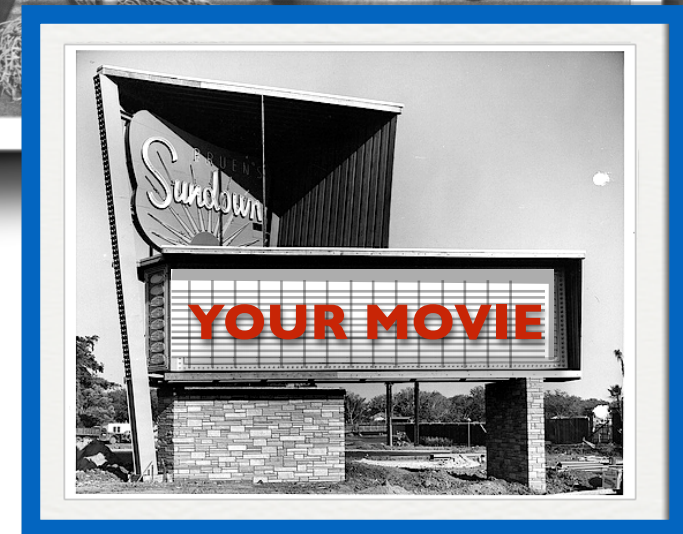


hi

Day 25, 05.12.2018

Quantum Mechanics 4

housekeeping



Gotta come to class

question about anything? I'll make a movie for you:

Quantum Mechanics:

Readings: Oerter, Cosmic Perspective, and Hobson

 ***Hobson_QM1.pdf & Hobson_QM2.pdf are chapters 12 & 13 out of Hobson***

Homework #11 is all from MasteringPhysics - ~~due saturday, rather than friday~~

yeah. but I messed up and slipped visibility and due dates by a week...
fixed last night, so due Monday, the 16th

honors project began

https://qstbb.pa.msu.edu/storage/Homework_Projects/honors_project_2018/

contains:

- the first instructions: the plan & tutorial

- the second instructions – v2 uploaded, added a missing student

- the data, assigned by name in the second instructions

dates:

- complete first part, March 16

- analyze data by April 24 and hand in complete writeup at the final exam

Here's

What You

Missed About

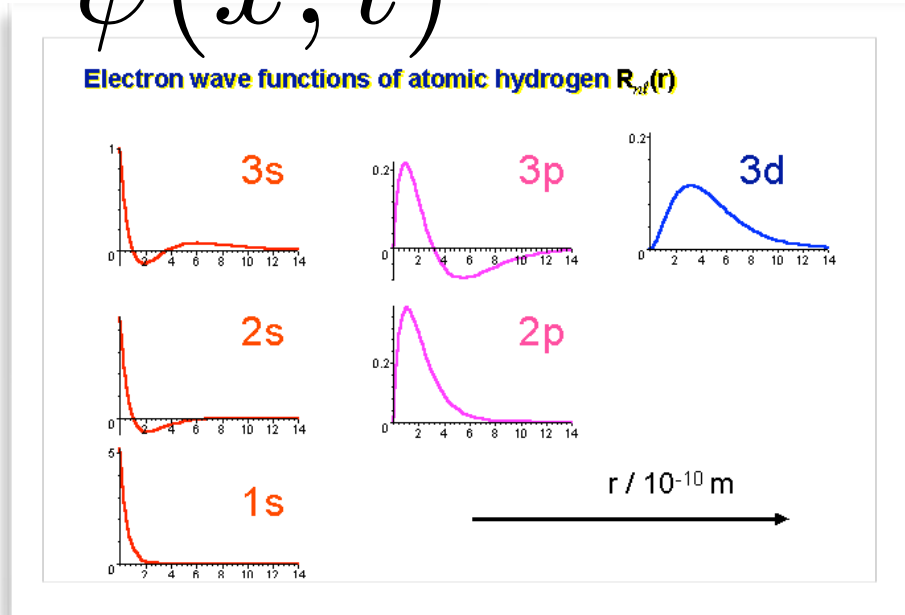
electrons



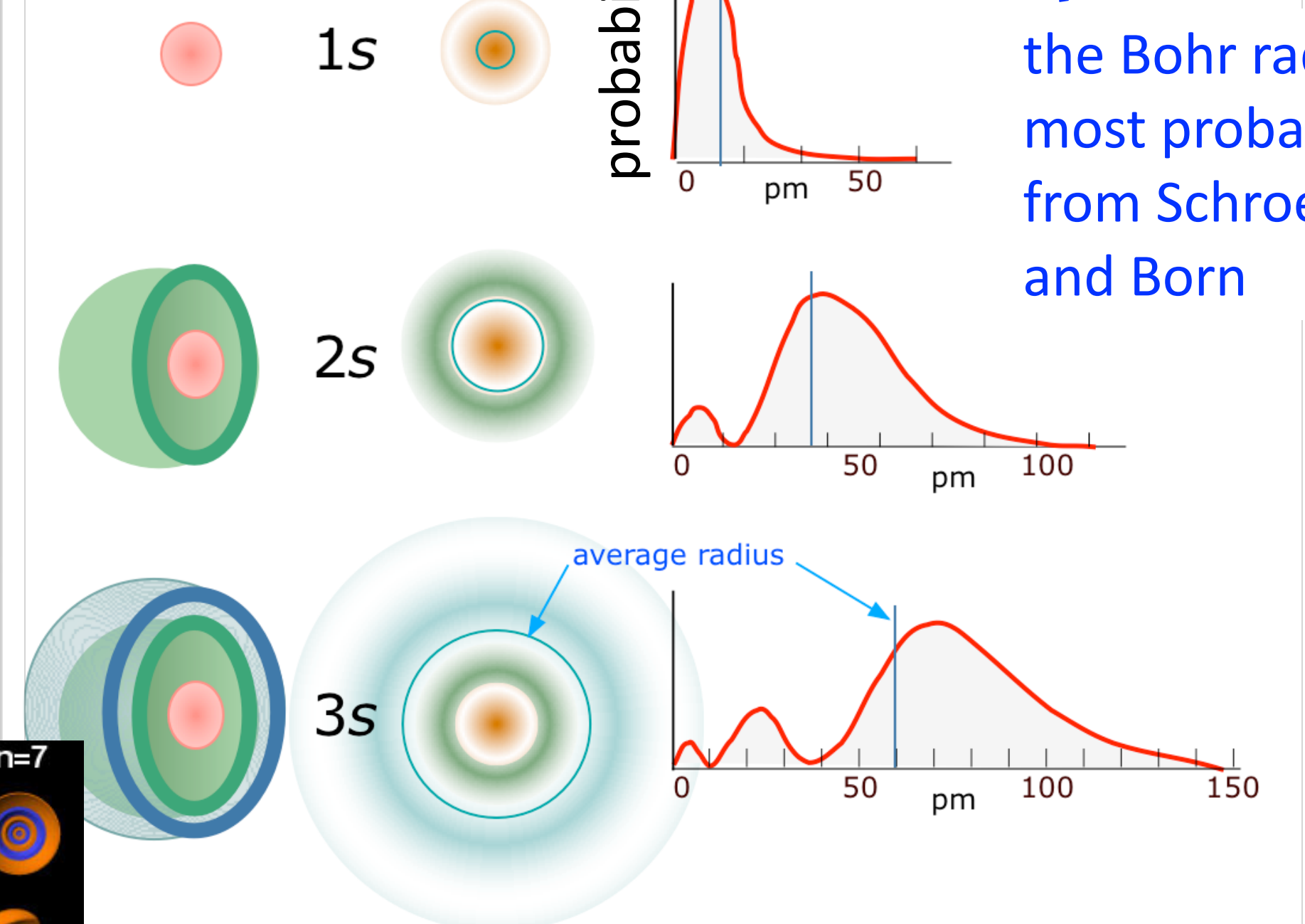
slice through the wavefunctions of Hydrogen

Solve Schroedinger equation and get wavefunctions:

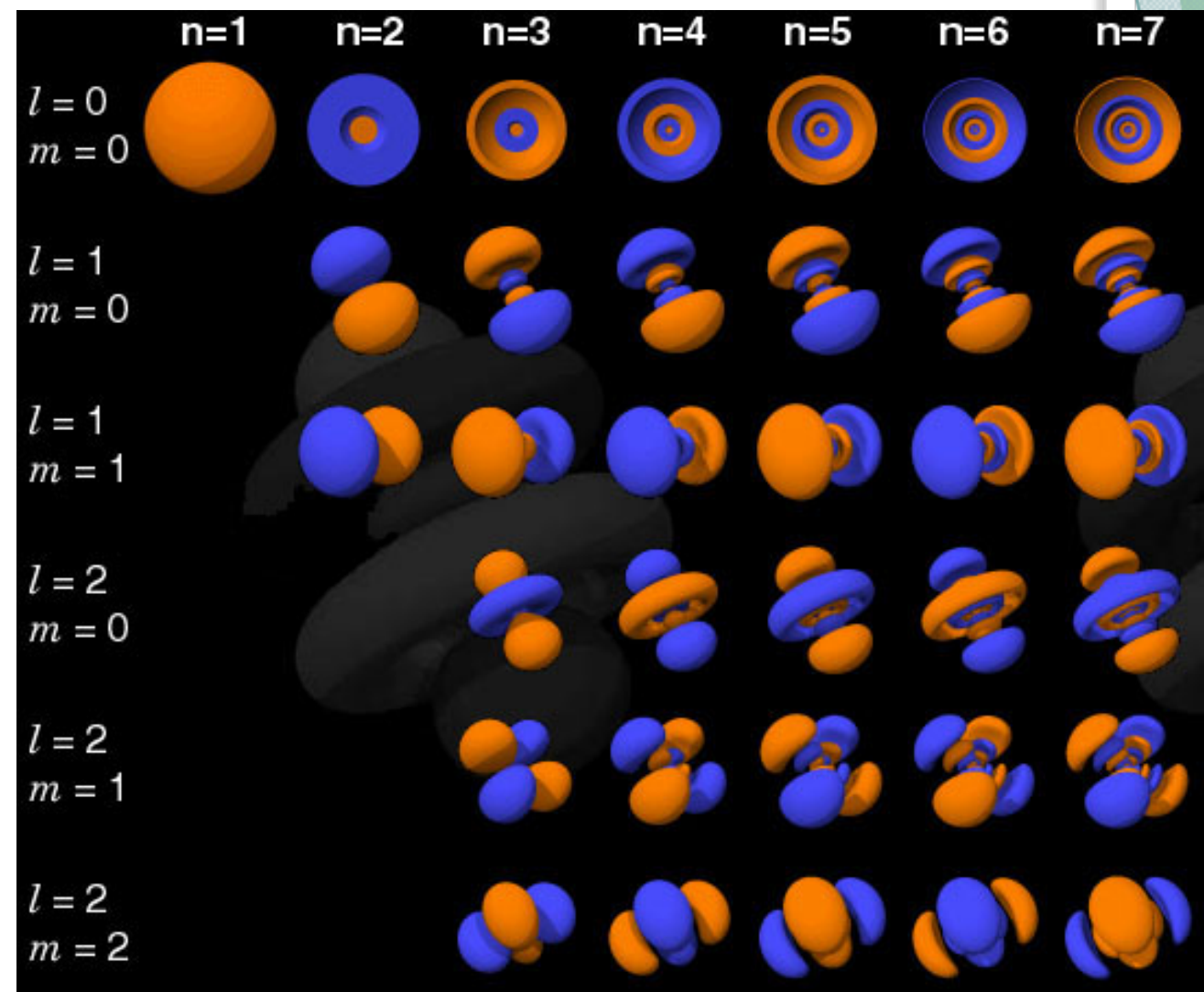
$$\psi(x, t)$$



then square them:



the Bohr radius - the most probable radius from Schroedinger and Born



relation alert:

Heisenberg Uncertainty Relation

refers to: $\Delta x \Delta p \geq h$ & $\Delta t \Delta E \geq h$

an inherent property of Nature

example:

lots of things!

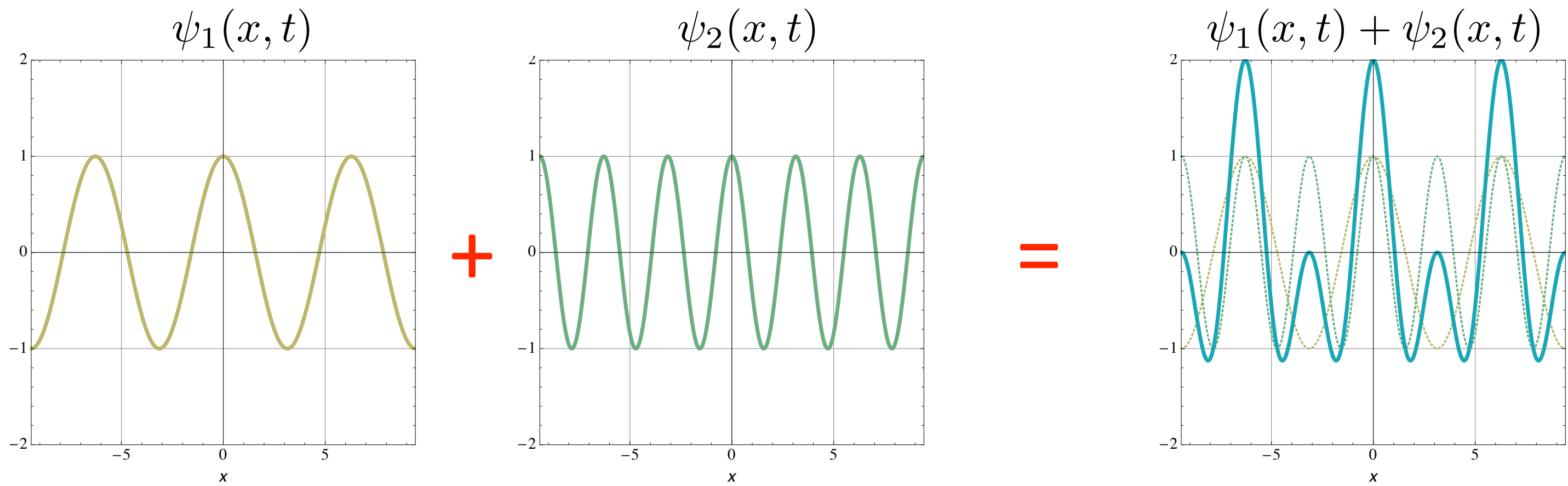
a new way

A measurement cannot be made of both precise position and precise momentum:
Objects in Nature don't possess those properties.

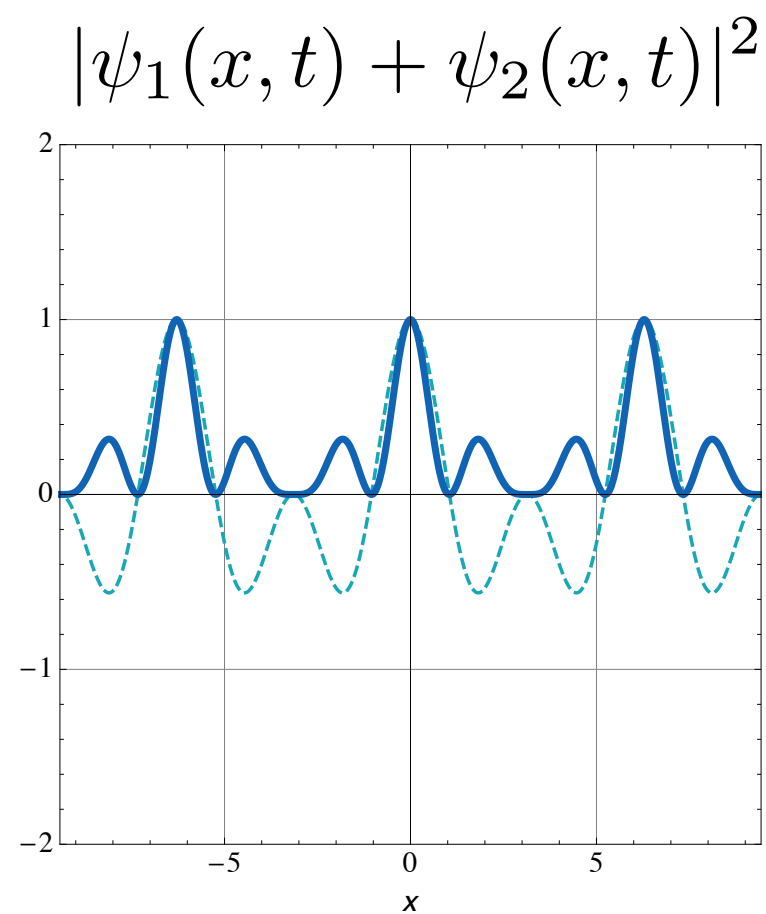
Of thinking and doing science

we lose another classical, unchallenged scenario





But, remember that what's real about the quantum fields is the square: $|\psi(x, t)|^2$



notice the peaking

(I've changed the heights)

a classical particle (dot) and its wavefunction

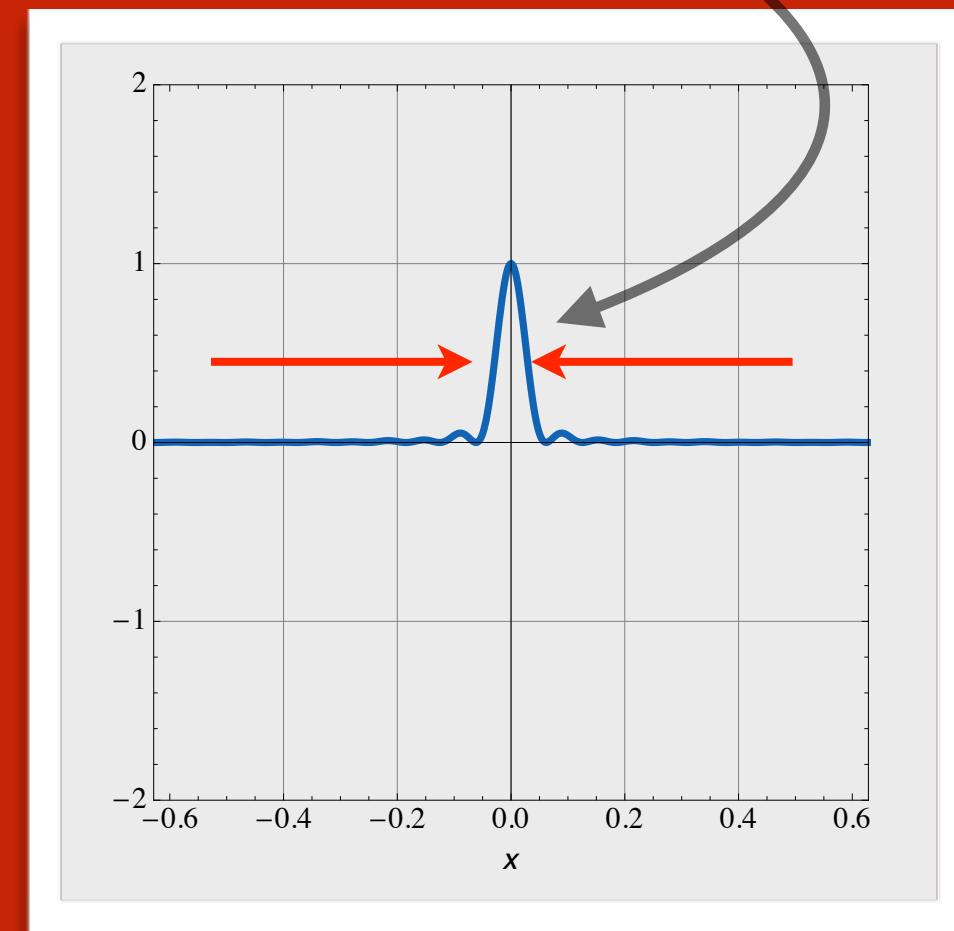
waves of different
wavelengths?
different momenta

Heisenberg Uncertainty Relation at
work again

called "wavepackets"

$$p = \frac{h}{\lambda}$$

the wave combinations localize
the state...with some spread in x



all of the wave combinations means all of the
momenta contribute: an spread in p .

the wavefunctions are everywhere

They're waves, after all.

make a measurement...there

Only then is it real.

the electron is there with probability $|\psi|^2$

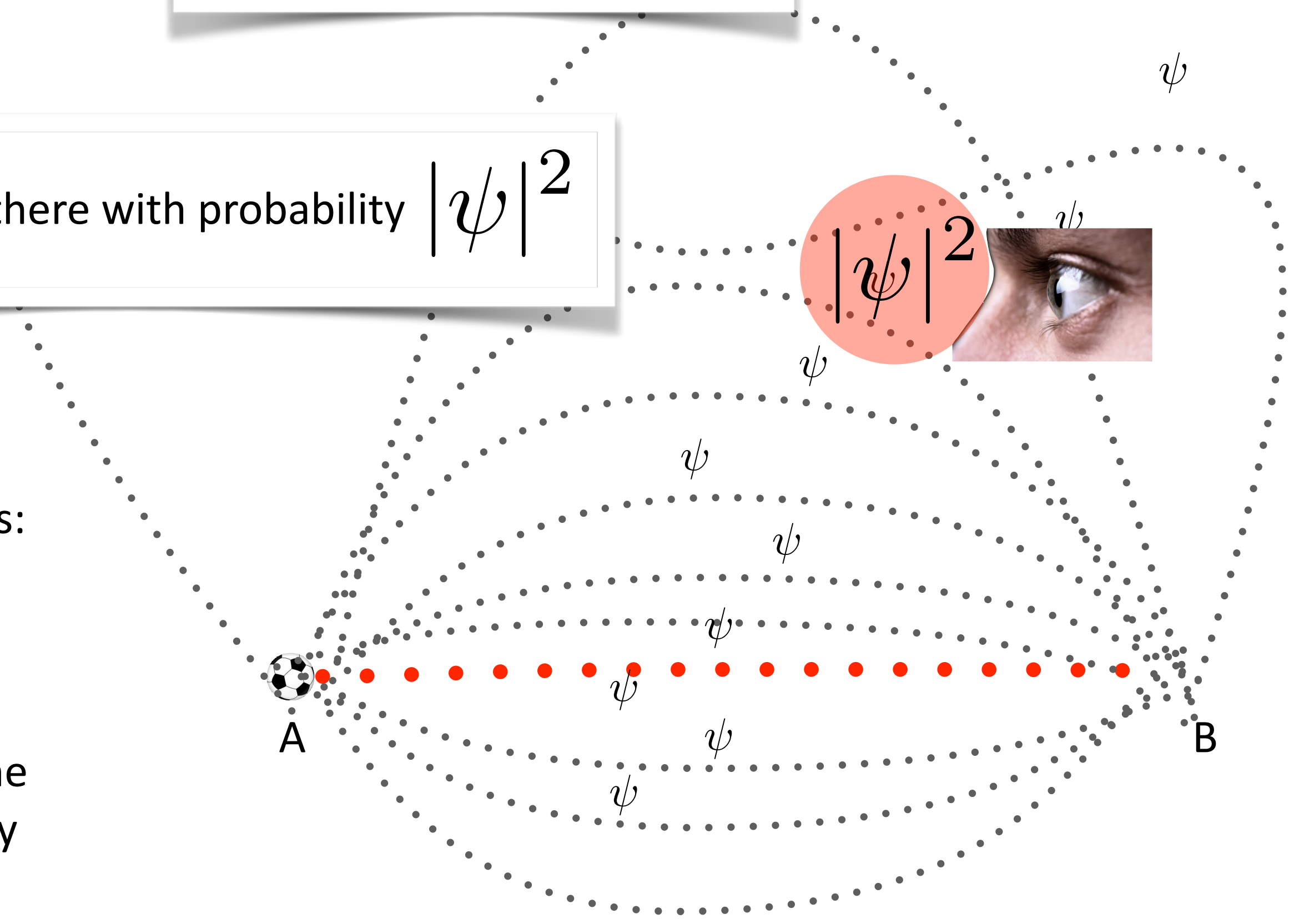


Feynman's picture was one of particles: which take all possible paths

We can calculate the wavefunction at any point, very precisely...it's completely deterministic

The trajectory of a big object?

Overwhelmingly probable quantum likelihood: the classical path



so where is a quantum

before it's measured?

anywhere? everywhere?

yeah.

this is how we have to think about it:

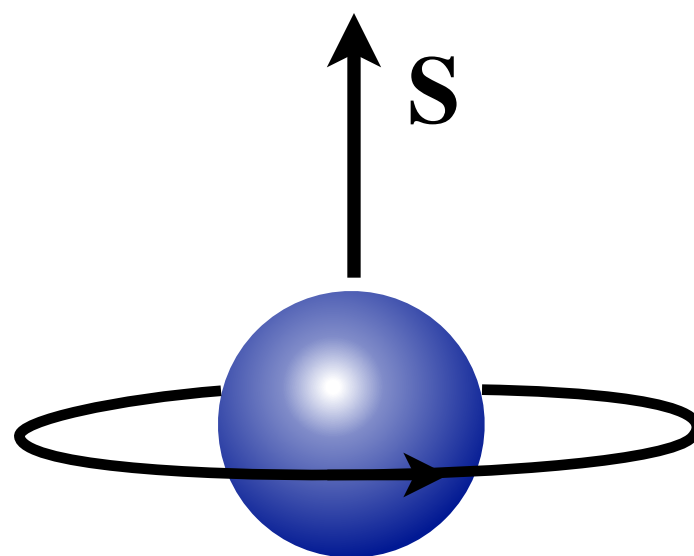
before measurement: alive-dead state -
superposition state of both

after measurement: is **either** alive **or** dead

OBTW:
electrons
are little
magnets

They behave in a
magnetic field as
if they are little
spinning current
spheres

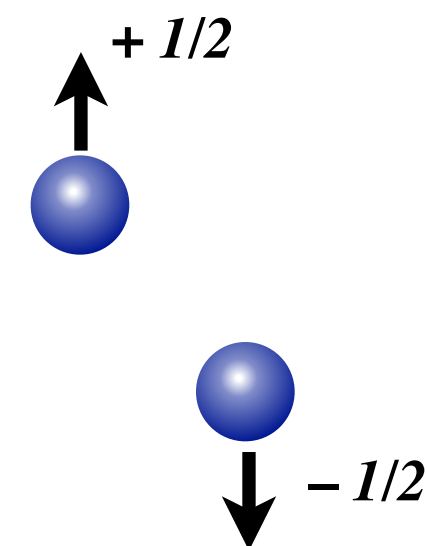
The electron **itself** is *like* a spinning charge...



Electrons have an **intrinsic** angular momentum, "S": "spin"

$$S_z = m_s \frac{h}{2\pi}$$

We say
"spin, plus 1/2" or "spin up"
and
"spin, minus 1/2" or "spin down"



jargon alert:

fermion

refers to:

any particle with half-integer spin

entomology:

from Fermi's theoretical work on the behavior of large numbers of Fermions

example:

electron, proton, neutron

jargon alert:

boson

refers to:

any quantum object with integer spin

etymology:

from Satyendra Nath Bose, who worked on the effects of multiple boson aggregates

example:

photon, pion, Higgs Boson

spin is a defining quality of an electron

electron

symbol:

e

charge:

$-1e$

mass:

$m_e = 9.0 \times 10^{-31} \text{ kg} \sim 0.0005 \text{ p}$

spin:

$1/2$

category:

fermion, lepton

particle:

proton

symbol:

p

charge:

$+1e$

mass:

$m_p = 1.6726 \times 10^{-27} \text{ kg} = 1 \text{ p}$

spin:

$1/2$

category:

fermion, hadron

again, an inherent angular momentum and a defining property of photons

particle:

photon

symbol:

γ

charge:

0

mass:

$m_{\gamma} = 0$

spin:

1

category:

boson, aka Intermediate Vector Boson

“

I think I can safely say that nobody understands quantum mechanics.

Richard Feynman

But we can calculate with Quantum Mechanics very, very well.

We're all highly skilled Quantum *Mechanics*



shifting gears

antimatter



here's a number:

0

0

zero

the # of successfully combined models of
Quantum Mechanics and Relativity
prior to 1928

remember the
relativistic
energy
relationship

and compare it to
the non-
relativistic one

Classical

$$E = \frac{1}{2}mv^2$$

$$p = mv$$

$$v = \frac{p}{m}$$

Relativistic

$$E^2 = (m_0c^2)^2 + (pc)^2$$

that square is problematic since it suggests:

$$E = \pm \sqrt{(m_0c^2)^2 + (pc)^2}$$

translated to Schroedinger QM:

**negative energies for freely
moving electrons**



negative energies for unbound systems

a disaster

negative energies for unbound systems

a disaster

negative energies for unbound systems

a disaster

there's no bottom!



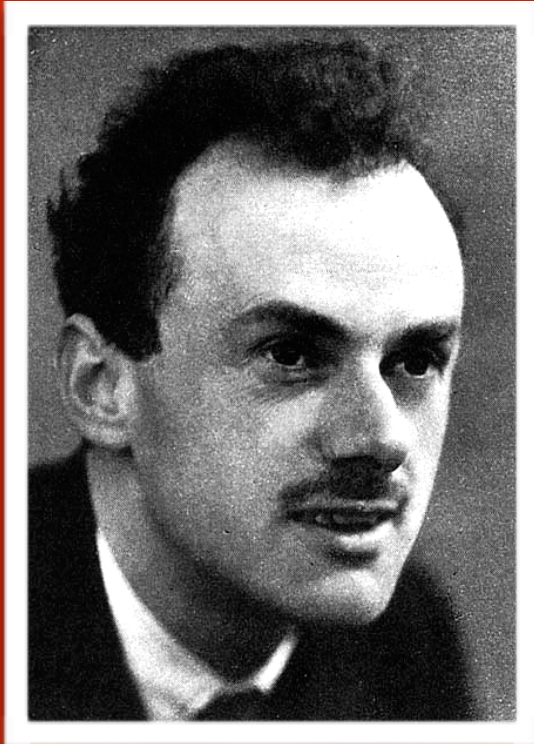
worse!

Quantum Mechanics using Relativity:

required not only negative energies

negative probabilities!

1928



Paul Dirac

1902 – 1984



“

At the question period after a Dirac lecture at the University of Toronto, somebody in the audience remarked: "Professor Dirac, I do not understand how you derived the formula on the top left side of the blackboard."

"This is not a question," snapped Dirac, "it is a statement."

hilarious interview with the Wisconsin State Journal from 1929 on the blog.

Dirac's Mathematical Imagination

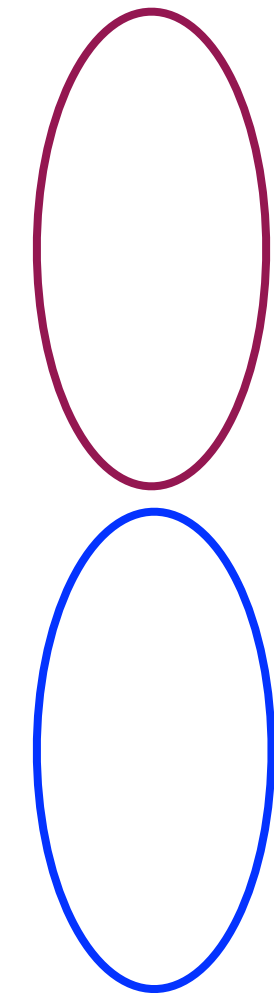
Dirac embraced the negative energy

Solved the negative probability

Dirac set out to find an equation that would solve both problems

Dirac's imagination

The "Dirac Equation" is the correct equation for electrons: Probabilities turn out okay, but required interpretation of negative energies



negative
electric charge
+ Energy

positive
electric charge
- Energy

Dirac's result

required: 4 quantum fields, rather than 1 $\psi_{up}(E), \psi_{down}(+E)$

2 have positive energy, 2 have negative energy $\psi_{up}(E), \psi_{down}(-E)$

each pair is related precisely to spin

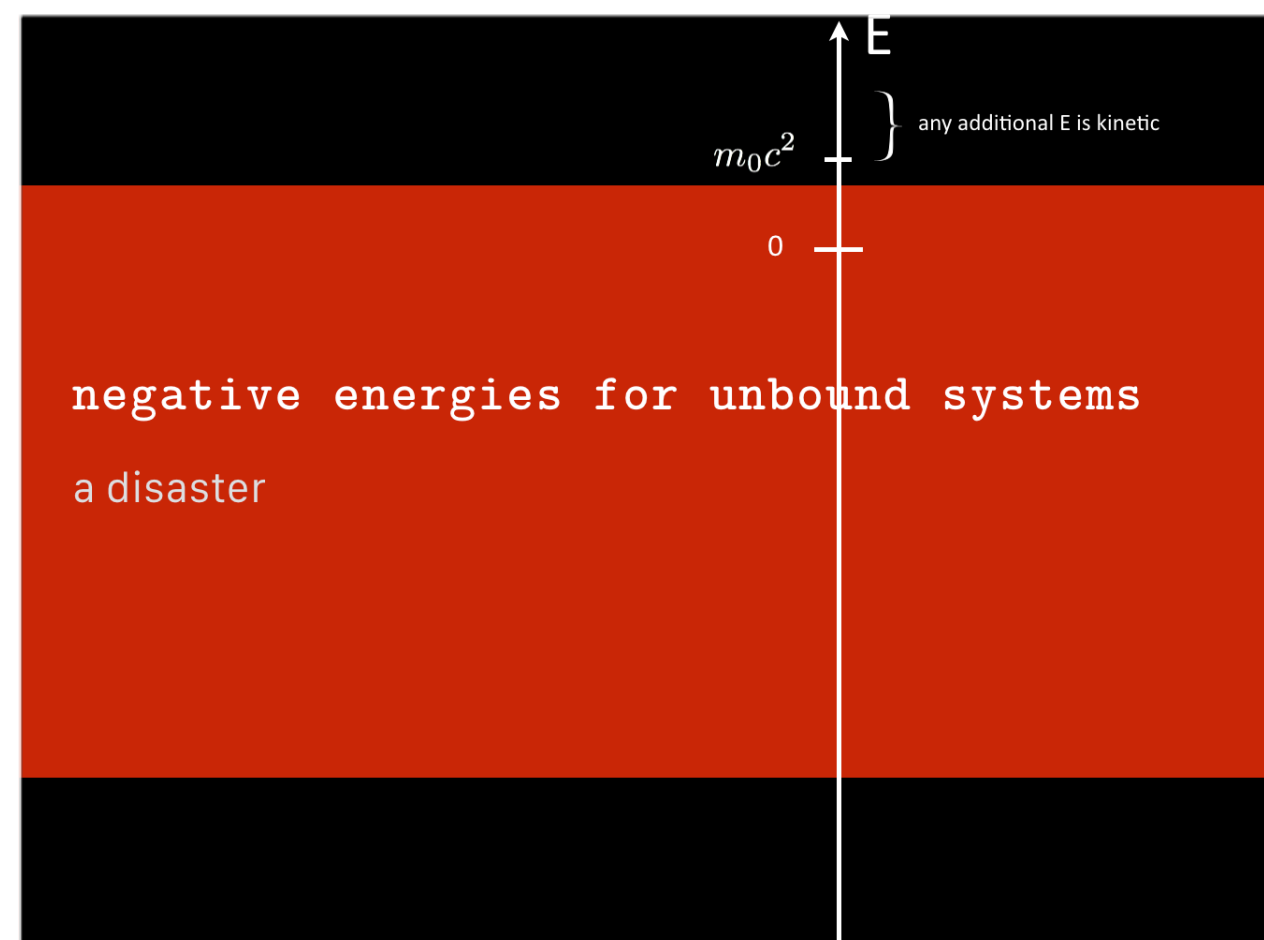
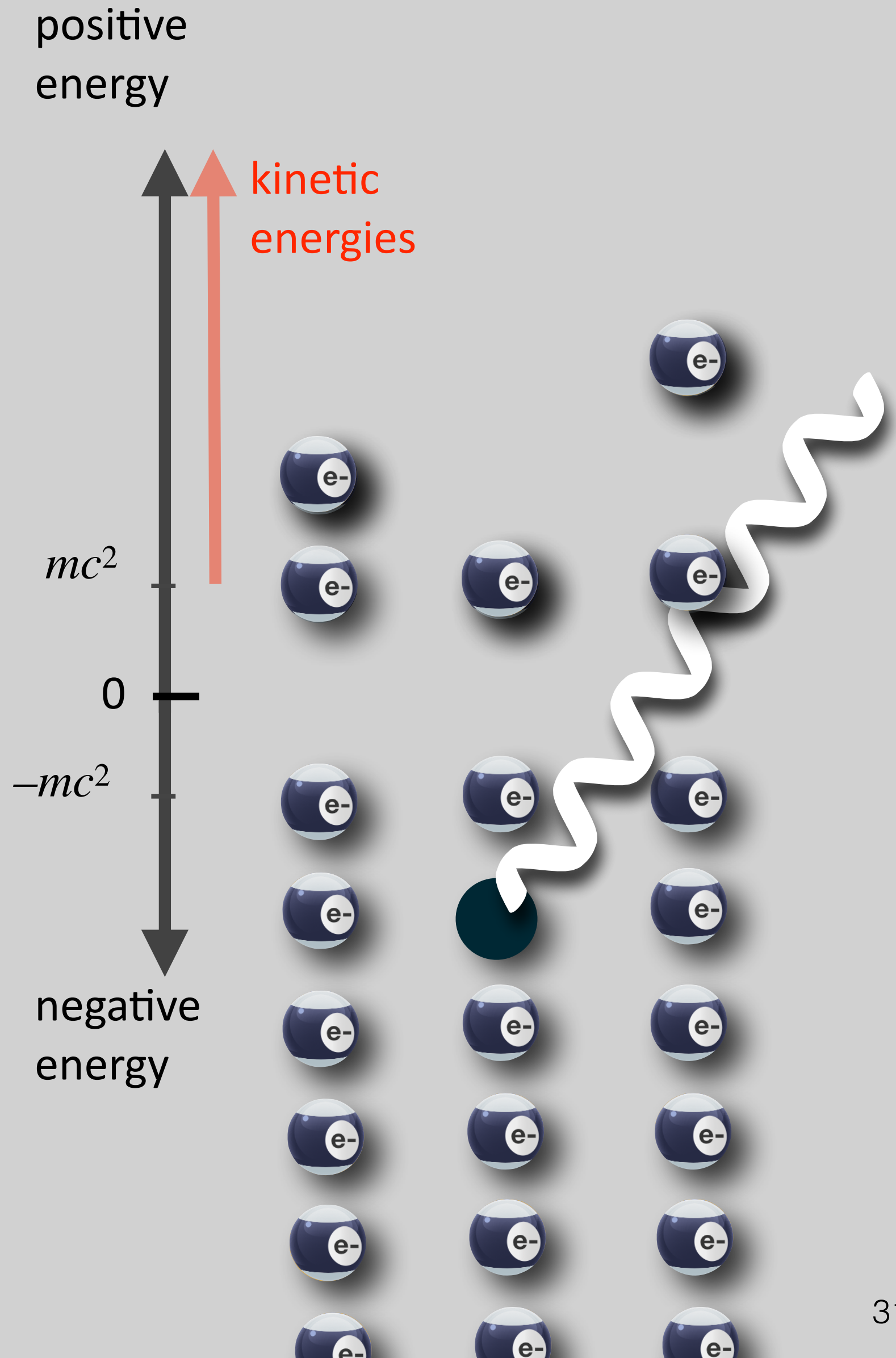
**Dirac showed that spin is a wholly relativistic effect
...it just popped out of his equation.**

still negative energies?

"solved" it with
Pauli's Exclusion
Principle

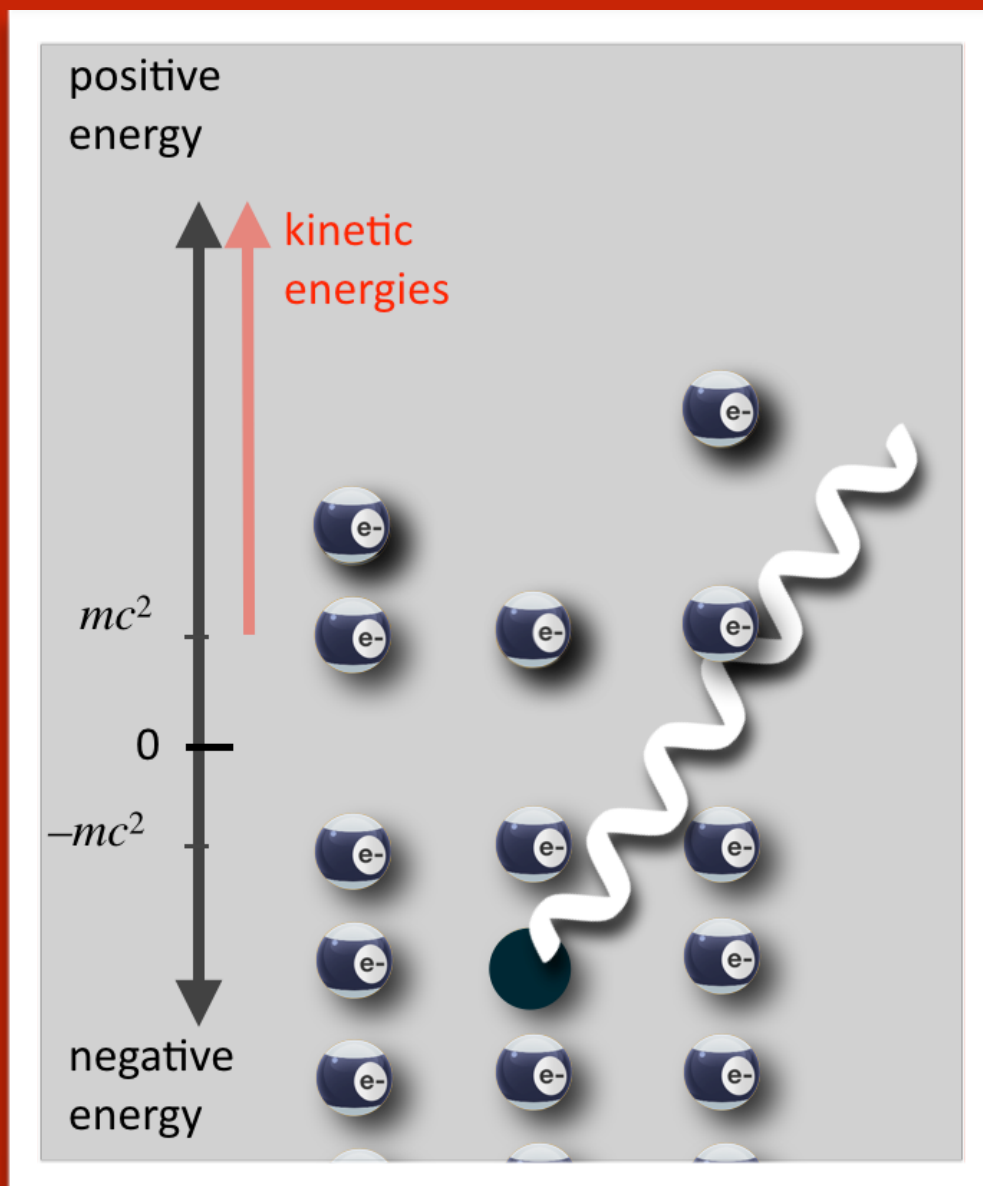


His vacuum is
full of negative
energy
electrons



start
with
nothing

$$E_\gamma > 2 m_e c^2$$



NOTHING
+ **Energy**



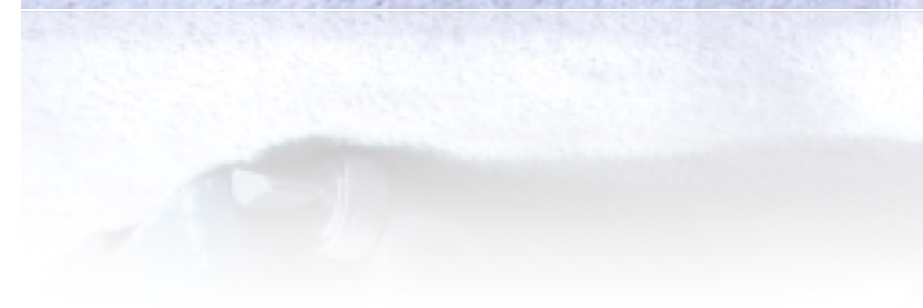
Let's talk about
Nothing.

Dirac began this
discussion

which continues today

in particle physics

and in cosmology



what is this?

$\psi(-E)$ a positively charged object with negative energy?

At first, he thought: "proton"

nah. A bolder idea: an anti-electron. The Positron.

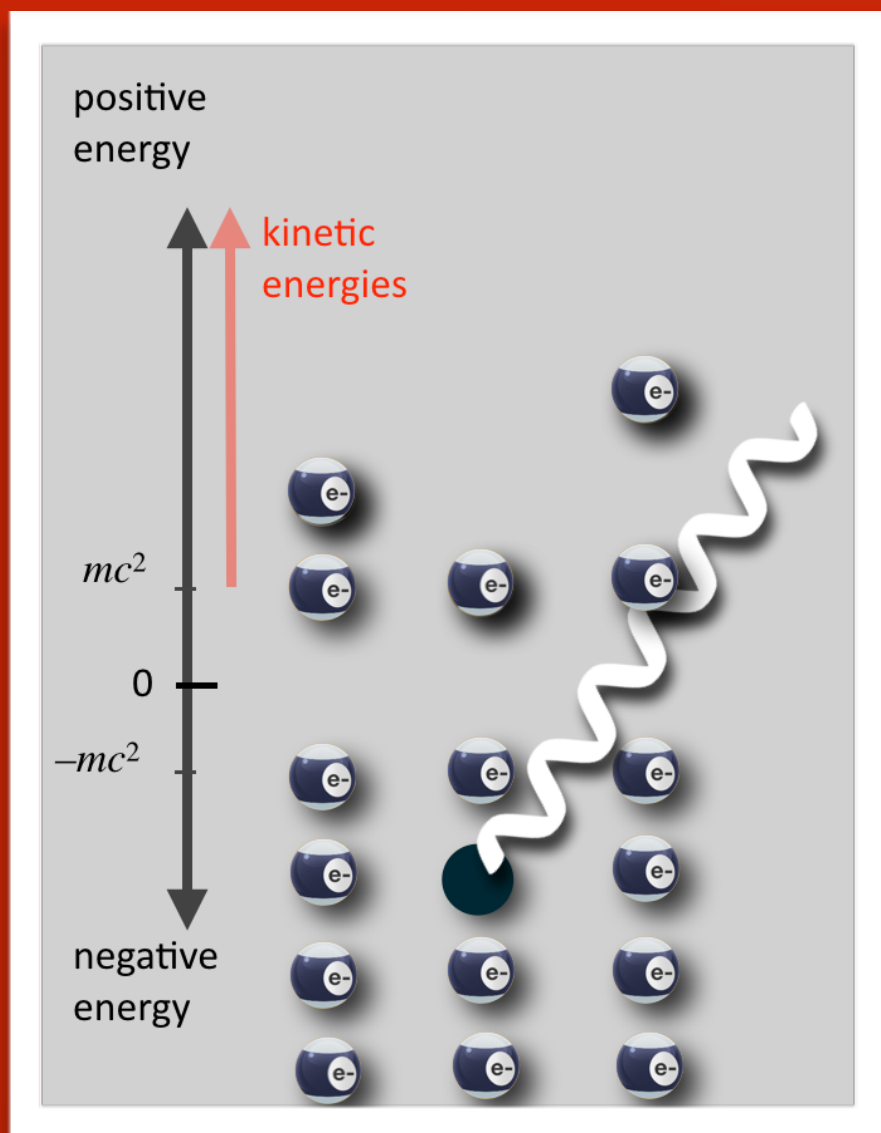
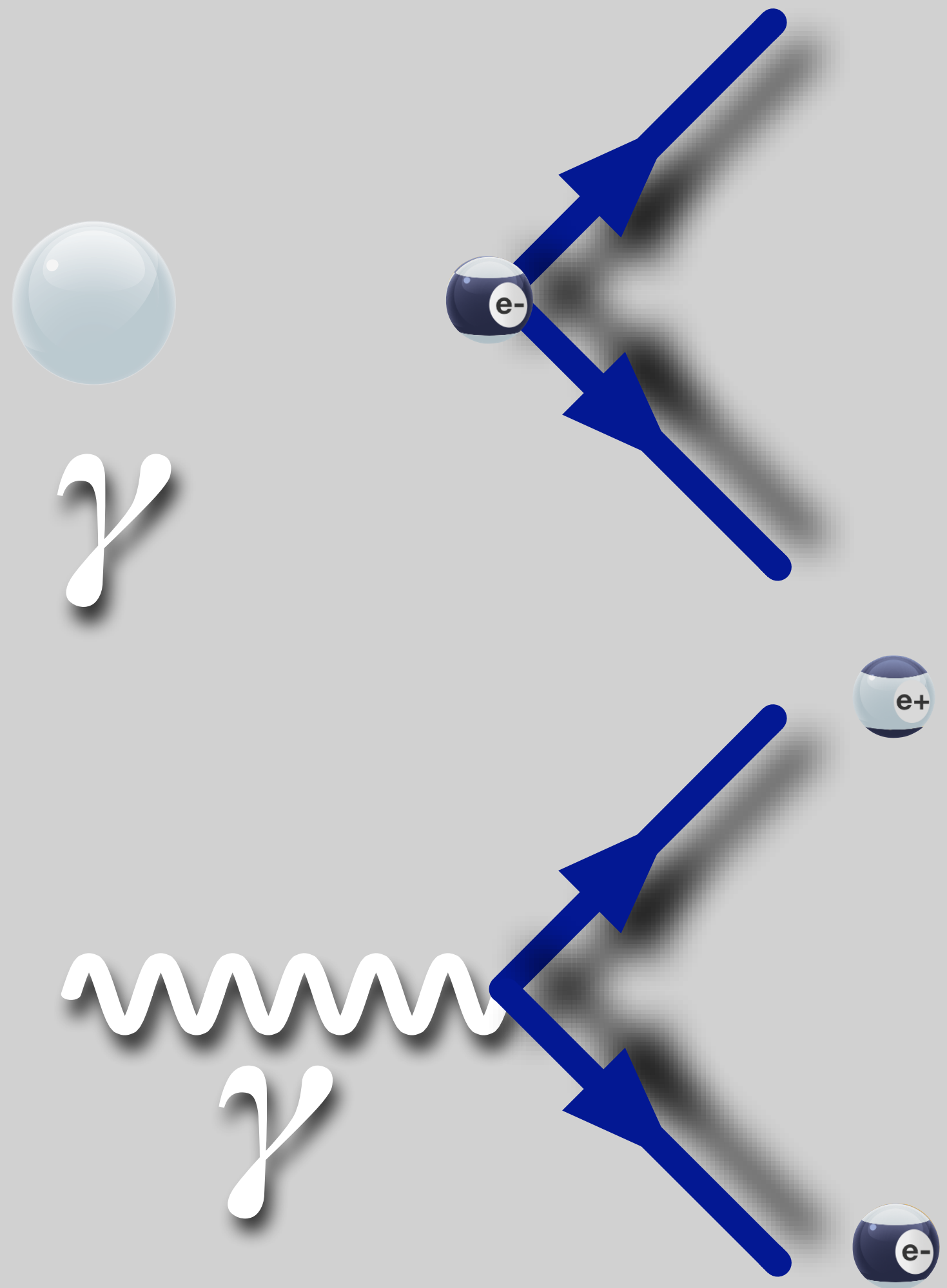


Yes...antimatter.

modern
interpretat



a photon
poof-disappears



The antimatter story has a
happy ending:

1932

Cosmic Rays

very high energy
protons from
space



~2 per minute per fingernail



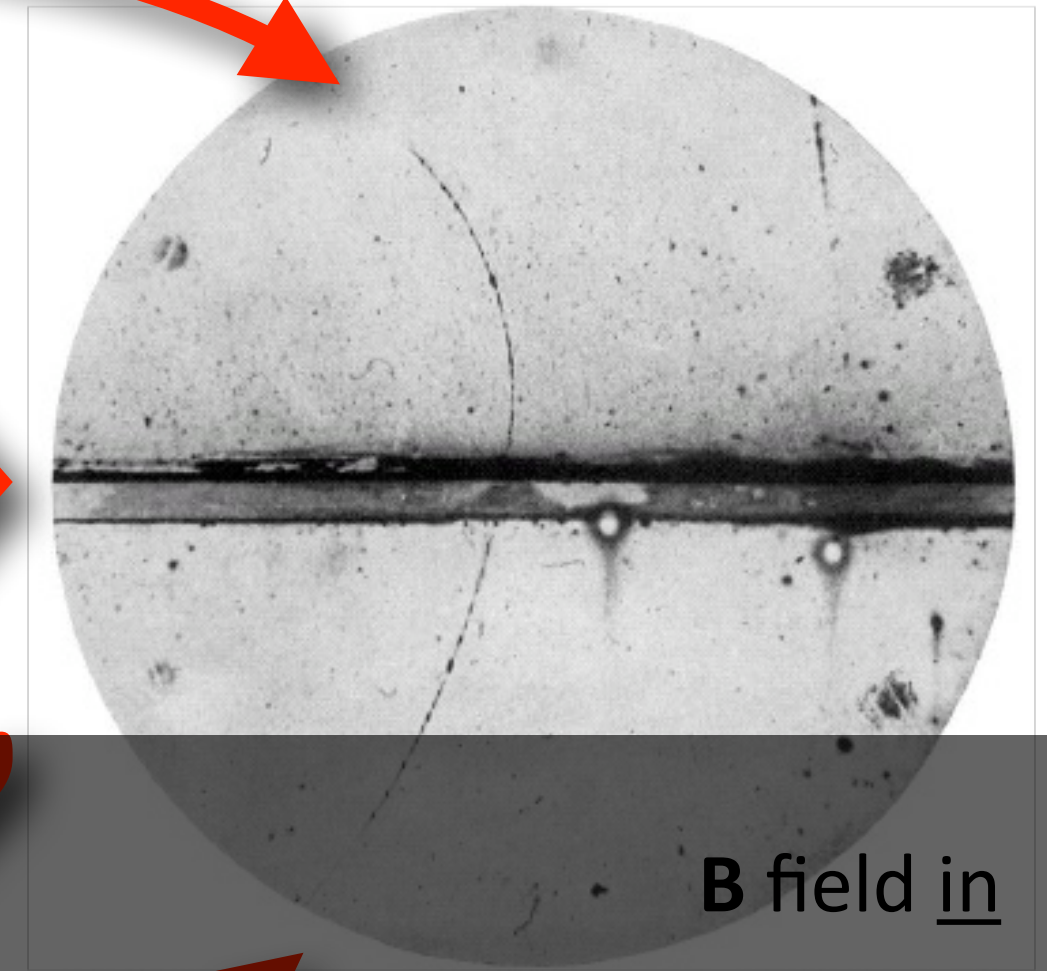
Carl Anderson

look at this track...

clever...put in a lead plate to cause particles to lose energy



sharper curvature at top



B field in

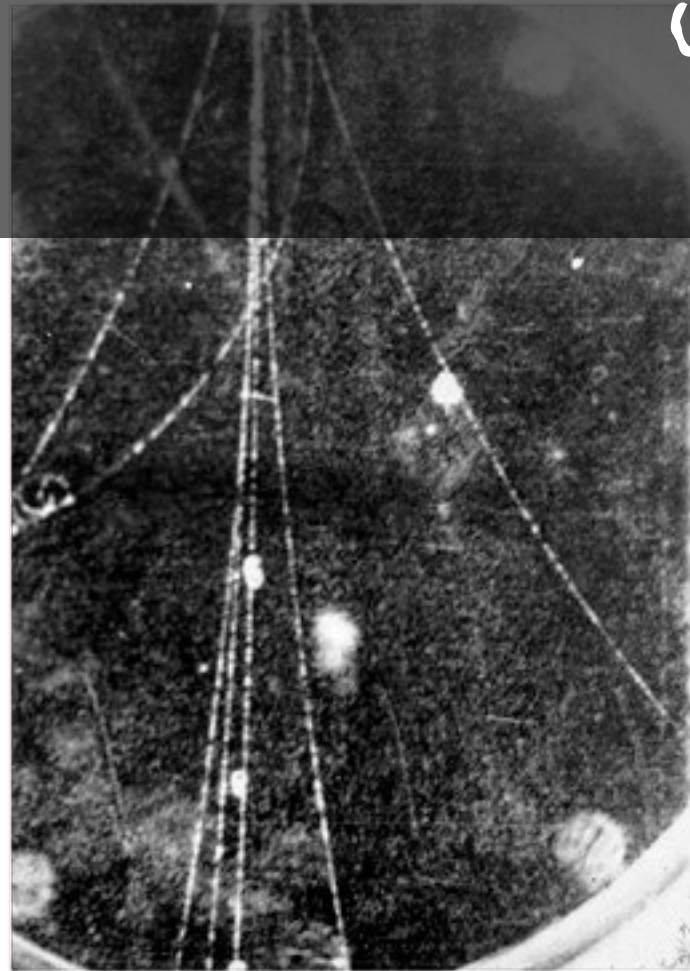
DOWN and negative?

UP and positive?

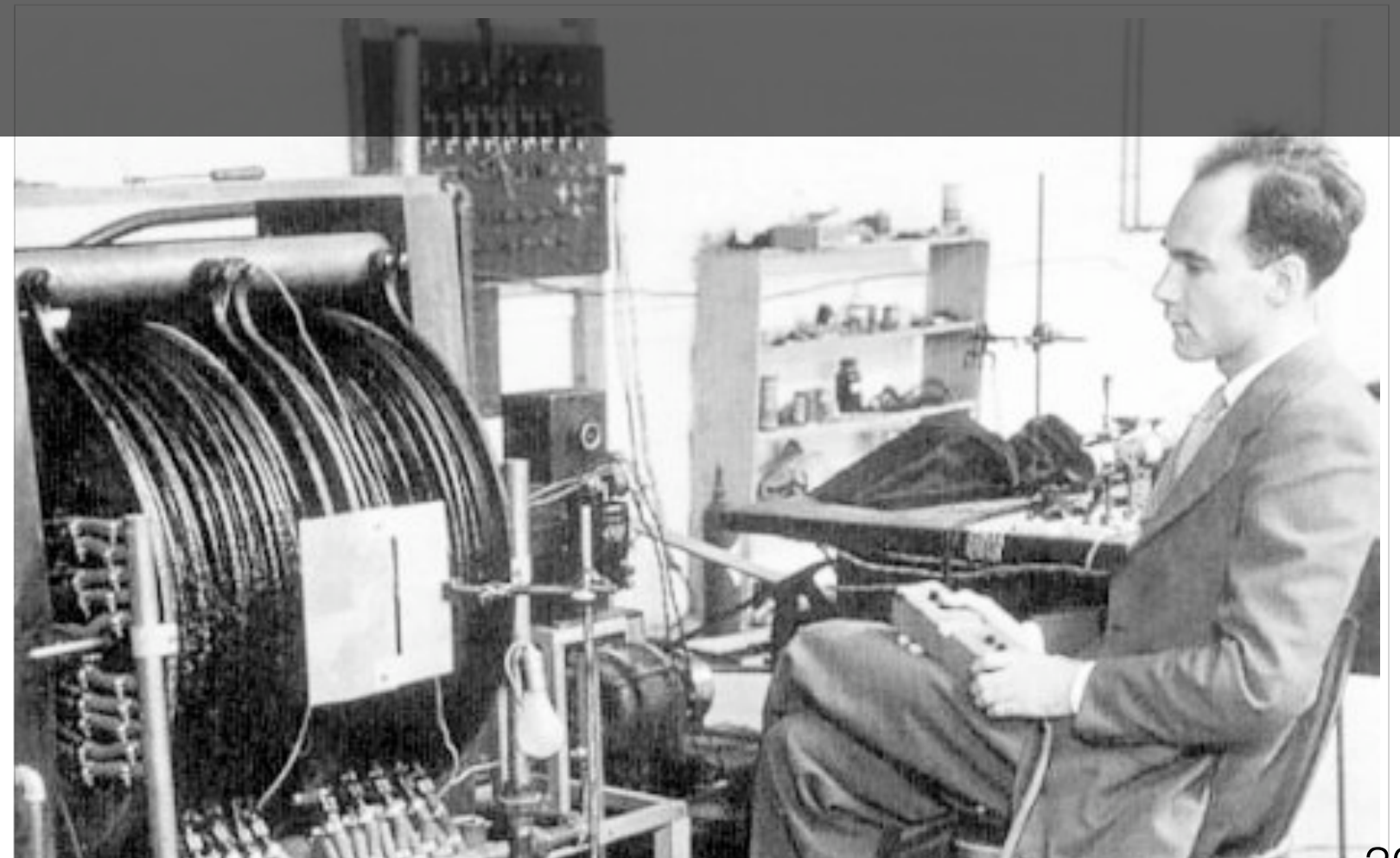


Yes...antimatter.

Right on schedule: 1932



B field in



the bar over the top will mean
“antiparticle”

anti-electron, aka “positron”

symbol:

\bar{e} or e^+

charge:

$+1e$

mass:

$m_e = 9.0 \times 10^{-31} \text{ kg} \sim 0.0005 \text{ p}$

spin:

$1/2$

category:

anti-fermion, anti-lepton

antimatter

is a fact of life

every particle has it's anti-particle partner

same mass, different electrical charge

Dirac Nobel

at the age of 31



Nobelprize.org
The Official Web Site of the Nobel Prize

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1933
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The Nobel Prize in Physics 1933
Erwin Schrödinger, Paul A.M. Dirac

The Nobel Prize in Physics 1933
Erwin Schrödinger
Paul A.M. Dirac



Erwin Schrödinger **Paul Adrien Maurice Dirac**

The Nobel Prize in Physics 1933 was awarded jointly to Erwin Schrödinger and Paul Adrien Maurice Dirac *"for the discovery of new productive forms of atomic theory"*

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Carl Anderson and Victor Hess

Anderson was 31

The screenshot shows the Nobelprize.org website. At the top, the logo and tagline "The Official Web Site of the Nobel Prize" are visible. A navigation menu includes "Nobel Prizes", "Alfred Nobel", "Educational", "Video Player", and "Nobel Organizations". The breadcrumb trail reads "Home / Nobel Prizes / Nobel Prize in Physics / The Nobel Prize in Physics 1936".

On the left, a sidebar menu lists various categories under "Nobel Prize in Physics", including "All Nobel Prizes in Physics", "Facts on the Nobel Prize in Physics", "Prize Awarder for the Nobel Prize in Physics", "Nomination and Selection of Physics Laureates", "Nobel Medal for Physics", "Articles in Physics", "Video Interviews", and "Video Nobel Lectures".

The main content area features a timeline from 1901 to 2012, with 1936 selected. Below the timeline, there are options for "Sort and list Nobel Prizes and Nobel Laur" and "Prize category: Physics".

The central focus is the page for "The Nobel Prize in Physics 1936", which lists the laureates: Victor F. Hess and Carl D. Anderson. Below their names are two black and white portrait photographs. Underneath the photos are their names: "Victor Franz Hess" and "Carl David Anderson".

A paragraph below the portraits states: "The Nobel Prize in Physics 1936 was divided equally between Victor Franz Hess 'for his discovery of cosmic radiation' and Carl David Anderson 'for his discovery of the positron'".

At the bottom of the page, there is a footer with the text "Nobelprize.org. 20 Mar 2013 laureates/1936/".



this is where it gets interesting

we need to establish a language for Dirac-like reactions

**“Relativistic Quantum Field Theory”
essentially invented by Paul Dirac**

notice a couple of things about what
appears in Dirac's equation

1. it's about more than one thing: two electrons and a
photon

"regular" Quantum Mechanics is about single objects only

2. stuff appears and stuff disappears

what's

nothing.

what's nothing

you'd maybe say:

no objects (particles...quanta)

zero energy

the Heisenberg Uncertainty Principle:

there's no state of Nature that can possess any precise value of, say, energy

and that includes Zero.

the Heisenberg Uncertainty Principle

will not allow a **void**.

but we *still* have a notion of the **vacuum**

it's the lowest energy state in Nature

where there are no real particles

understanding what's going on here

requires some mental fortitude

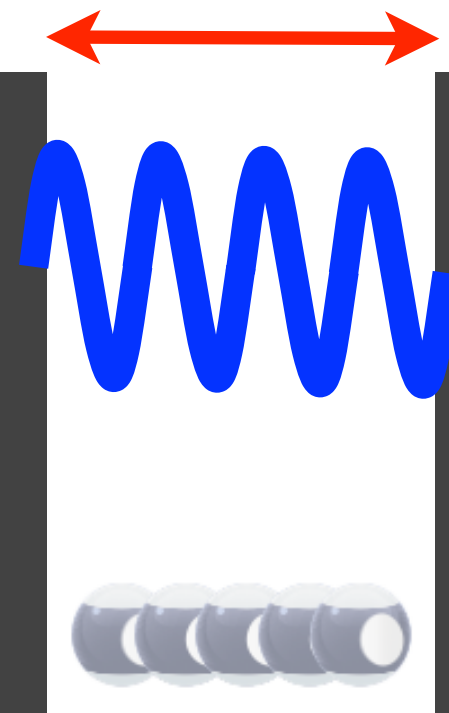
remember

trying to trap an electron?

let's make it all about nothing.



an electron...somewhere here:



make the trap smaller

remember

trying to trap an electron?

do nothing tighter



or an electron, somewhere here:



make the trap smaller to this value:

$$\Delta x \sim \frac{h}{m_e c}$$

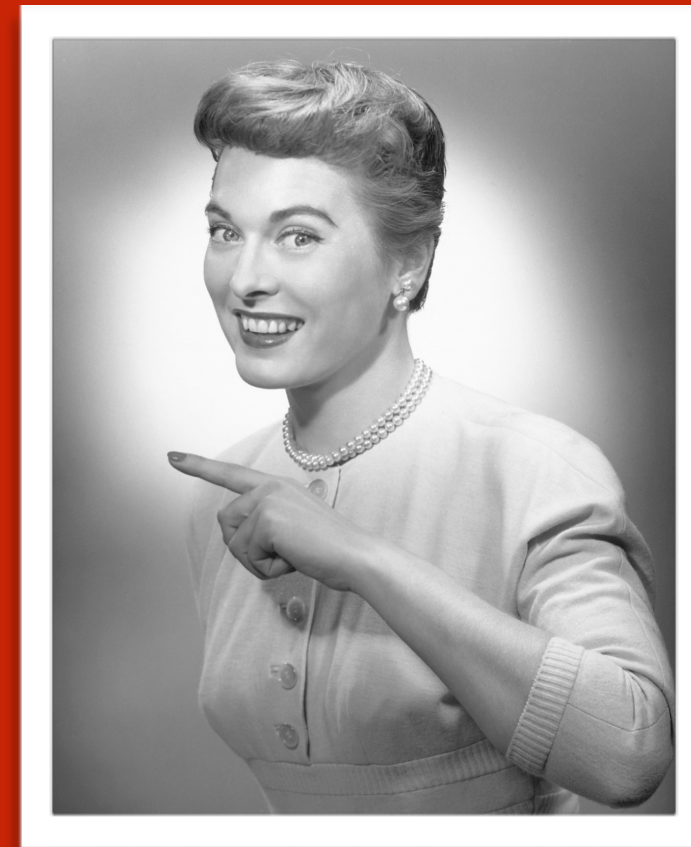
$$\sim 2.2 \times 10^{-12} \text{ m}$$

The size of a Hydrogen atom... $5 \times 10^{-11} \text{ m}$

The size of a proton... $\sim 10^{-15} \text{ m}$

an important

but simple calculation about
nothing



a very
important
“length”

Compton
Wavelength

we consider this to be
“the size of a
particle”

$$\Delta x \sim \frac{h}{m_e c} = \lambda_{\text{Compton}} = \lambda_C$$

$$\Delta x \Delta p \sim h$$

$$\frac{h}{m_e c} \Delta p \sim h$$

$$\frac{1}{m_e c^2} \Delta p c \sim 1$$

$$\Delta p c \sim m_e c^2$$

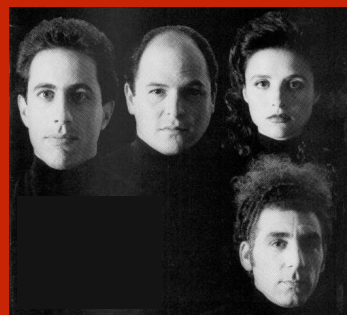
Remember: $E_T^2 = (m c^2)^2 + (p c)^2$

An energy equivalent to the mass energy...all by looking
closely at nothing.

remember

trying to trap an electron?

do nothing tighter



What's in Nothing with an electron?



electron + nothing, somewhere here:



make the trap smaller to this value:

$$\Delta x \sim \frac{h}{m_e c} = \lambda_C$$

$$\sim 2.2 \times 10^{-12} \text{ m}$$

but wait.

let's just do this in space...shrink to this critical size

the same thing happens

remember

trying to trap an
electron?

do nothing tighter



What's in Nothing?



or **NOTHING**, somewhere here:



make the trap smaller to this value:

$$\Delta x \sim \frac{h}{m_e c} = \lambda_C$$

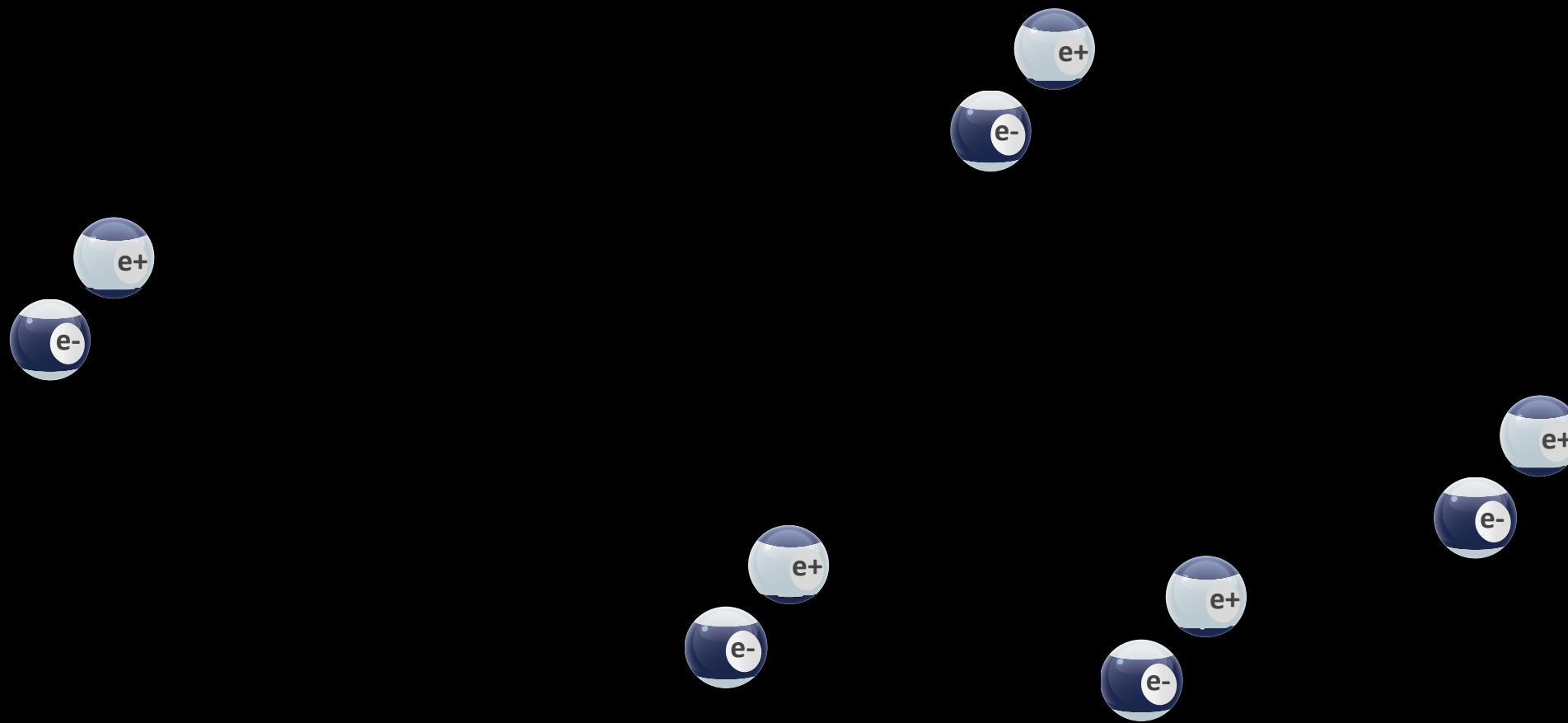
$$\sim 2.2 \times 10^{-12} \text{ m}$$

pop

the Uncertainty Principle requires

that particle-antiparticle pairs pop into and out of existence

all the time



uncertainty principle

+ the particular length of:

$$\lambda_C = \frac{h}{mc}$$

makes the vacuum very active.

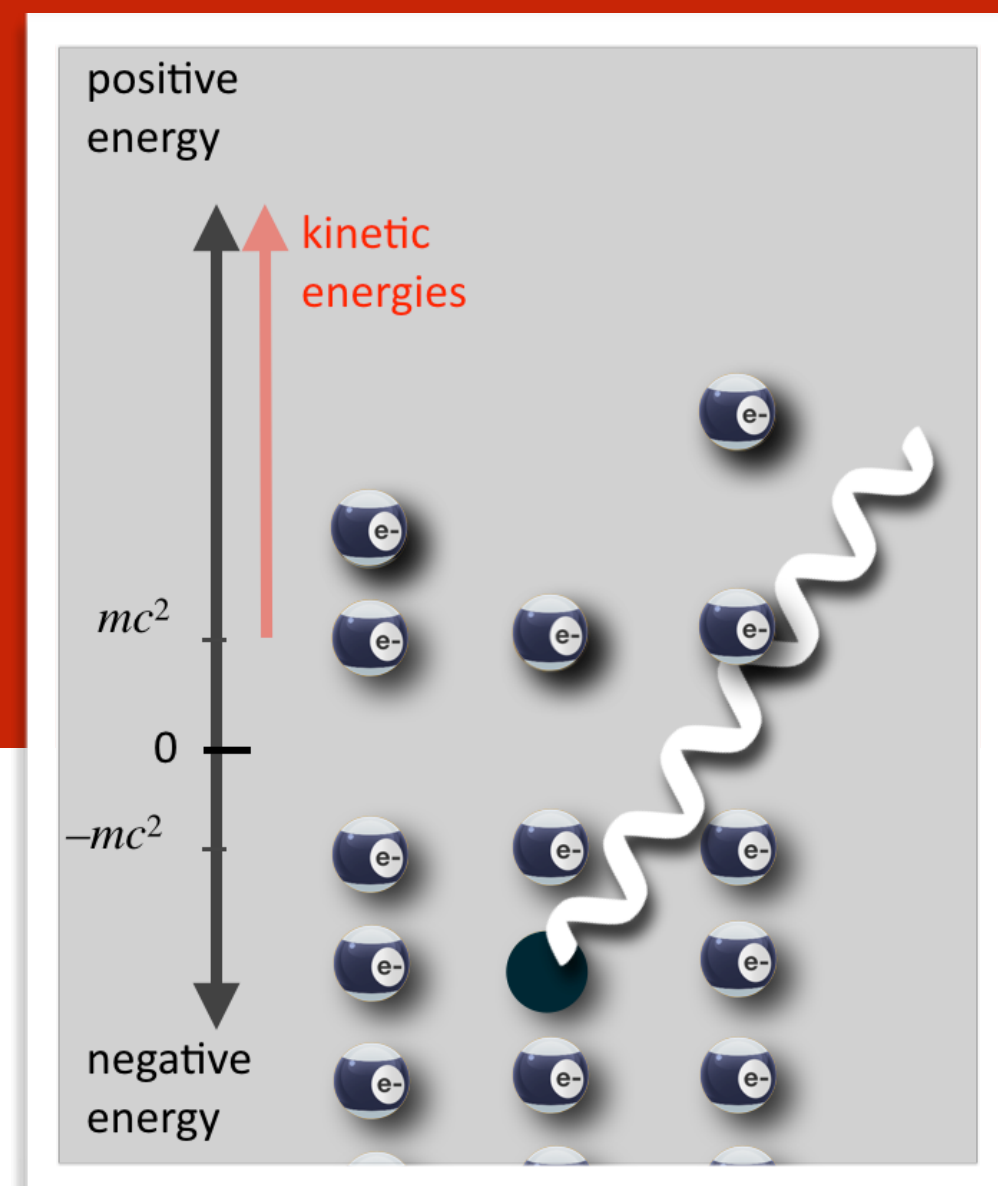


they are all popped out of the same stuff

the vacuum Field of the electron

electrons appear because they're coerced out of the vacuum

like by a photon



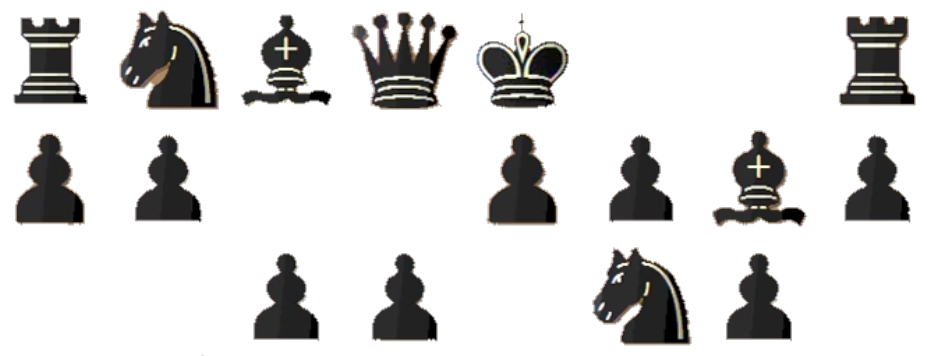
the quantum vacuum

a word about theories

let's play chess...

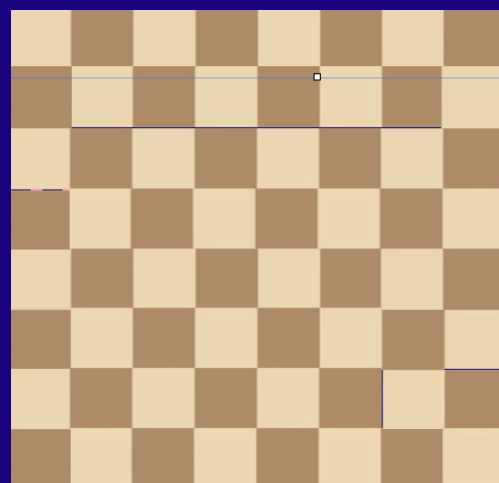
my model of chess

watch tons of matches



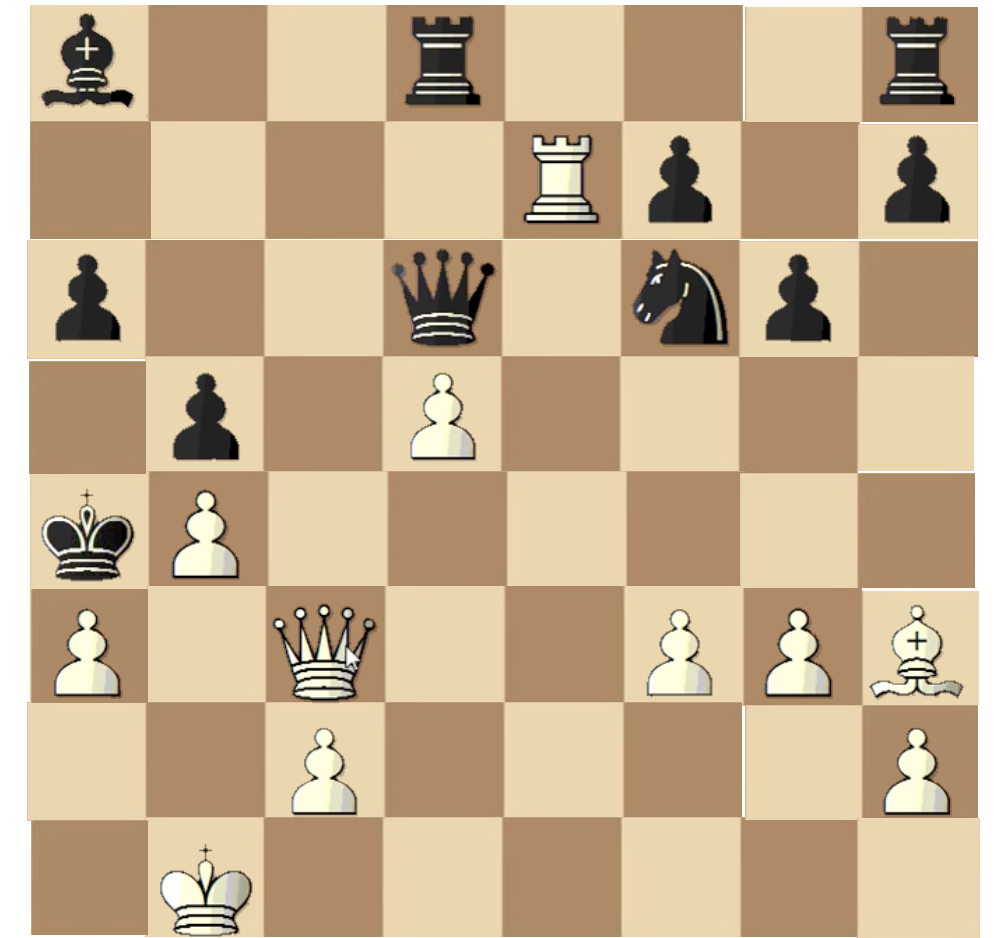
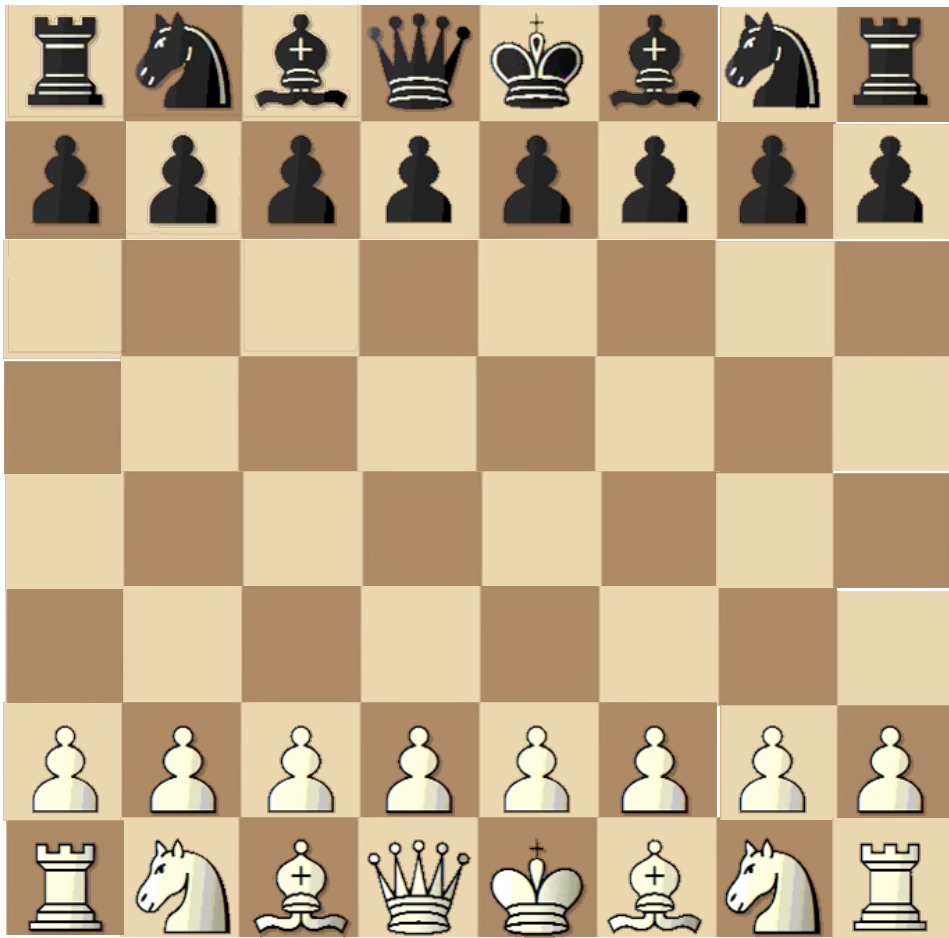
my model requires

the existence of a new entity



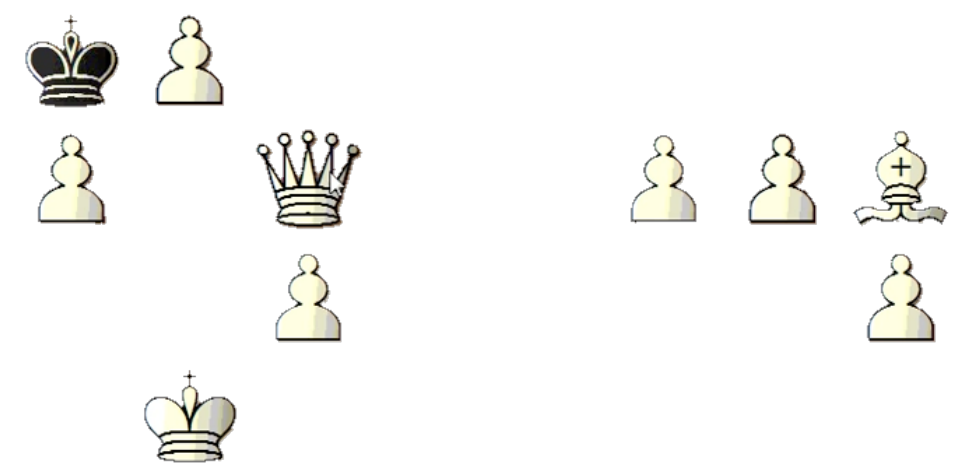
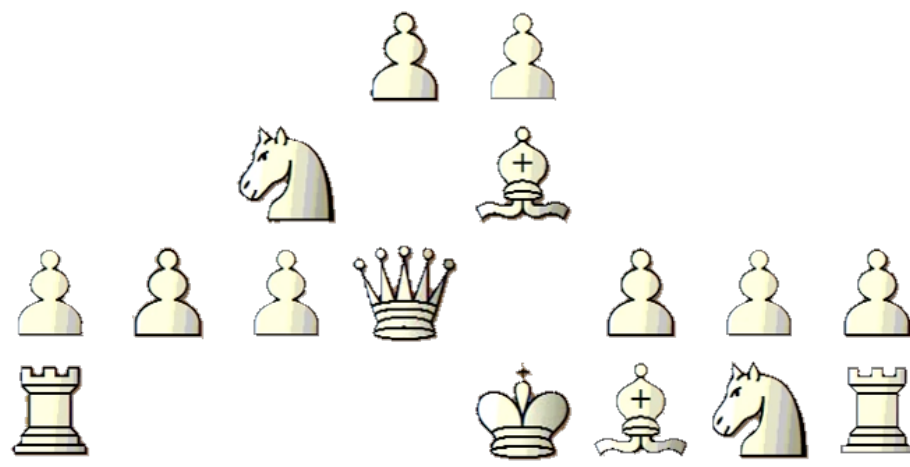
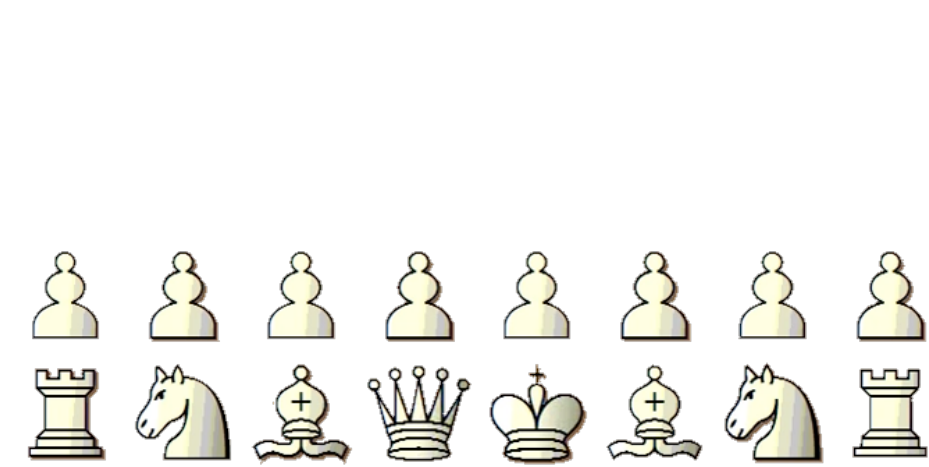
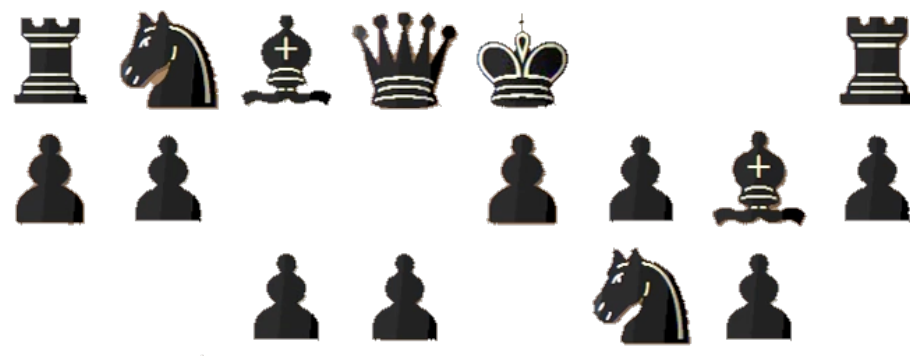
my model of chess

only with the board do the rules make sense



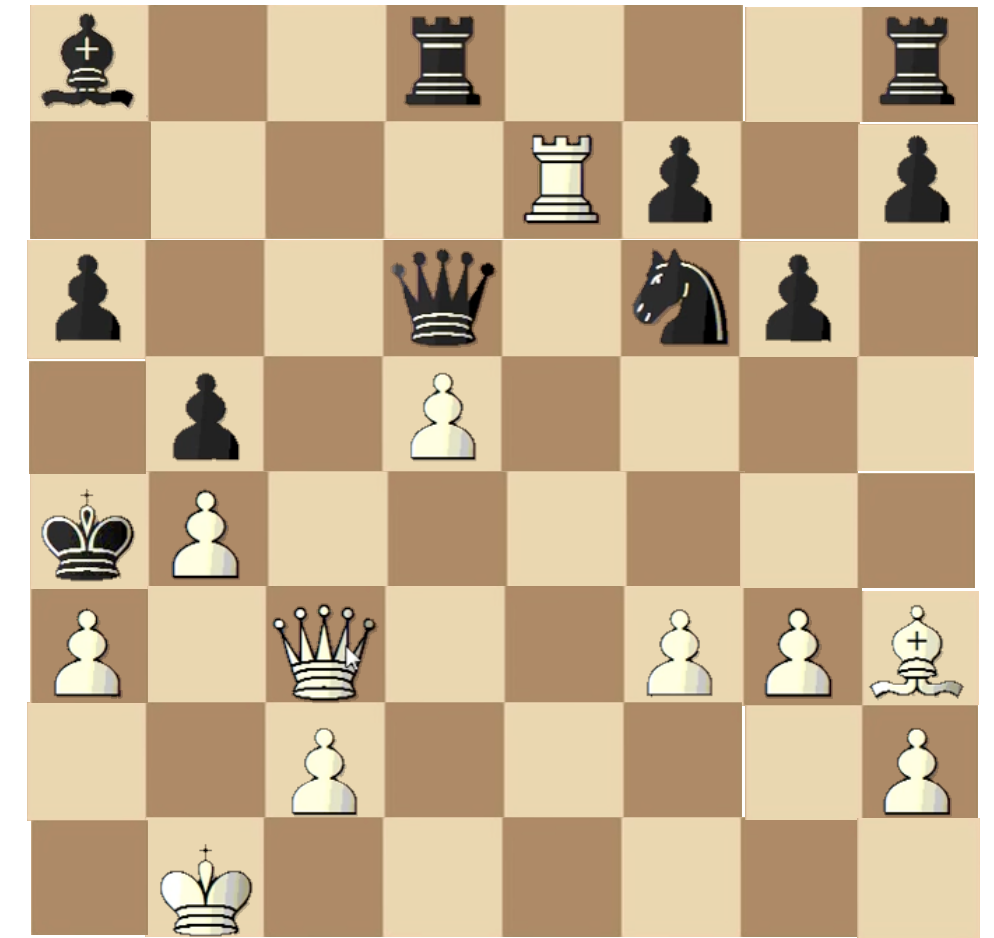
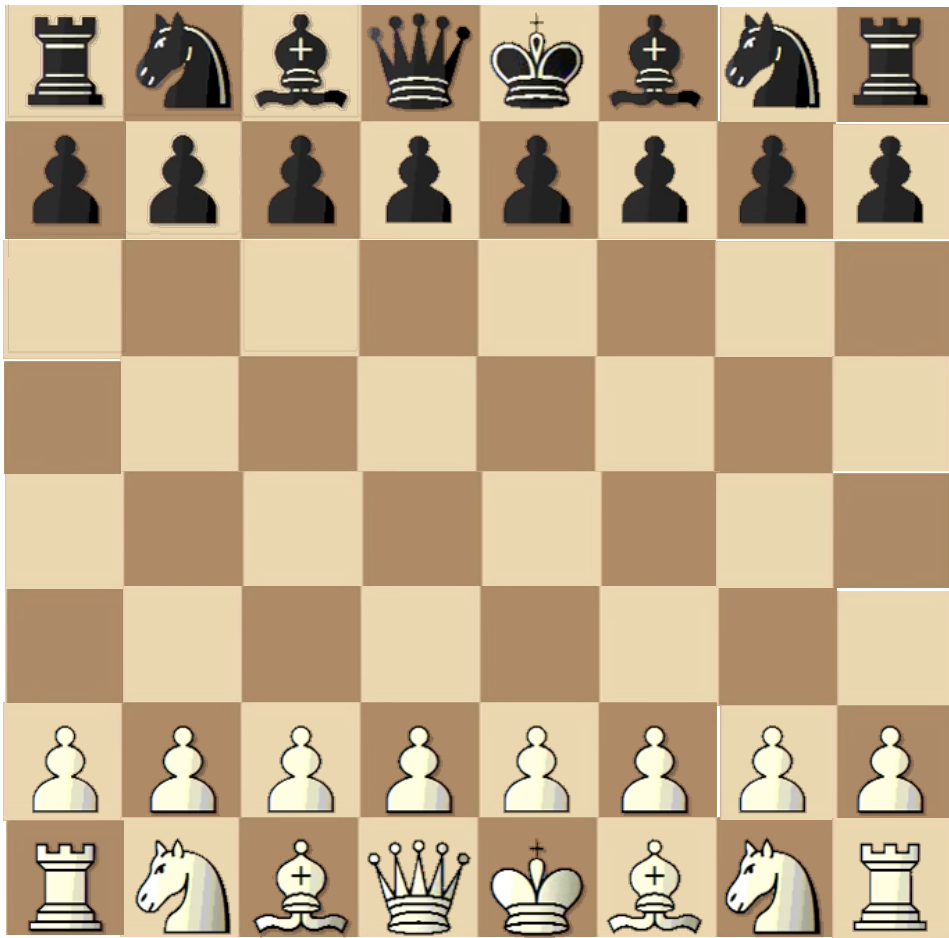
remember

what I see are the pieces



remember

what I need to be the case...is the board



The technical description:

"if it walks like a  and it quacks like a , then it must be a 

"



"

a successful physics model that requires an additional commitment!



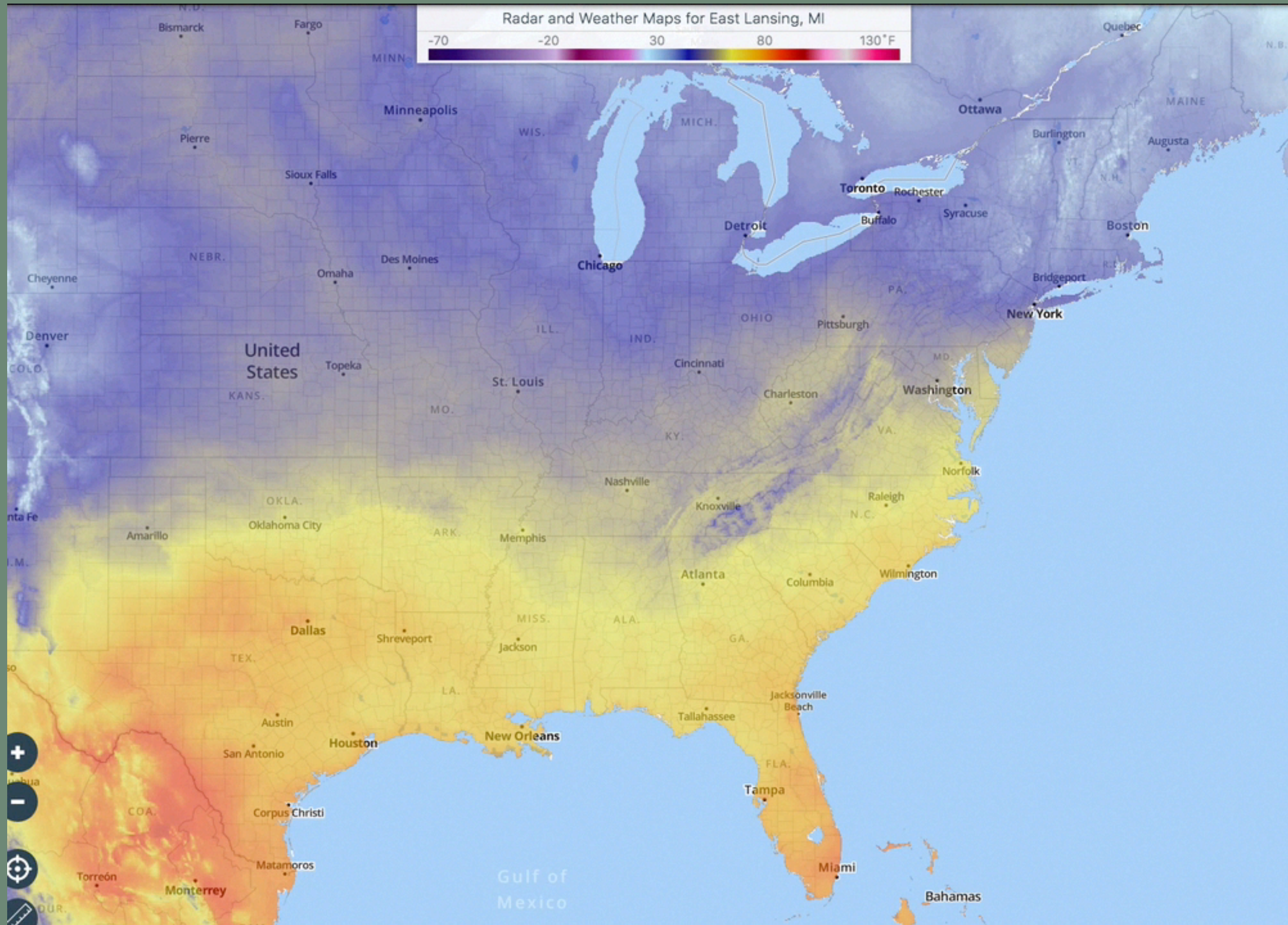
what about

fields?

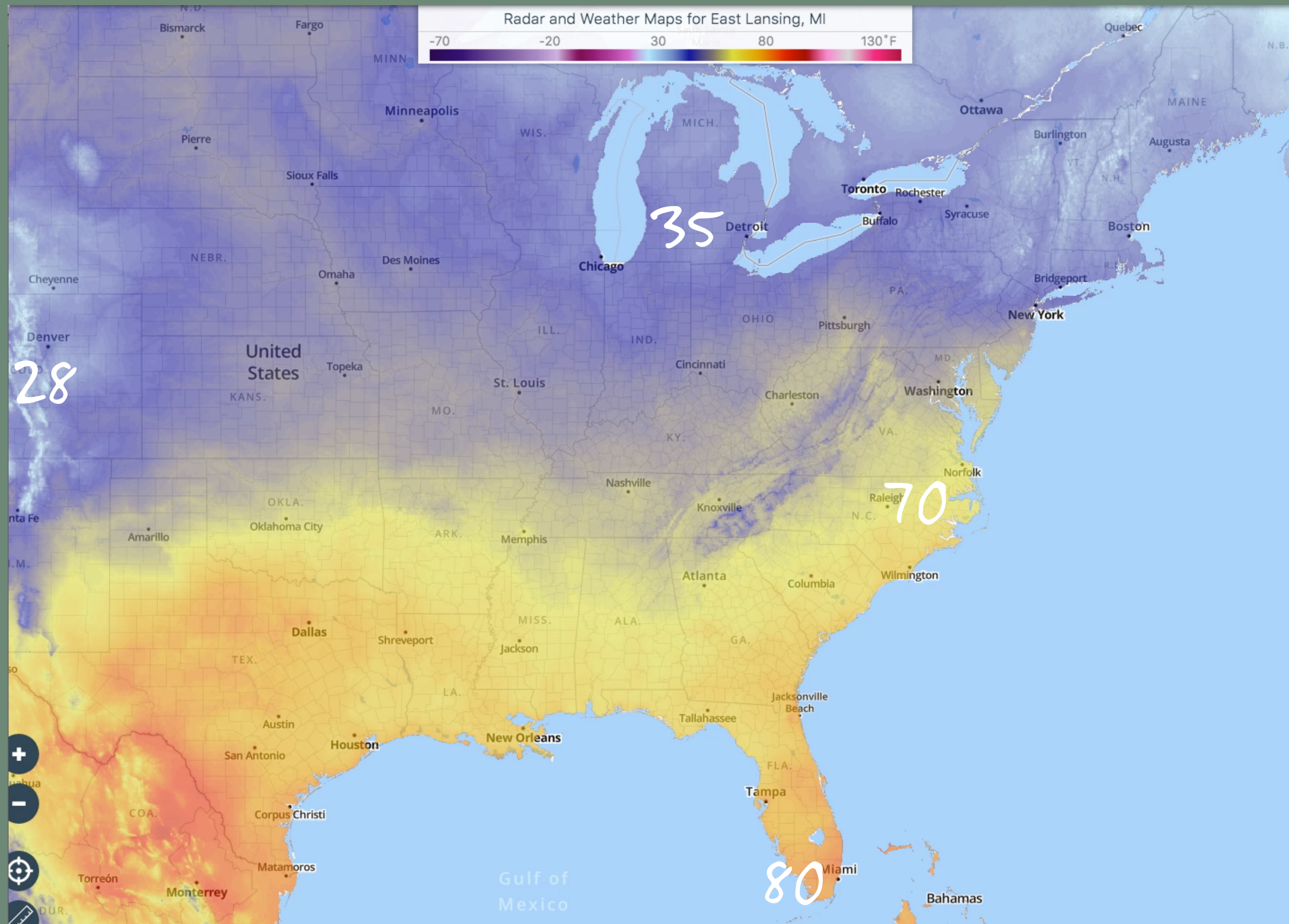
fields

a number in space

you know one
everywhere...a number

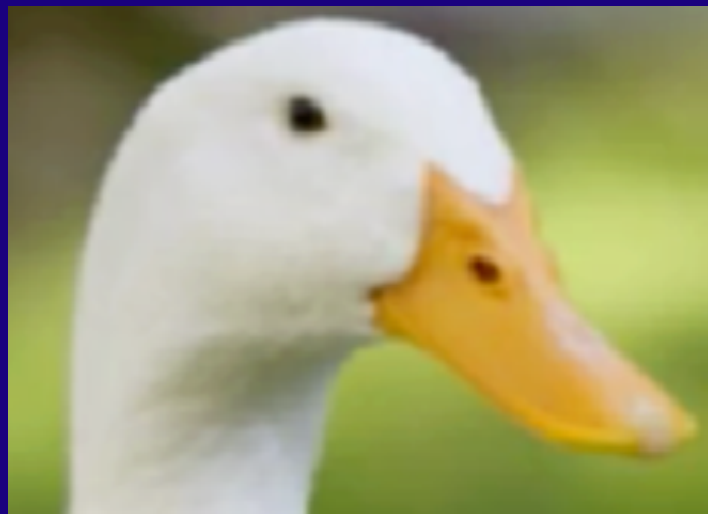


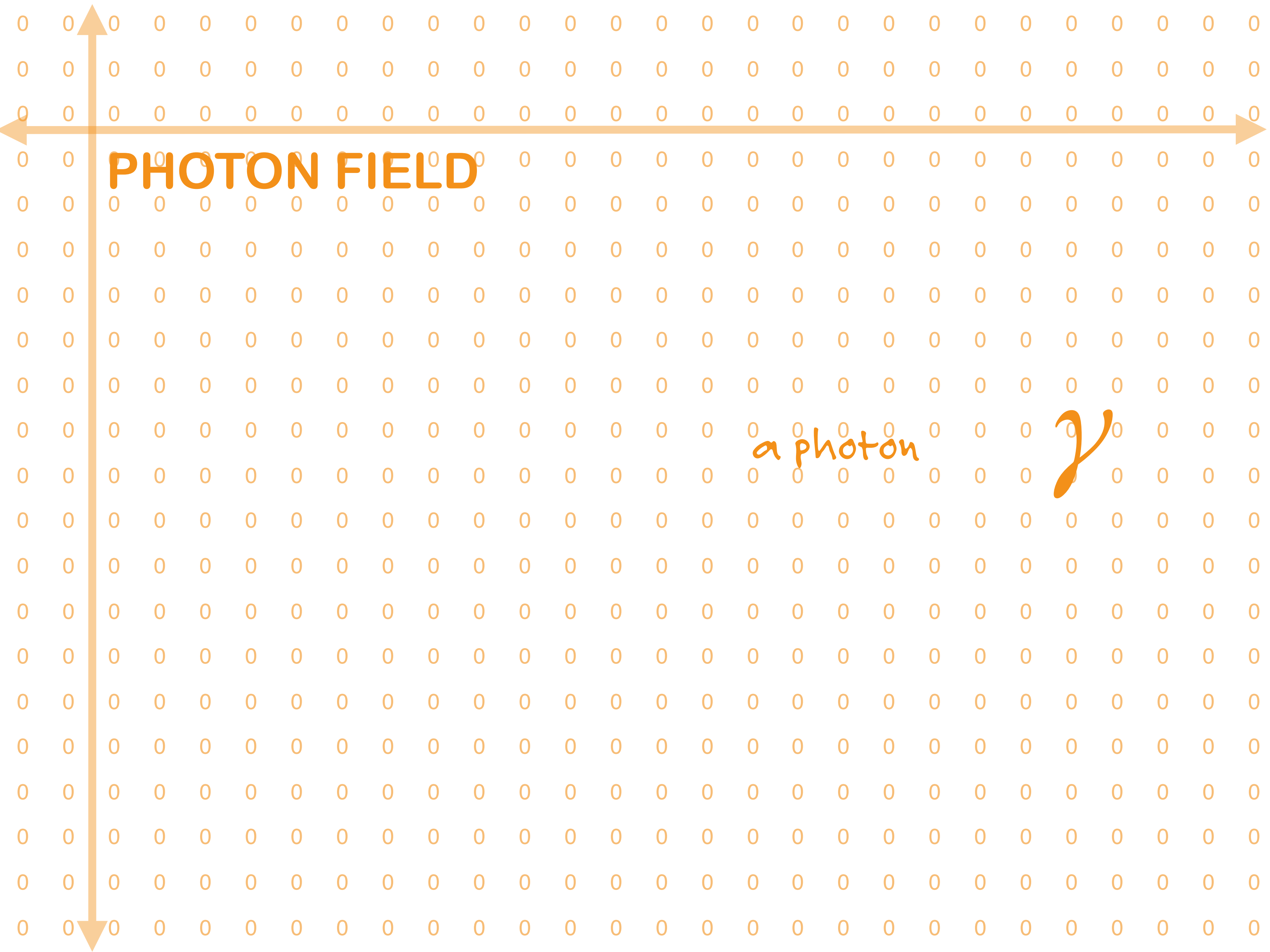
you know one
everywhere...a number



what's a particle?

it's localized wave in a field



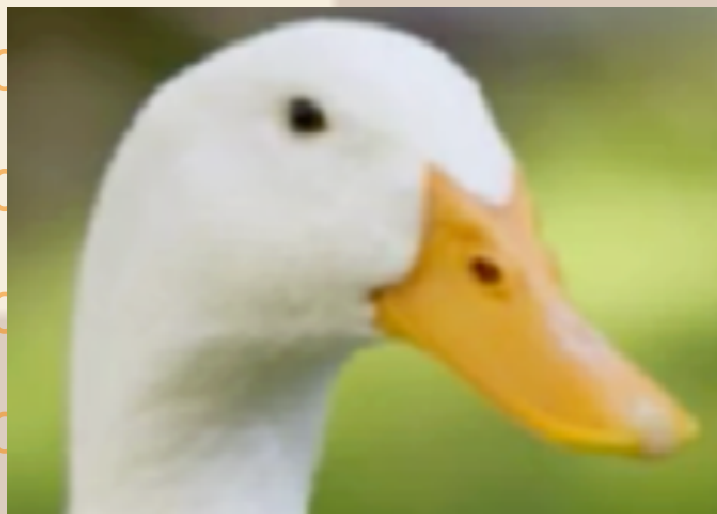


PHOTON FIELD

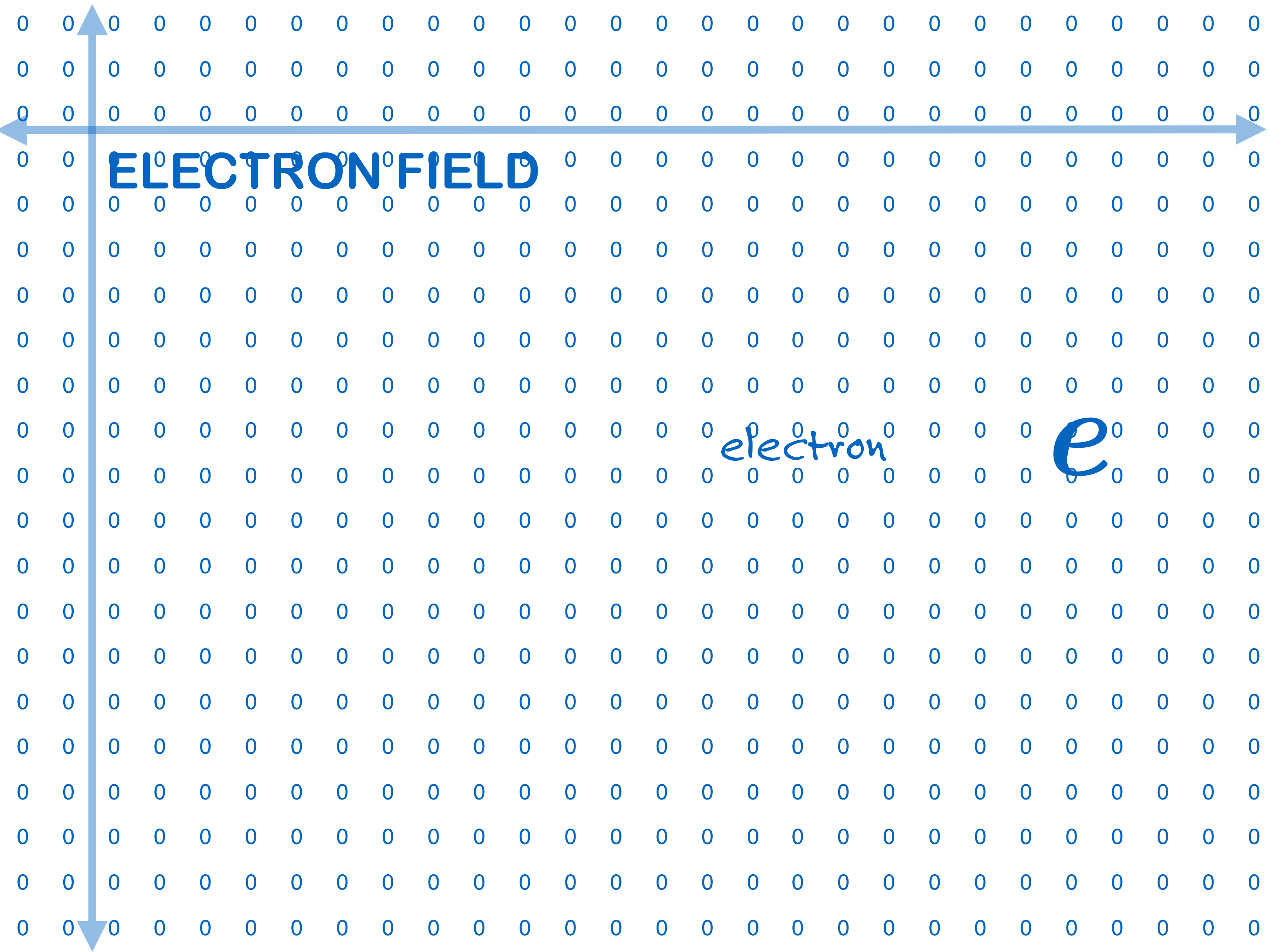
a photon

γ

PHOTON FIELD



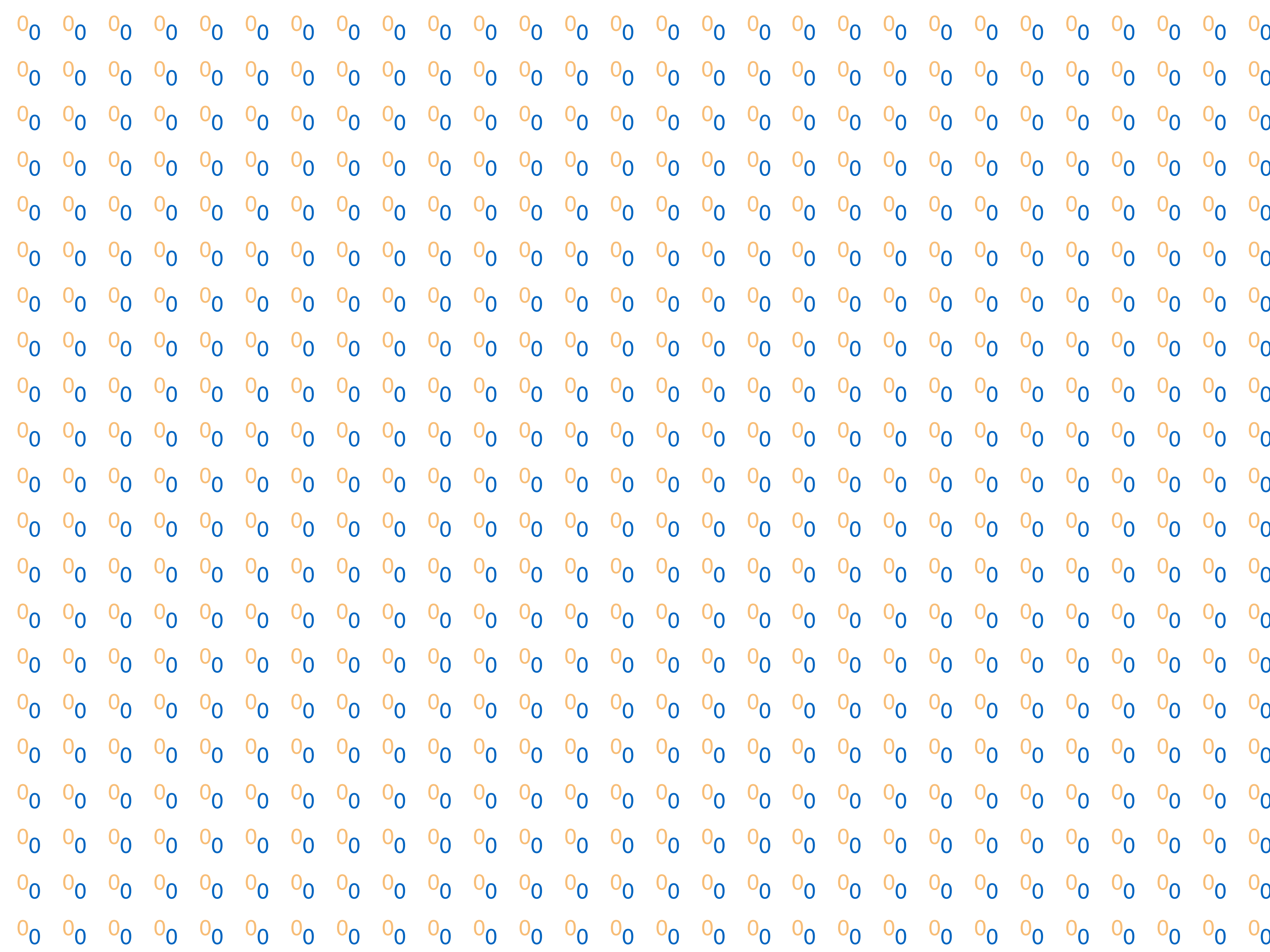
a photon γ



ELECTRON FIELD

electron

e



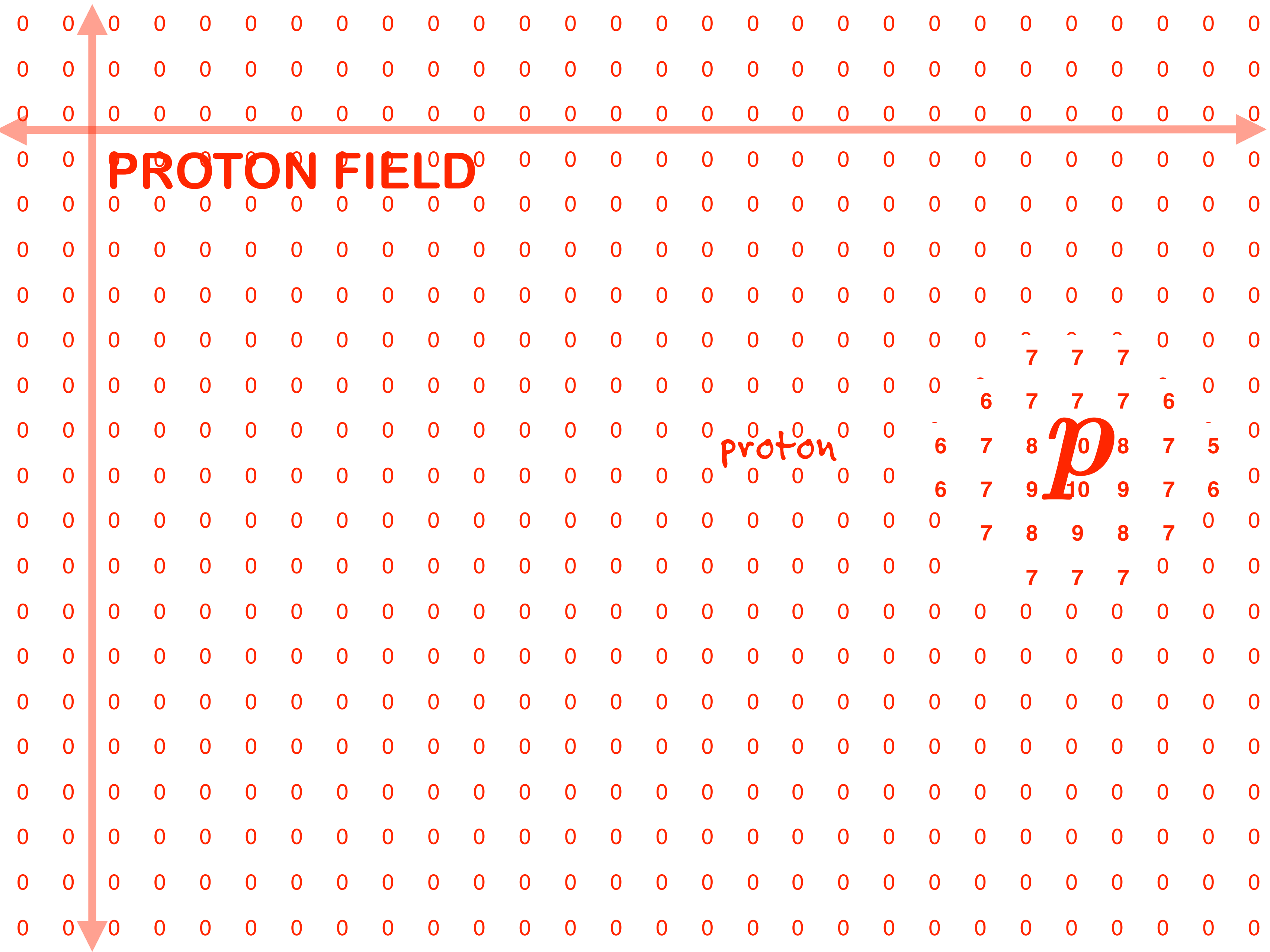
-2	0	0	-1	-2	-1	-2	2	1	-1	-2	1	2	0	-1	1	1	1	0	-1	-2	-2	-2	1	1	2	-1	0
0	-2	2	0	1	-1	-1	0	-2	1	2	2	-2	1	1	1	-2	-1	0	0	-1	-2	-1	0	-2	1	0	2
1	0	-2	-1	0	-2	-1	1	2	-2	2	1	0	2	1	2	2	1	-2	2	2	-2	0	2	2	1	2	0
1	0	0	-1	0	2	2	-1	-2	0	-2	2	1	2	-1	-1	-1	-1	-2	0	-2	-1	2	1	-1	-1	1	0
-1	-1	1	2	-2	-1	-1	-2	2	1	1	2	0	-2	-2	1	2	-1	2	-1	0	1	0	-2	0	1	2	2
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vacuum fluctuations

here's how

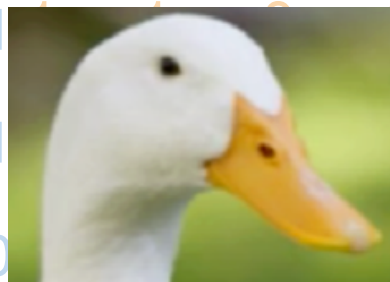
stuff happens

in this particle field theory model



electron

photon field "disturbance"



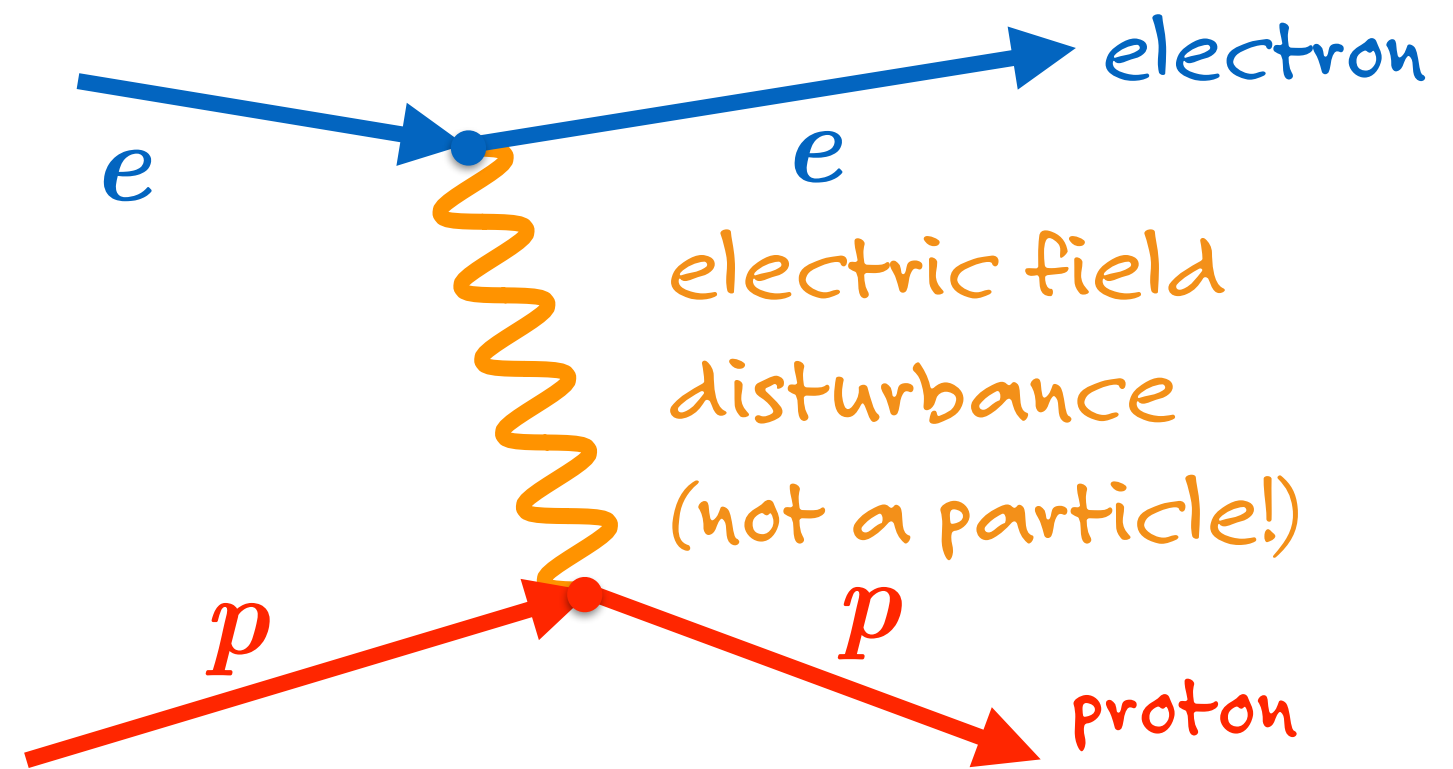
time



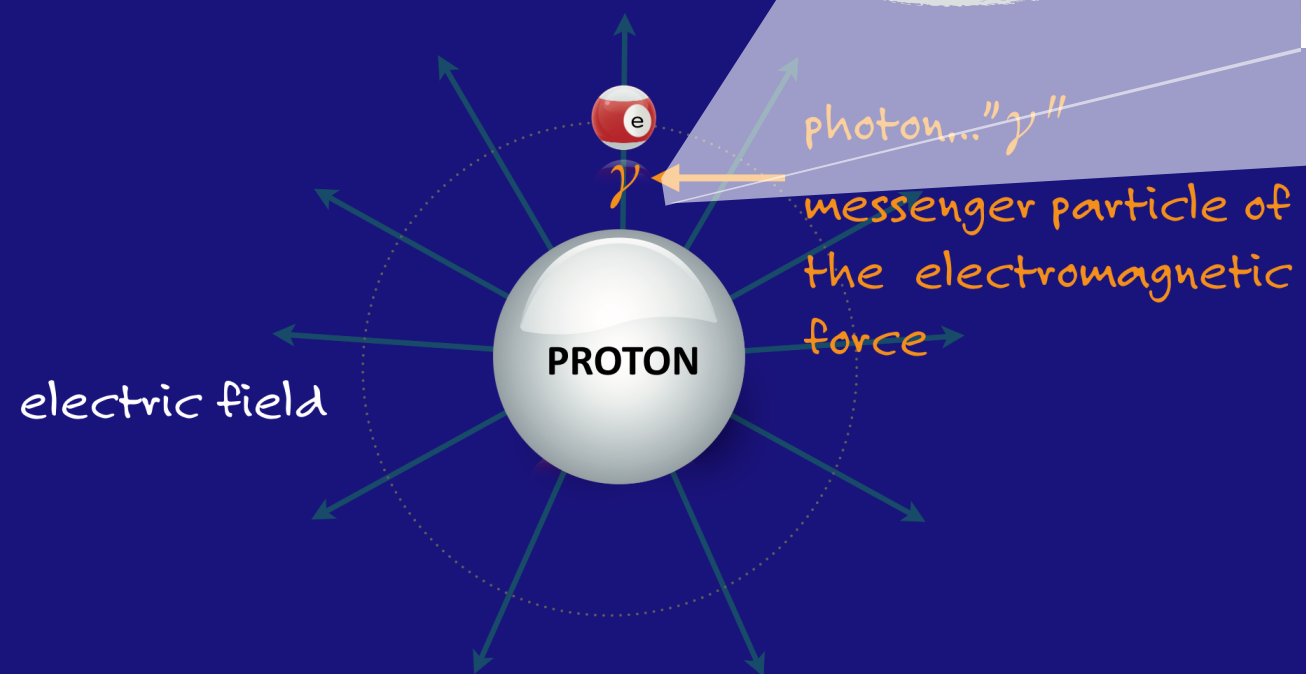
proton

our atom

2nd way



forces?
from particles, 1st way



particle field theory*

the best theory in history

never an
incorrect
prediction



outrageously
precise
agreement,
prediction and
measurement

*Quantum Electrodynamics

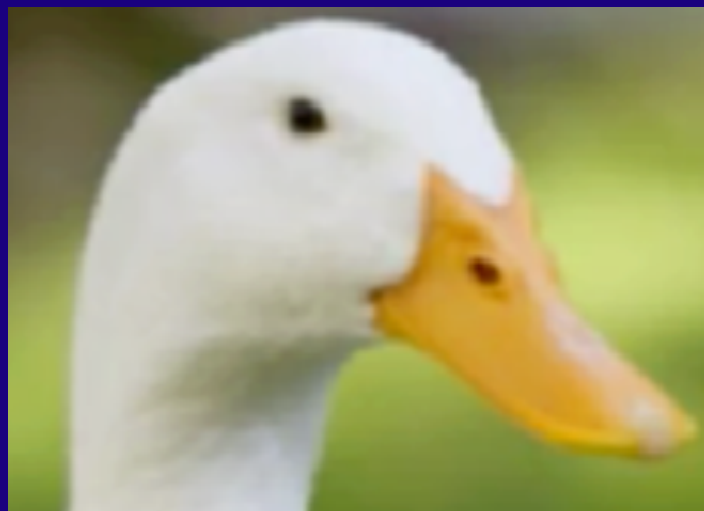
what's more fundamental?

a winner

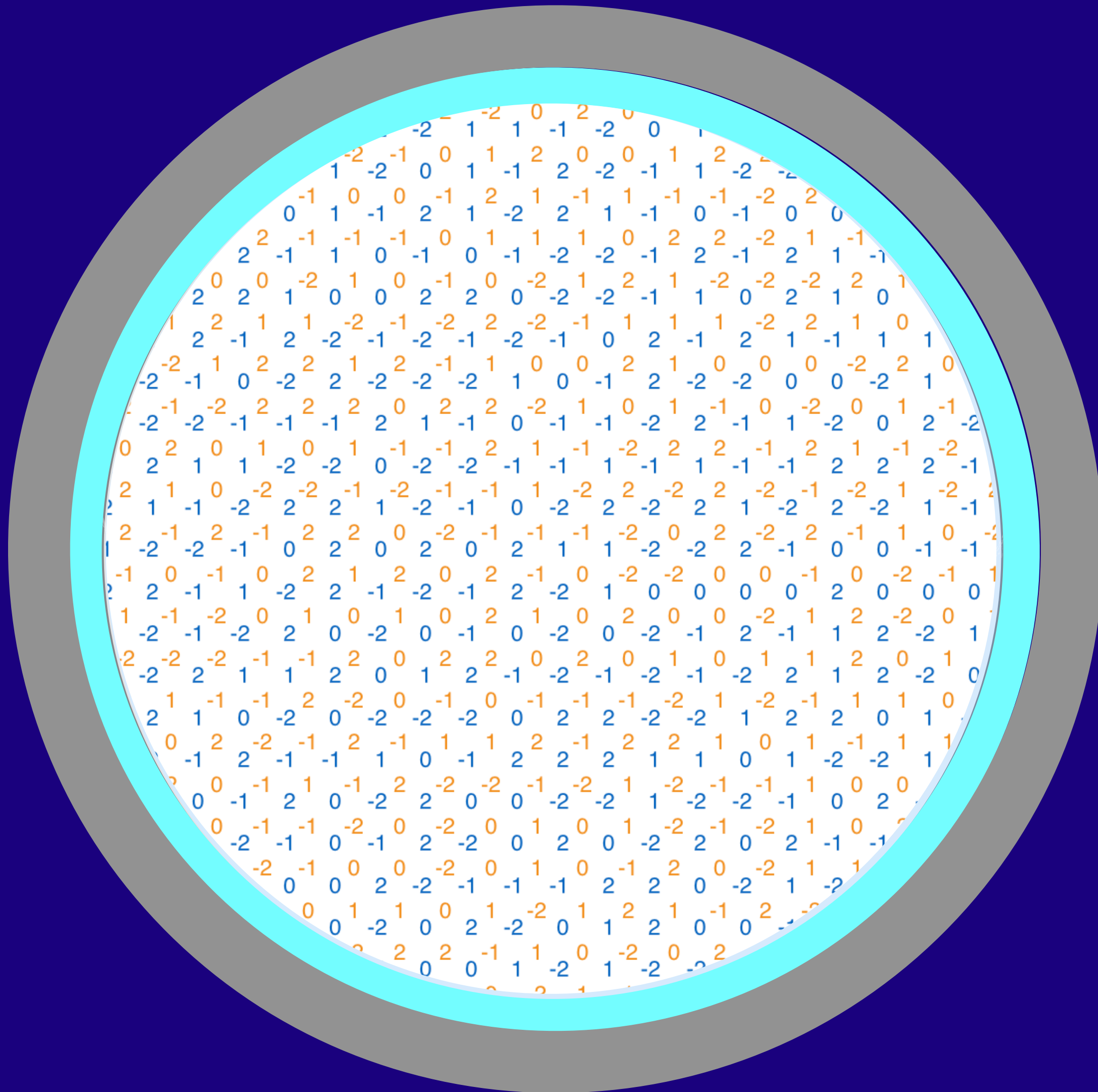


particles

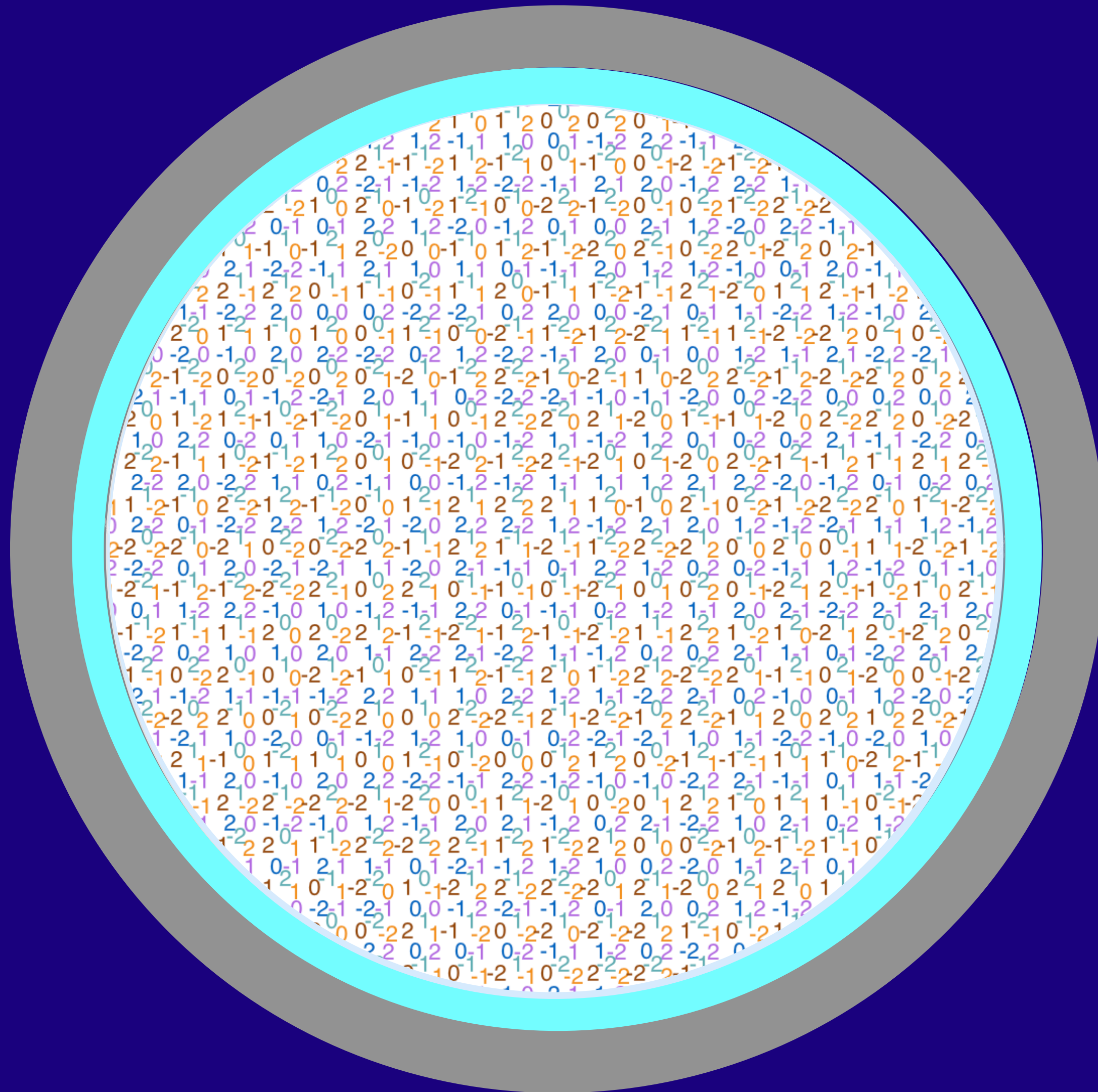
fields



the particle vacuum full of fields:



the particle vacuum full of fields for every "particles"



this has a name

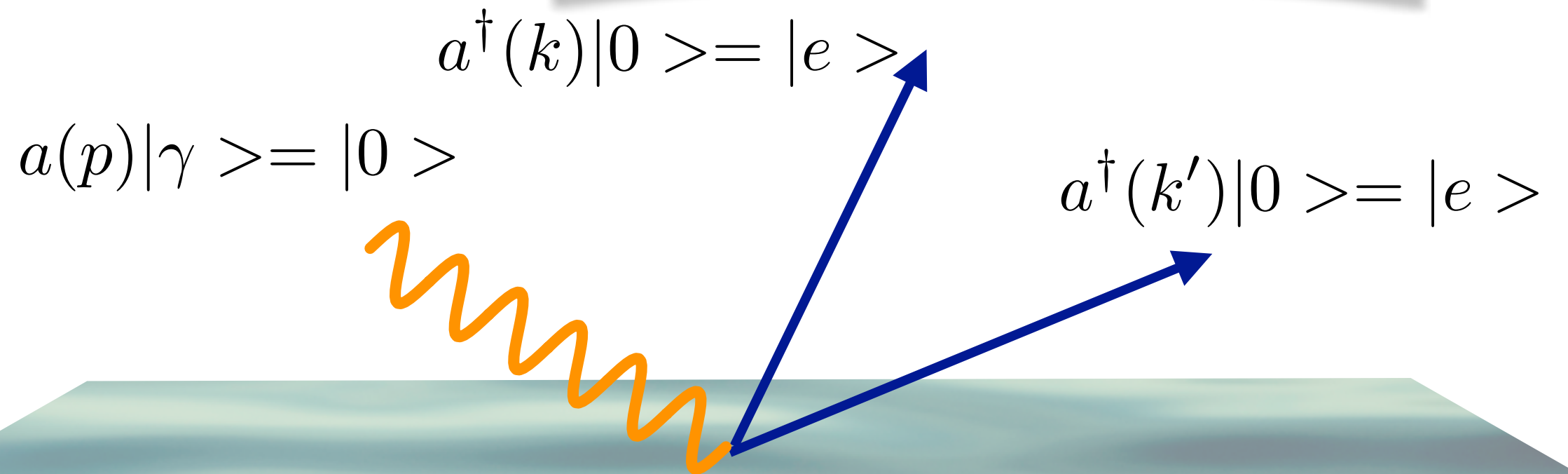
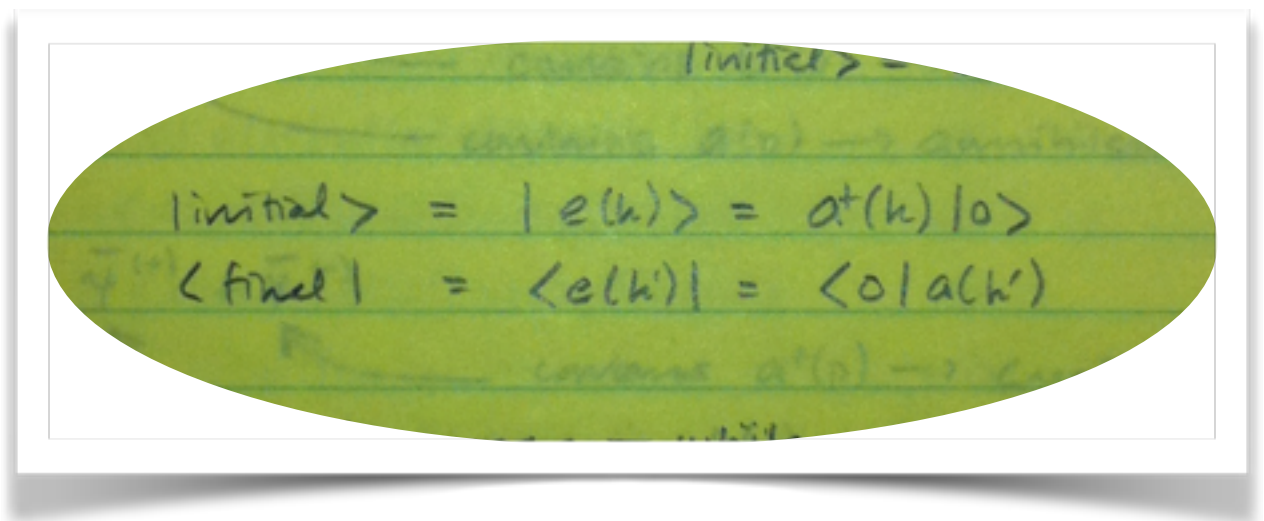


"the worst prediction in the history of physics"



a little
more
specific

what the
mathematics tells
us



it's not like the photon is now "in" the electron

the photon pops the electron- positron pair out of the Ur
electron field

and itself disappears back into the Ur photon field.

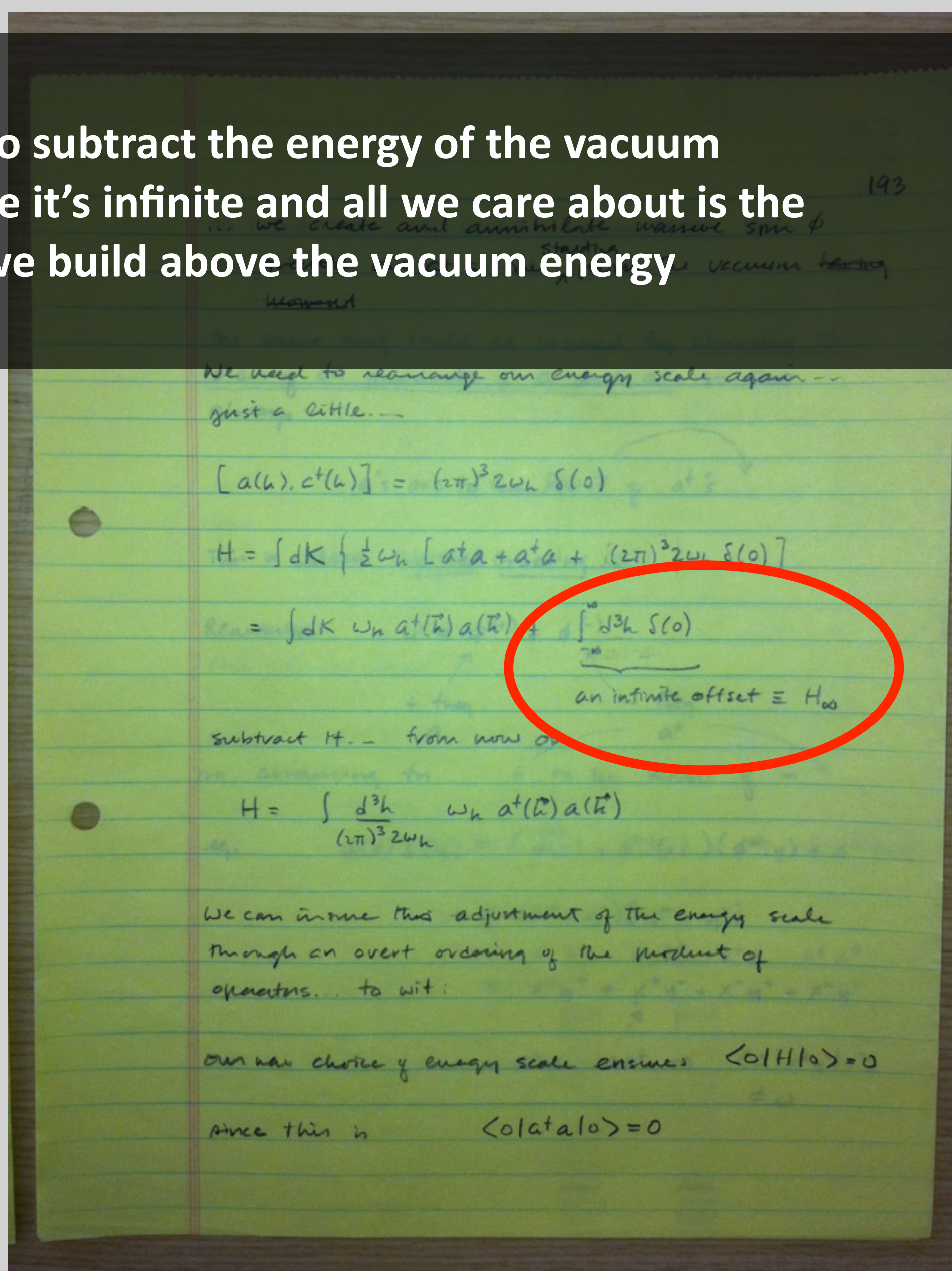
but what
does that
mean?

we have to subtract the energy of the vacuum
away... because it's infinite and all we care about is the
states we build above the vacuum energy

it means that the
vacuum is full of
energy

like a reservoir

particles are
created out of the
vacuum



Okay. So we don't like infinity

we subtract it away and worry about the difference
between

infinity and finite energies of real particles

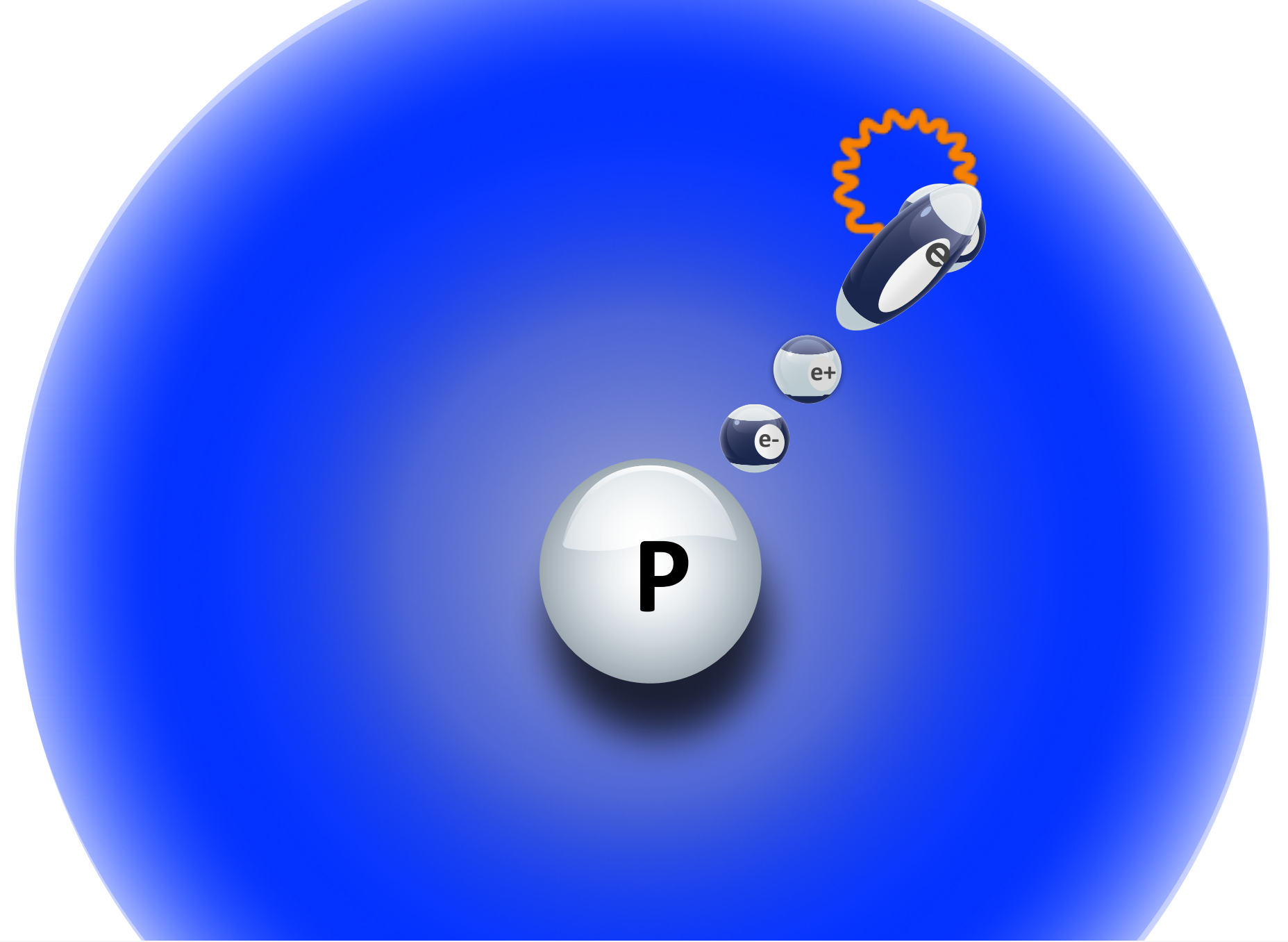
Seriously?

That seems to be the case.

This picture works with exquisite precision and accuracy.

but the
vacuum

is roiling with
particle-
antiparticle
“virtual pairs”
popping into and
out of existence
multiple ways we
know this.



A “regular” model of the hydrogen atom...needs
modification to take into account the effects of the vacuum

The electron cloud is spread out by the virtual
photon and the positron’s effects...and that changes
the emission spectrum of hydrogen:
The “**Lamb Shift**”...measured after WWII with
microwave technologies

1955 Nobel Prize

Willis E. Lamb

died just a few years
ago at the University
of Arizona




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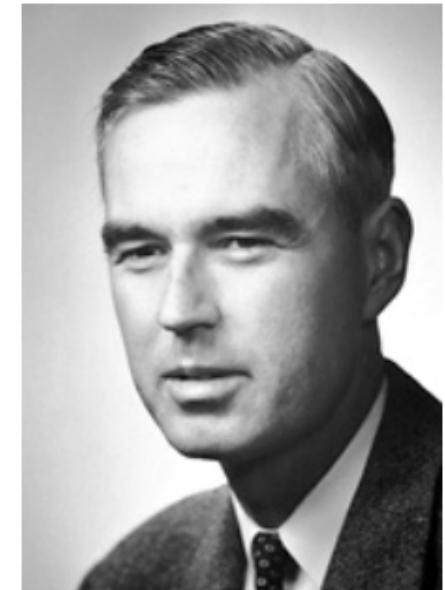
1901 2012

Sort and list Nobel Prizes and Nob Prize category:

 **The Nobel Prize in Physics 1955**
Willis E. Lamb, Polykarp Kusch

The Nobel Prize in Physics 1955

- Nobel Prize Award Ceremony
- Willis E. Lamb
- Polykarp Kusch



Willis Eugene Lamb Polykarp Kusch

The Nobel Prize in Physics 1955 was divided equally between Willis Eugene Lamb "for his discoveries concerning the fine structure of the hydrogen spectrum" and Polykarp Kusch "for his precision determination of the magnetic moment of the electron".

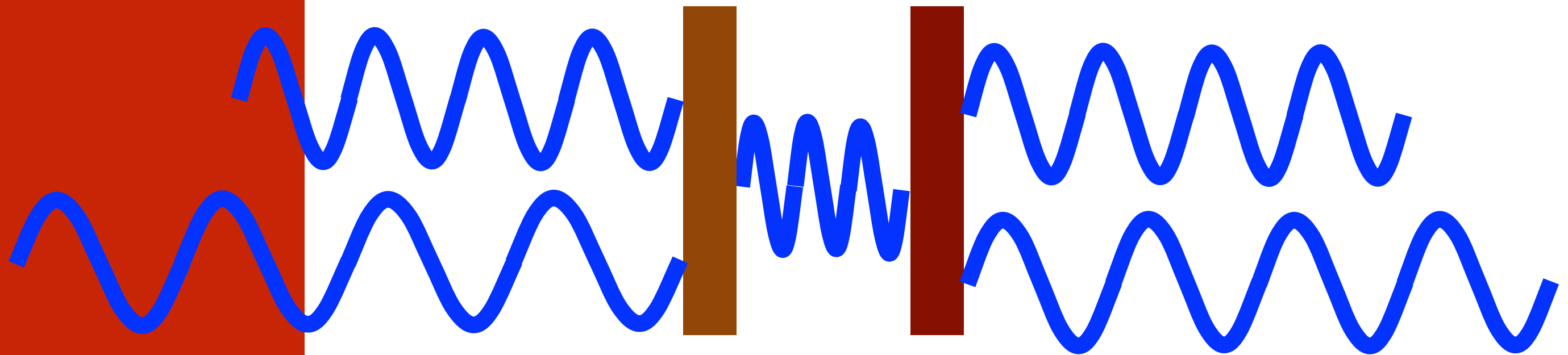
Photos: Copyright © The Nobel Foundation

the “Casimir Effect”

**two highly polished mirrors isolated
from all external effects**

The vacuum has all wavelengths of virtual waves from
particles and fields...but fewer can fit between the walls

...and the pressure from the outside, moves them closer
together



The amount is precisely predicted...and a few years ago the
experiments confirmed it convincingly in 2001

the vacuum

is a very complicated thing

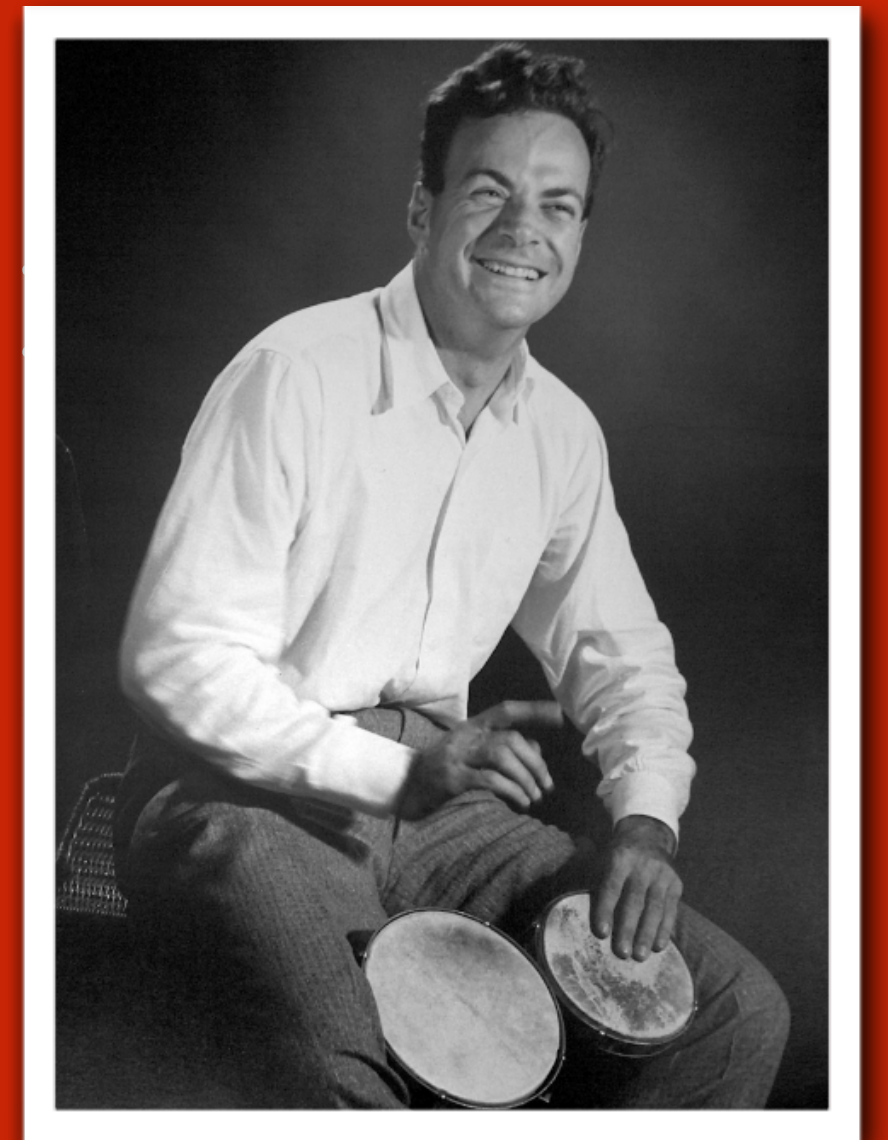
as we'll see when we get back to cosmology



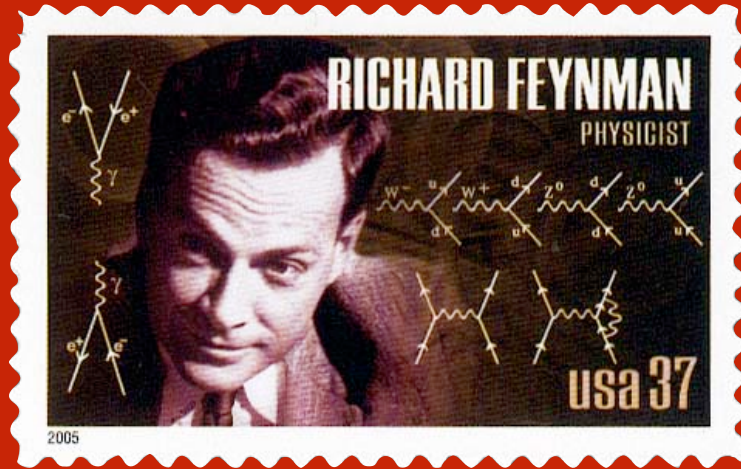
Feynman Diagrams

now for real.

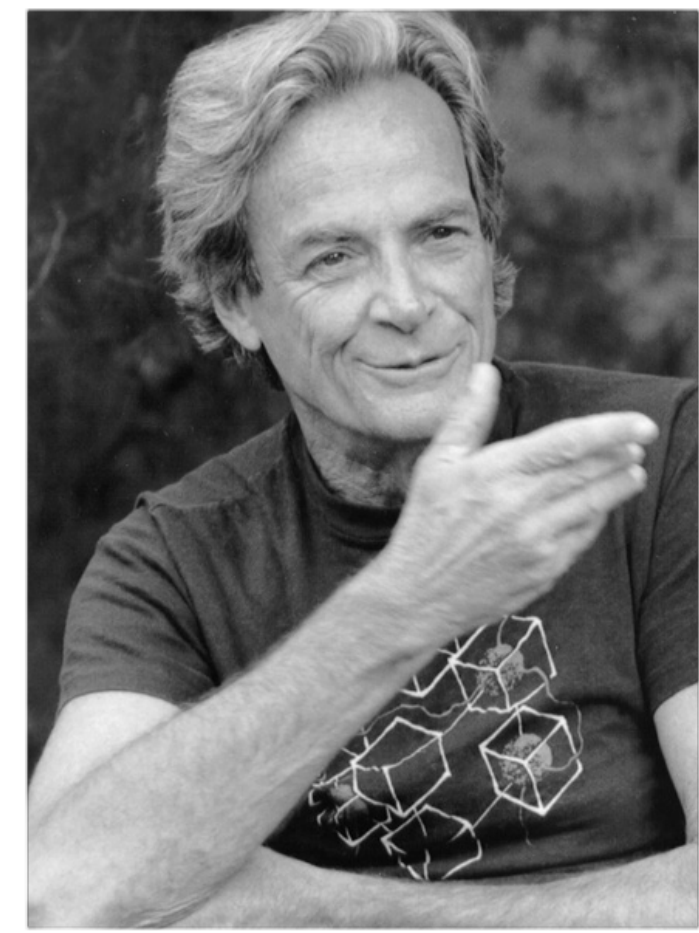
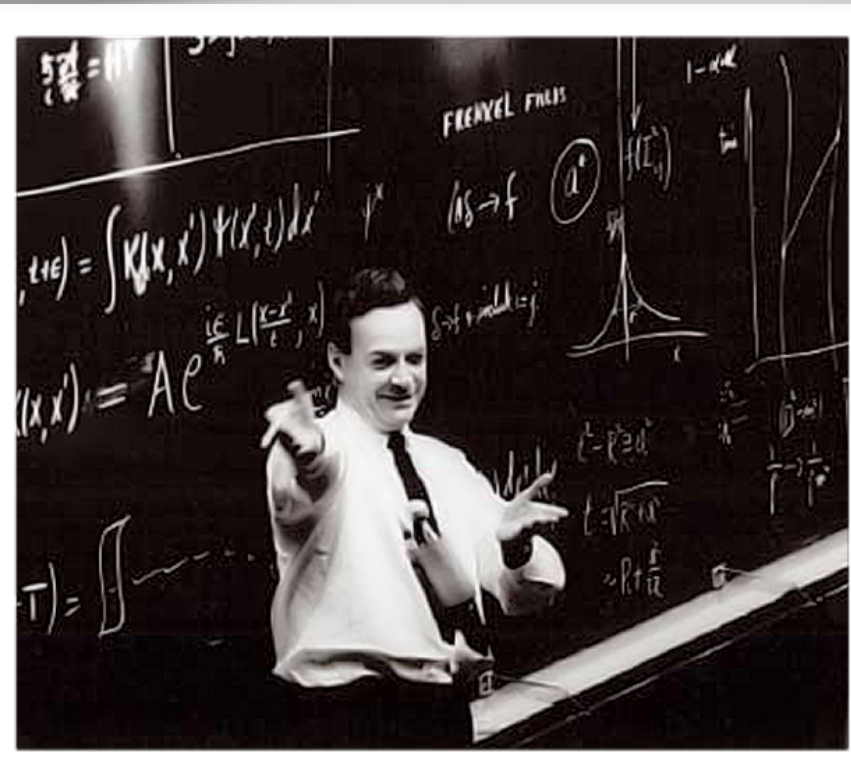
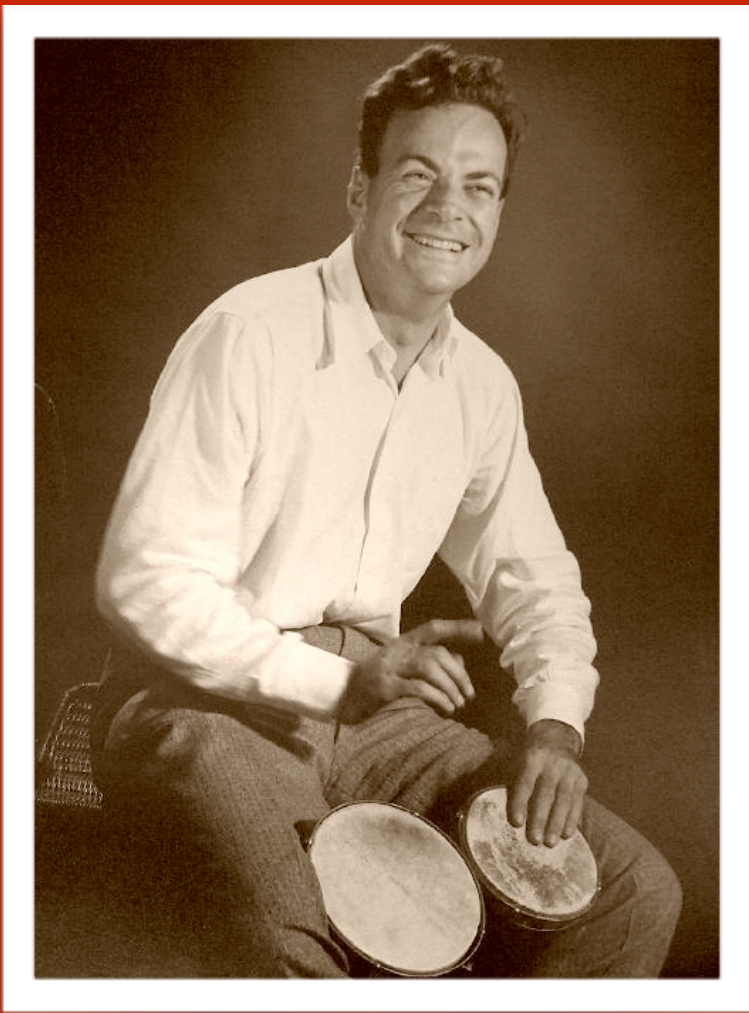
creation and annihilation of
can be embodied in Feynman Diagrams



hero worship



about as close as
we come



Richard Feynman, Sin-Itiro Tomonaga, Julian Schwinger

1965 Nobel



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The Nobel Prize in Physics 1965
Sin-Itiro Tomonaga, Julian Schwinger, Richard P. Feynman

The Nobel Prize in Physics 1965
Sin-Itiro Tomonaga
Julian Schwinger
Richard P. Feynman

Sin-Itiro Tomonaga **Julian Schwinger** **Richard P. Feynman**

The Nobel Prize in Physics 1965 was awarded jointly to Sin-Itiro Tomonaga, Julian Schwinger and Richard P. Feynman *"for their fundamental work in quantum electrodynamics, with deep-ploughing consequences for the physics of elementary particles"*.

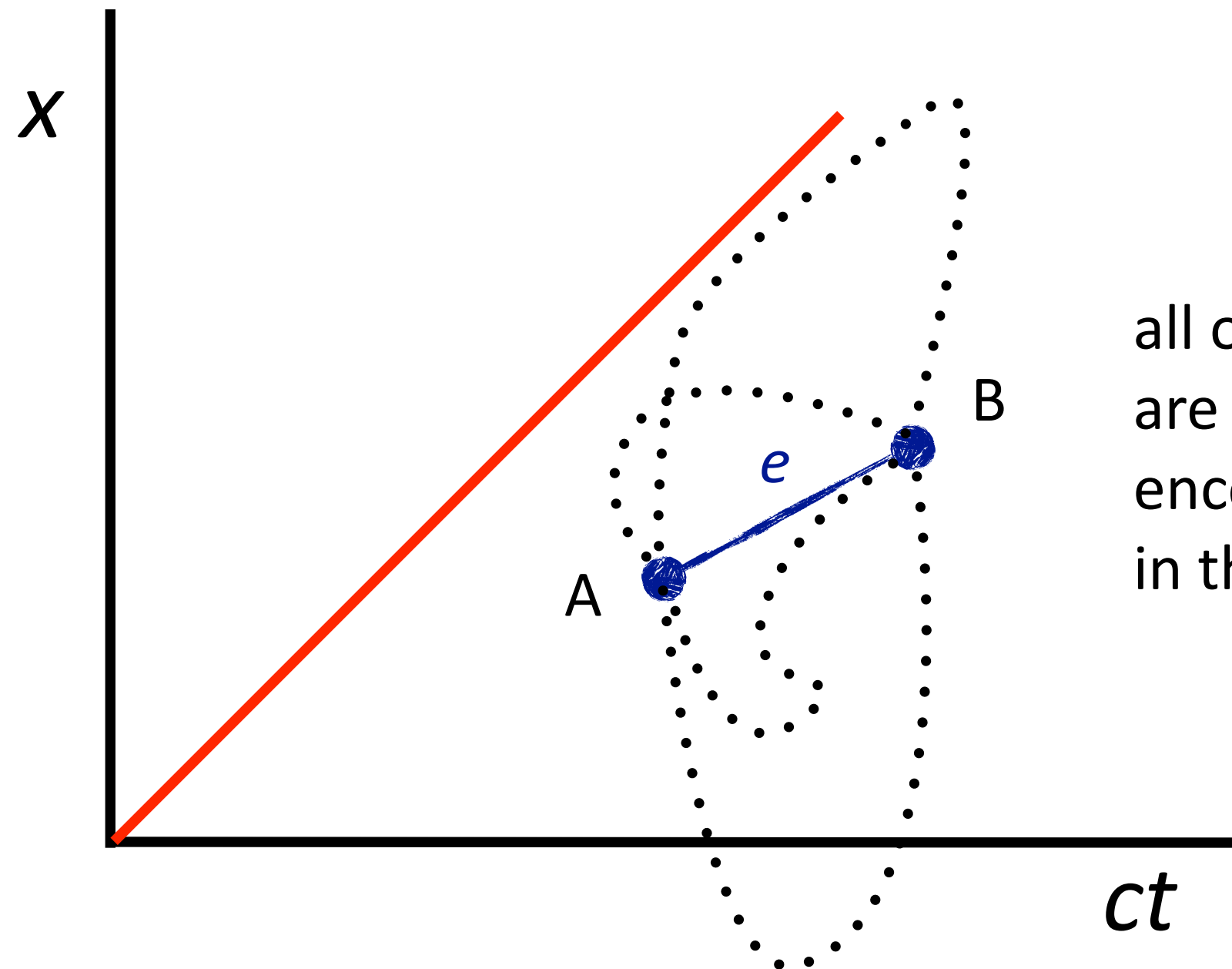
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MLA style: "The Nobel Prize in Physics 1965". Nobelprize.org. 23 Mar 2013
http://www.nobelprize.org/nobel_prizes/physics/laureates/1965/

the symbols of Feynman Diagrams

each line
represents an
entire "history" of
trajectories

to go from A to B, represent all
histories with a single line.



all of the paths
are
encompassed
in the one line

Feynman's lines include rules on how to calculate the
possibilities in a relativistically consistent way.

very
efficient

avoids lots of
technicalities.

When I teach these techniques to second year graduate students, I first do the calculation of Compton Scattering and do it without Feynman's tools.

... spread. (ask look at 204)

$$\langle 0 | T [\psi(x) \bar{\psi}(y)] | 0 \rangle = \bar{\psi}(x) A(x) \psi(y) A(y) A(x) A(y)$$

$$\begin{aligned} (\bar{\psi}^+ + \bar{\psi}^-)(\psi^+ + \psi^-) \\ \bar{\psi}^+ \psi^+ + \bar{\psi}^+ \psi^- + \bar{\psi}^- \psi^+ + \bar{\psi}^- \psi^- \\ \bar{\psi}^+ \psi^+ + \bar{\psi}^- \psi^- = 1 \cdot 0 \cdot 0 + 1 \cdot 0 \cdot 0 \\ \bar{\psi}^+ \psi^- + \bar{\psi}^- \psi^+ = 0 \cdot 0 + 0 \cdot 0 = 0 \end{aligned}$$

... depends on the actual process.

... classical Compton scattering $\epsilon \rightarrow \epsilon'$.

$$\begin{aligned} |initial\rangle &= |e, k\rangle \\ |final\rangle &= |e', k'\rangle \end{aligned}$$

... and $k^0 = -|k^0|$...

$$\frac{1}{k^2 - m^2} \rightarrow \frac{1}{(k^0 - \omega_k)(k^0 + \omega_k)}$$

... and exponential decays.

$$\int_C \frac{d\omega}{\omega} = -2\pi i \sum \text{residues of enclosed poles}$$

$$-2\pi i R = \int_C \frac{d\omega}{\omega} = 2\pi i$$

The pole is simple, so the residue is

$$\frac{1}{(k^0 - \omega_k)(k^0 + \omega_k)} \Big|_{k^0 = \omega_k} = \frac{1}{2\omega_k}$$

$$\langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle$$

$$\begin{aligned} \int d^3x \int d^3y \\ \int d^3x \int d^3y \\ \int d^3x \int d^3y \end{aligned}$$

$$\langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle = \langle 0 | T [\bar{\psi}(x) \psi(x)] | 0 \rangle$$

$$\langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle = \langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle$$

$$\frac{1}{k^2 - m^2}$$

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$$\langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle$$

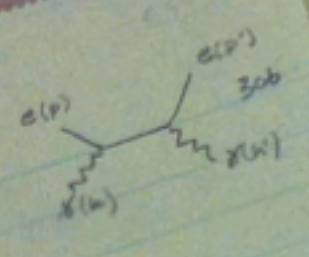
$$\langle 0 | \bar{\psi}(x) \psi(x) | 0 \rangle$$

Compton scattering

$$\begin{aligned}
 &= \frac{(-ie)^2}{2} \int d^4x_1 \int d^4x_2 \langle e^0 | \delta_0^{\mu\nu} \delta_{im} \\
 &\psi_1^+(x_2) A_\nu(x_1) A_\mu(x_2) \psi_2(x_1) \bar{\psi}_2(x_2) \\
 &\psi_1^+(x_1) A_\nu(x_1) A_\mu(x_2) \bar{\psi}_2(x_1) \psi_2(x_2) \rangle |e^0\rangle \\
 &= \int dP_1 a^{(1)}(P_1) u_m^{(1)}(P_1) e^{-iP_1 \cdot x_1} \\
 &= \int dP_2 a^{(2)}(P_2) u_m^{(2)}(P_2) e^{-iP_2 \cdot x_2} \\
 &= \int dK_1 \sum_\lambda \epsilon_{\nu\lambda}(k_1) A_{\lambda\nu}(k_1) e^{-ik_1 \cdot x_1} \\
 &= \int dK_2 \sum_\lambda \epsilon_{\mu\lambda}(k_2) A_{\lambda\mu}(k_2) e^{-ik_2 \cdot x_2}
 \end{aligned}$$

Put all the stuff, into the brackets, separate out the four space terms for do photons.

$$\begin{aligned}
 J^{(1)}(e^0 \rightarrow e^1) &= \frac{(-ie)^2}{2} \int d^4x_1 \int d^4x_2 \langle e^0 | A_\mu \\
 &= \int dP_1 \int dP_2 \langle 0 | a^{(1)}(P_1) a^{(2)}(P_2) a^{(1)\dagger}(P) a^{(2)\dagger}(p) | 0 \rangle \\
 &= \int dP_1 \int dP_2 \langle 0 | a^{(1)}(P) a^{(2)}(p) | 0 \rangle e^{-iP \cdot x_1 - ip \cdot x_2} \\
 &= \int dP_1 \int dP_2 \langle 0 | a^{(1)}(P) a^{(2)}(p) | 0 \rangle e^{-iP \cdot x_1 - ip \cdot x_2}
 \end{aligned}$$



$\psi_a(y)$

Do the standard $\{a(n), a^\dagger(m)\}$ - twice

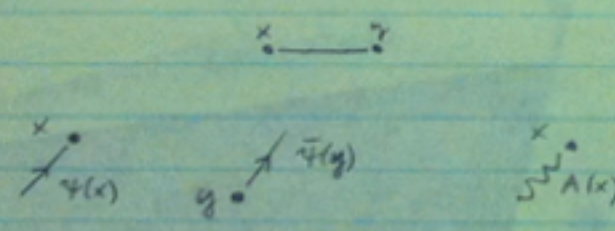
$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2}
 \end{aligned}$$

first

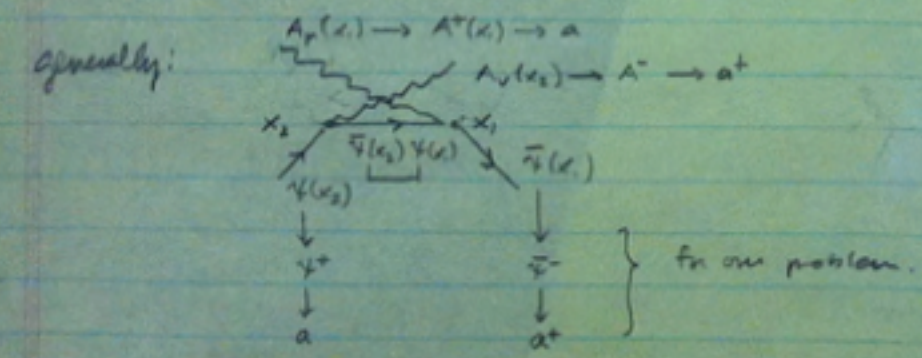
$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2}
 \end{aligned}$$

Earlier, I attached a graphical meaning to the Wick expansion terms, let's recap that according to what we've calculated.

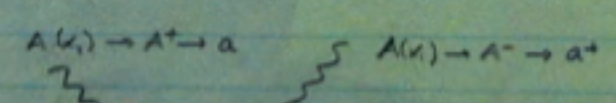
$$\bar{\psi}(x) \psi(y) = \langle 0 | T[\bar{\psi}(x) \psi(y)] | 0 \rangle$$



So, the first (1) or (3) graph would be



and (2) or (4)



- twice (which cancels the funny minus signs) - So we get

$$A_\mu^+(x_1) | \gamma \rangle = \int dK_1 \sum_\lambda \epsilon_{\nu\lambda}(k_1) e^{-ik_1 \cdot x_1} (\pi)^3 2\omega_1 \delta_{\lambda\lambda} \delta(\vec{k}-\vec{k}') | 0 \rangle$$

and we can do the momentum integrals giving

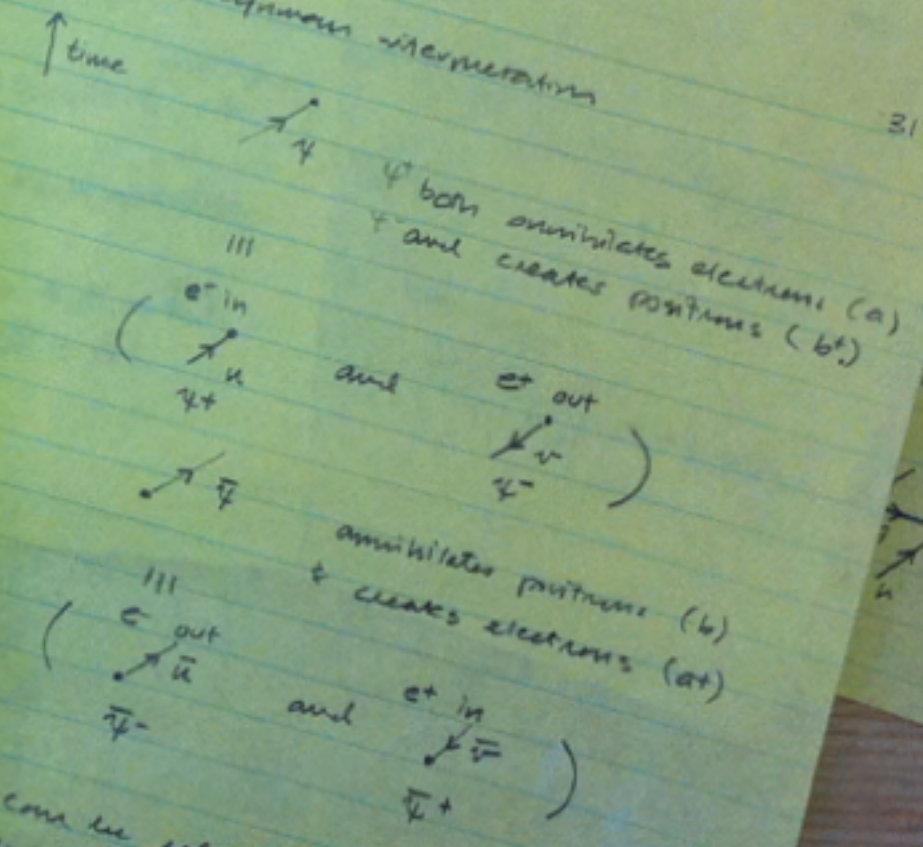
$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2}
 \end{aligned}$$

therefore, 4 terms overall.

$$\langle e^0 | = -e^{\int d^3x_1 \int d^3x_2}$$

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2} \\
 &= \int d^3x_1 \int d^3x_2 \delta(\vec{p}-\vec{p}') \delta_{\lambda_1} \delta_{\lambda_2}
 \end{aligned}$$

Remember the Feynman interpretation



can be related with the frequency indicated.

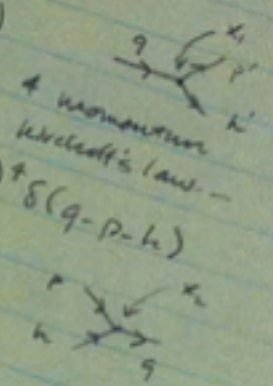
Feynman rules eliminate all ψ and $\bar{\psi}$ and go directly to

$$\int d^4x_1 e^{iP_1 \cdot x_1} e^{-iQ_1 \cdot x_1} e^{iK_1 \cdot x_1}$$

$$\begin{aligned}
 &= (2\pi)^4 \delta(P'-Q+K) \\
 &= (2\pi)^4 \delta(Q-P-K)
 \end{aligned}$$

in (1)

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1}
 \end{aligned}$$



$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1}
 \end{aligned}$$

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1}
 \end{aligned}$$

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
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 \end{aligned}$$

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
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 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1}
 \end{aligned}$$

$$\begin{aligned}
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1} \\
 &= \int d^3x_1 \int d^3x_2 e^{-iP_1 \cdot x_1} e^{iQ_1 \cdot x_1} e^{-iK_1 \cdot x_1}
 \end{aligned}$$

his rules
eliminate
all of
that

and I can just
write down the
"answer"

appropriately labeled, each line tells us what to put into a long equation for further solving

23 Compton - finish! 325

Now let's do the Compton calculation as if we know the rules all along. I want the cross section for $\gamma e \rightarrow \gamma e$ for an initial electron at rest.

$\gamma(k) + e(p) \rightarrow \gamma(k') + e(p')$ to 2nd order.

$(-ie) \bar{u}^{(s)}(p') \epsilon_{\nu\alpha} (-ie\gamma^\alpha) \frac{i}{\not{p} + \not{k} - m} (-ie\gamma^\nu) u^{(s)}(p)$

$(-ie) \bar{u}^{(s)}(p') (-ie\gamma^\nu) \epsilon_{\nu\alpha} \frac{i}{\not{p} - \not{k}' - m} (-ie\gamma^\alpha) u^{(s)}(p)$

so,

$T_{fi} = -e^2 \bar{u}^{(s)}(p') \not{\epsilon}' \frac{1}{\not{p} + \not{k} - m} \not{\epsilon} u^{(s)}(p)$
 $- e^2 \bar{u}^{(s)}(p') \not{\epsilon} \frac{1}{\not{p} - \not{k}' - m} \not{\epsilon}' u^{(s)}(p)$

simplify:

$$\frac{1}{\not{p} + \not{k} - m} = \frac{\not{p} + \not{k} + m}{(\not{p} + \not{k})^2 - m^2} = \frac{\not{p} + \not{k} + m}{p^2 + k^2 + 2p \cdot k - m^2}$$

$$= \frac{\not{p} + \not{k} + m}{2p \cdot k}$$

$$\frac{1}{\not{p} - \not{k}' - m} = \frac{\not{p} - \not{k}' + m}{(\not{p} - \not{k}')^2 - m^2} = \frac{\not{p} - \not{k}' + m}{p^2 + k'^2 - 2p \cdot k' - m^2}$$

24
24

but the pictures themselves are
visually...informative

and I'm going to try to tell you how to do this
without the geeky mathematics

theoretical papers

each diagram is a complicated calculation

m_{π_T} smaller than 260 GeV. The parameter ϵ refers to the proportion of the top quark mass generated by the extended technicolor which is taken in the range of $\epsilon \sim (0.01, 0.1)$.

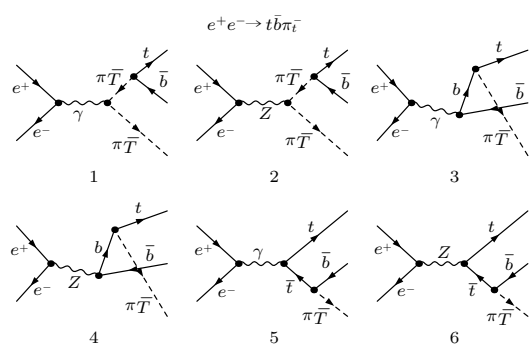


Fig. 1. Diagrams for $e^+e^- \rightarrow t\bar{b}\pi^-$.

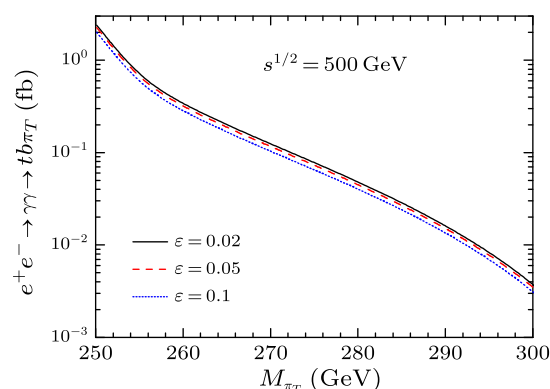


Fig. 2. Dependence of the cross section for $e^+e^- \rightarrow t\bar{b}\pi^-$ ($t\bar{b}\pi^-$) on the top-pion mass m_{π_T} at the ILC with energy of 500 GeV.

By using the laser back-scattering technique on electron beam, an e^+e^- LC which has the c.m. energy of hundreds of GeV to several TeV can be transformed to be a photon collider.^[19–21] By integrating over the photon luminosity in an e^+e^- linear collider, the total cross section for the process $e^+e^- \rightarrow t\bar{b}\pi^-$ can be obtained in the form

$$\sigma(s) = \int_{\frac{E_0}{\sqrt{s}}}^{x_{\max}} dz \frac{d\mathcal{L}_{\gamma\gamma}}{dz} \hat{\sigma}(\gamma\gamma \rightarrow t\bar{b}\pi^-, \text{ at } \hat{s} = z^2s), \quad (12)$$

where $E_0 = m_t + m_b + m_{\pi_T}$, $\sqrt{s}(\sqrt{\hat{s}})$ is the e^+e^- ($\gamma\gamma$) c.m. energy, and $\frac{d\mathcal{L}_{\gamma\gamma}}{dz}$ is the distribution function of photon luminosity, which is defined as

$$\frac{d\mathcal{L}_{\gamma\gamma}}{dz} = 2z \int_{z^2/x_{\max}}^{x_{\max}} \frac{dx}{x} F_{\gamma/e}(x) F_{\gamma/e}\left(\frac{z^2}{x}\right). \quad (13)$$

For the initial unpolarized electrons and laser photon beams, the energy spectrum of the back scattered

photon is given by^[22]

$$F_z = \frac{1}{D(\xi)} \left[1 - x + \frac{1}{1-x} - \frac{4x}{\xi(1-x)} + \frac{4x^2}{\xi^2(1-x)^2} \right], \quad (14)$$

where $x = 2\omega/\sqrt{s}$ is the fraction of the energy of the incident electron carried by the back-scattered photon, the maximum fraction of energy carried by the back-scattered photon is $x_{\max} = 2\omega_{\max}/\sqrt{s} = \xi/(1+\xi)$, and

$$D(\xi) = \left(1 - \frac{4}{\xi} - \frac{8}{\xi^2}\right) \ln(1+\xi) + \frac{1}{2} + \frac{8}{\xi} - \frac{1}{2(1+\xi)^2}, \quad (15)$$

$$\xi = \frac{2\sqrt{s}\omega_0}{m_e^2}, \quad (16)$$

where m_e and $\sqrt{s}/2$ are the mass and energy of the electron, ω_0 is the laser photon energy. In our evaluation, we choose ω_0 such that it maximizes the backscattered photon energy without spoiling the luminosity through e^+e^- pair creation. Then we have $\xi = 2(1+\sqrt{2})$, $x_{\max} \simeq 0.83$, and $D(\xi) \simeq 1.84$, as used in Ref. [23].

The processes $\gamma\gamma \rightarrow t\bar{b}\pi^-$ ($t\bar{b}\pi^-$) occurs through the u- and t-channel involving charged top-pion bremsstrahlungs originated from different positions on quark lines. The Feynman diagrams are drawn in Fig. 3, but the corresponding diagrams with interchange of the two incoming photons are not shown.

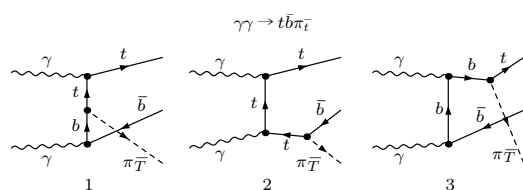


Fig. 3. Diagrams for $\gamma\gamma \rightarrow t\bar{b}\pi^-$.

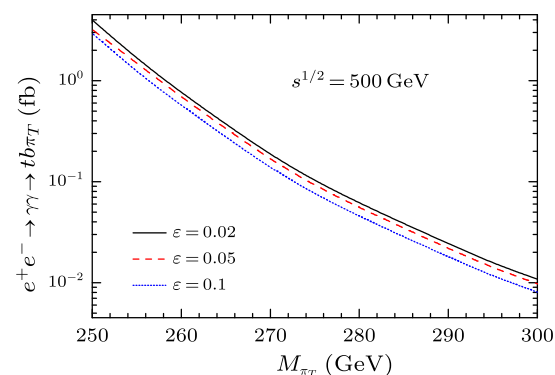


Fig. 4. Dependence of the cross section for $e^+e^- \rightarrow \gamma\gamma \rightarrow t\bar{b}\pi^-$ ($t\bar{b}\pi^-$) on the top-pion mass m_{π_T} at the ILC with energy of 500 GeV.

We show the cross section for $e^+e^- \rightarrow \gamma\gamma \rightarrow t\bar{b}\pi^-$ ($t\bar{b}\pi^-$) at the ILC with energy of 500 GeV as a

2.1.1 Lepton and heavy quark pair decays of the SM Higgs particle

In lowest order the leptonic decay width of the SM Higgs boson is given by [10, 37]

$$\Gamma[H \rightarrow l^+l^-] = \frac{G_F M_H}{4\sqrt{2}\pi} m_l^2 \beta^3 \quad (6)$$

with $\beta = (1 - 4m_l^2/M_H^2)^{1/2}$ being the velocity of the leptons. The branching ratio of decays into τ leptons amounts to about 10% in the intermediate mass range. Muonic decays can reach a level of a few 10^{-4} , and all other leptonic decay modes are phenomenologically unimportant.

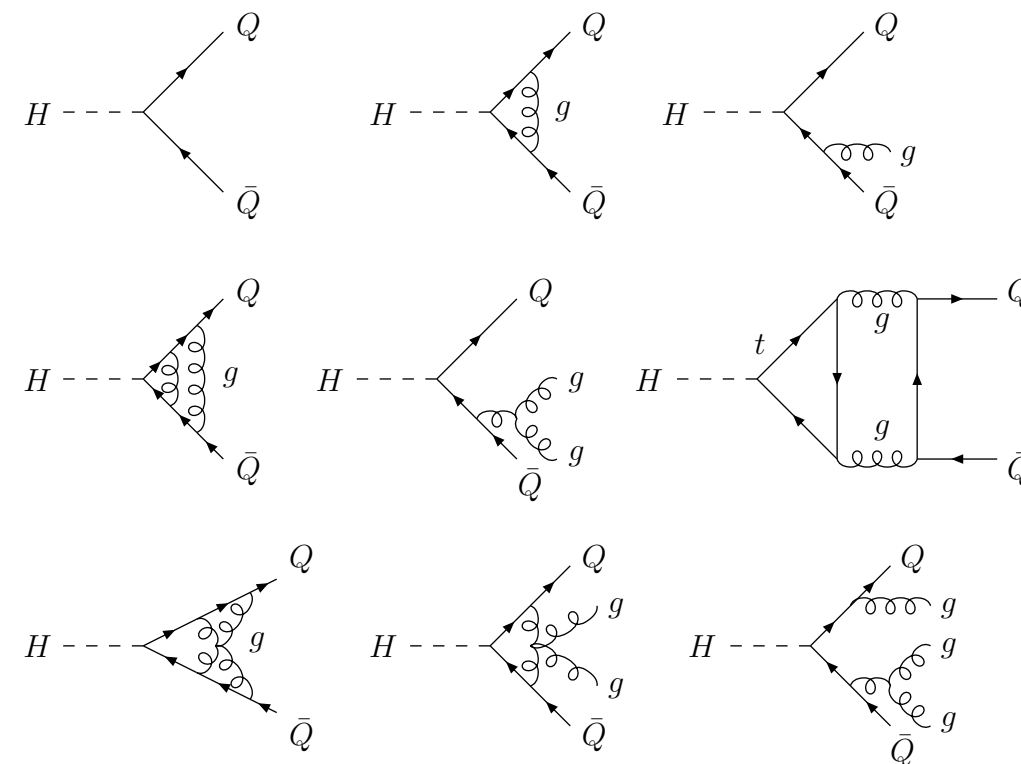
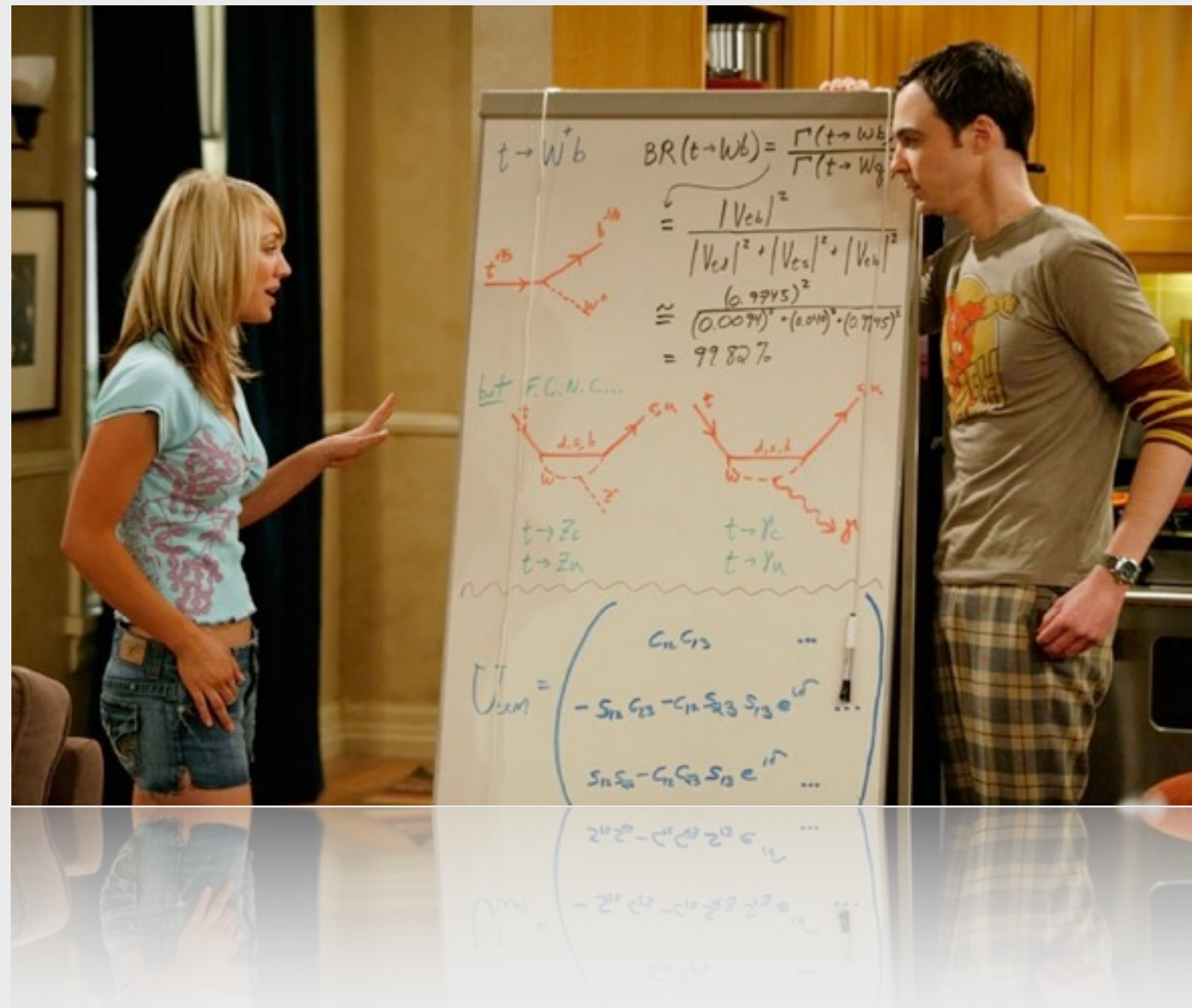


Figure 3: Typical diagrams contributing to $H \rightarrow Q\bar{Q}$ at lowest order and one-, two- and three-loop QCD.

For large Higgs masses the particle width for decays to b, c quarks [directly coupling to the SM Higgs particle] is given up to three-loop QCD corrections [typical diagrams are depicted in Fig. 3] by the well-known expression [38–40]

$$\Gamma[H \rightarrow Q\bar{Q}] = \frac{3G_F M_H}{4\sqrt{2}\pi} \bar{m}_Q^2(M_H) [\Delta_{\text{QCD}} + \Delta_t] \quad (7)$$

we really do use Feynman Diagrams



Feynman's approach is really sneaky and really cute

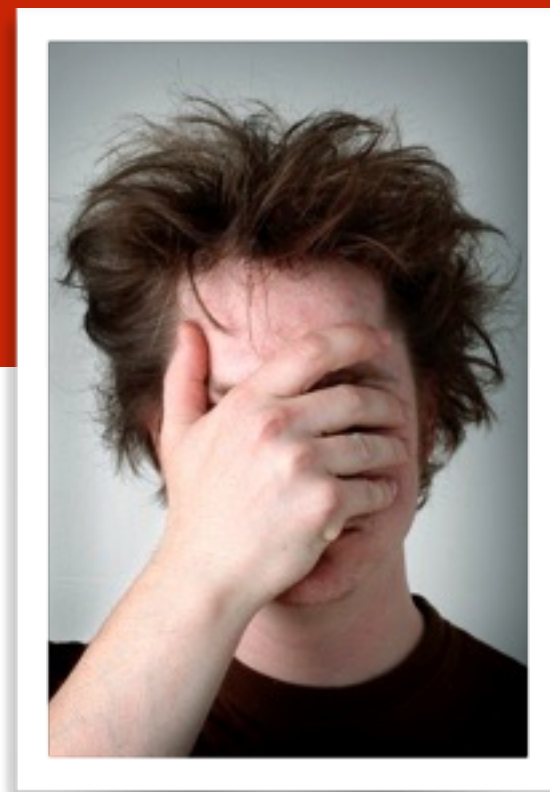
energy and time appear together in the equations:

In essence, this:

either energy solution: $(\pm E)(t)$

just the -E solution: $(-E)(t)$

move the - sign: $(E)(-t)$



Get a whole new interpretation of antimatter